MULTIDIMENSIONAL CONTINUED FRACTIONS, DYNAMICAL RENORMALIZATION AND KAM THEORY

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ABSTRACT. The disadvantage of 'traditional' multidimensional continued fraction algorithms is that it is not known whether they provide simultaneous rational approximations for generic vectors. Following ideas of Dani, Lagarias and Kleinbock-Margulis we describe a simple algorithm based on the dynamics of flows on the homogeneous space $SL(d, \mathbb{Z}) \setminus SL(d, \mathbb{R})$ (the space of lattices of covolume one) that indeed yields best possible approximations to any irrational vector. The algorithm is ideally suited for a number of dynamical applications that involve small divisor problems. We explicitly construct renormalization schemes for (a) the linearization of vector fields on tori of arbitrary dimension and (b) the construction of invariant tori for Hamiltonian systems.

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1. INTRODUCTION

This paper has two main goals. The first one is to introduce a new multidimensional continued fraction algorithm that is ideally suited for different dynamical applications. The algorithm can be used in order to effectively deal with small divisors whenever quasi-periodicity with several frequencies is an essential feature of a problem. Our second goal is to demonstrate the strength of the algorithm by applying it to KAM theory. We use a renormalization approach to prove two theorems of KAM type. The method, being conceptually very simple, is also very general, and allows us to consider a wide class of frequency vectors. For reasons of clarity we restrict our attention to vectors satisfying an explicit Diophantine condition (valid for a set of vectors of full Lebesgue measure); generalizations to more general frequency vectors follow straightforwardly from the approach presented here, and will be detailed in a separate publication.

The classical continued fraction algorithm produces, for every irrational $\alpha \in \mathbb{R}$, a sequence of rational numbers p_n/q_n that approximate α up to an error of order $1/q_n^2$. The first objective of this paper is to develop a multidimensional analogue that allows us to approximate any irrational $\alpha \in \mathbb{R}^{d-1}$ by rational vectors. The drawback of all 'traditional' multidimensional continued fraction algorithms is that it is not even known whether they provide simultaneous rational approximations for almost Lebesgue almost all α . In the case $d \geq 4$ the only result in this direction is a recent computer-assisted proof of the almost everywhere strong convergence for ordered Jacobi-Perron algorithm [10, 11]. However, even in this case an explicit description of the set of bad vectors seems difficult. For example, the existence of 'noble' vectors, that is vectors corresponding to a periodic continued-fraction expansion, for which approximations do not converge, is rather unsatisfactory.

The algorithm we employ here does not suffer from such pathologies. Following Lagarias' seminal ideas in [19], our approach is based on the dynamics of the geodesic¹ flow on the homogeneous space $\Gamma \setminus G$ with $G = \operatorname{SL}(d, \mathbb{R})$ and $\Gamma = \operatorname{SL}(d, \mathbb{Z})$. Notice that $\Gamma \setminus G$ may be identified with the space of lattices in \mathbb{R}^d of covolume one or, equivalently, with the Teichmüller space of flat *d*-dimensional tori.

The problem of multidimensional continued fractions may be formulated in the following way. Given a vector $\boldsymbol{\alpha} \in \mathbb{R}^{d-1}$ find a sequence of matrices $T^{(n)} \in \text{GL}(d, \mathbb{Z})$, $n \in \mathbb{N}$, such that the 'cocycle' corresponding to the products $P^{(n)} = T^{(n)}T^{(n-1)} \dots T^{(1)}$ exponentially contracts in the direction of the vector $\boldsymbol{\omega} = \begin{pmatrix} \boldsymbol{\alpha} \\ 1 \end{pmatrix} \in \mathbb{R}^d$ and exponentially expands in all other directions. Thus, the cocycle should have one negative Lyapunov exponent and d-1 positive. In this spirit, our algorithm comprises the following steps:

(1) With every $\boldsymbol{\alpha} \in \mathbb{R}^{d-1}$ associate the orbit $\{C(t) : t \ge 0\} \subset \mathrm{SL}(d, \mathbb{R})$, where

$$C(t) = \begin{pmatrix} 1_{d-1} & \boldsymbol{\alpha} \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 1_{d-1} e^{-t} & 0 \\ 0 & e^{(d-1)t} \end{pmatrix},$$
(1.1)

¹The term 'geodesic' is slightly inaccurate when d > 2. The orbits of the flow on $\Gamma \backslash G$ we are discussing here correspond in fact to geodesics on the unit cotangent bundle of the space $\Gamma \backslash G / \operatorname{SO}(d)$ for a certain family of initial conditions. Only for d = 2 the cotangent bundle can be identified with $\Gamma \backslash G$.

 1_{d-1} denotes the $(d-1) \times (d-1)$ unit matrix.

- (2) Fix a fundamental domain \mathcal{F} of $\mathrm{SL}(d,\mathbb{Z})$ in $\mathrm{SL}(d,\mathbb{R})$. Given a sequence of times $t_1 < t_2 < \ldots \rightarrow \infty$, use classical reduction theory to find the matrices $P^{(n)} \in \mathrm{SL}(d,\mathbb{Z})$ that map the points $C(t_n)$ to \mathcal{F} .
- (3) Define the nth continued fraction map by

$$\boldsymbol{\alpha}^{(n-1)} \mapsto \boldsymbol{\alpha}^{(n)} = \frac{T_{11}^{(n)} \boldsymbol{\alpha}^{(n-1)} + \boldsymbol{t}_{12}^{(n)}}{{}^{\mathsf{T}} \boldsymbol{t}_{21}^{(n)} \boldsymbol{\alpha}^{(n-1)} + \boldsymbol{t}_{22}^{(n)}},\tag{1.2}$$

where $\boldsymbol{\alpha}^{(0)} = \boldsymbol{\alpha}$ and

$$T^{(n)} = \begin{pmatrix} T_{11}^{(n)} & \boldsymbol{t}_{12}^{(n)} \\ \top \boldsymbol{t}_{21}^{(n)} & t_{22}^{(n)} \end{pmatrix} \in \mathrm{SL}(d, \mathbb{Z})$$
(1.3)

is the *n*th transfer matrix defined by $P^{(n)} = T^{(n)}P^{(n-1)}$.

Since the action of $\mathrm{SL}(d,\mathbb{R})$ on \mathbb{R}^{d-1} by fractional linear transformation defines a group action, we have

$$\boldsymbol{\alpha}^{(n)} = \frac{P_{11}^{(n)} \boldsymbol{\alpha} + \boldsymbol{p}_{12}^{(n)}}{{}^{\mathsf{T}} \boldsymbol{p}_{21}^{(n)} \boldsymbol{\alpha} + p_{22}^{(n)}},\tag{1.4}$$

where

$$P^{(n)} = T^{(n)}T^{(n-1)}\cdots T^{(1)} = \begin{pmatrix} P_{11}^{(n)} & \boldsymbol{p}_{12}^{(n)} \\ {}^{\mathsf{T}}\boldsymbol{p}_{21}^{(n)} & p_{22}^{(n)} \end{pmatrix} \in \mathrm{SL}(d,\mathbb{Z}).$$
(1.5)

Dani [5] and Kleinbock-Margulis [15] observed that Diophantine properties of α translate to divergence properties of the corresponding orbit { $\Gamma C(t) : t \geq 0$ } $\subset \Gamma \setminus G$ in the cusps of $\Gamma \setminus G$. We exploit these results to show that, under mild Diophantine conditions on α (satisfied by a set of α of full Lebesgue measure, cf. Sec. 2.6), there is a sequence of times t_n such that the transfer matrices $T^{(n)}$ are uniformly hyperbolic in a sense made precise in Sec. 2.8. This fact allows us to develop renormalization schemes for vector fields and Hamiltonian flows that had previously been constructed only in dimension one [21] or for very special choices of α [1, 16, 17, 20]. In particular, we obtain renormalization-based proofs of the following theorems.

We say $\boldsymbol{\omega} \in \mathbb{R}^d$ is *Diophantine* if there are constants $\epsilon > 0, C > 0$ such that

$$\|\boldsymbol{k}\|^{(d-1)(1+\epsilon)}|\boldsymbol{k}\cdot\boldsymbol{\omega}| > C, \tag{1.6}$$

for all $\mathbf{k} \in \mathbb{Z}^d - \{\mathbf{0}\}$. Note that we may assume without loss of generality that $\boldsymbol{\omega}$ is of the form $\boldsymbol{\omega} = \begin{pmatrix} \alpha \\ 1 \end{pmatrix}$ with $\boldsymbol{\alpha} \in \mathbb{R}^{d-1}$. Condition (1.6) then translates to a standard Diophantine condition on $\boldsymbol{\alpha}$, see Sec. (2.6) for details.

Theorem 1.1. For any real analytic vector field \boldsymbol{v} on \mathbb{T}^d , $d \geq 2$, sufficiently close to a constant vector field with Diophantine vector $\boldsymbol{\omega} \in \mathbb{R}^d$, there is $\epsilon > 0$, an analytic curve $(-\epsilon, \epsilon) \ni s \mapsto \boldsymbol{p}^s \in \mathbb{R}^d$ and an analytic conjugacy h isotopic to the identity between the flow generated by $\boldsymbol{v} + \boldsymbol{p}^s$ and the linear flow $\phi_t(\boldsymbol{x}) = \boldsymbol{x} + t(1+s)\boldsymbol{\omega}$ on \mathbb{T}^d , $t \geq 0$, for each $|s| < \epsilon$. Moreover, the maps $\boldsymbol{v} \mapsto h$ and $\boldsymbol{v} \mapsto \boldsymbol{p}$ are analytic.

Let us emphasize that the result holds for all analytic vector fields close to a constant one without any additional conditions, such as preservation of volume etc. The second theorem deals with the case of Hamiltonian vector fields. Let $B \subset \mathbb{R}^d$, $d \geq 2$, be an open set containing the origin, and let H^0 be a real-analytic Hamiltonian function

$$H^{0}(\boldsymbol{x},\boldsymbol{y}) = \boldsymbol{\omega} \cdot \boldsymbol{y} + \frac{1}{2} {}^{\mathsf{T}} \boldsymbol{y} Q \boldsymbol{y}, \quad (\boldsymbol{x},\boldsymbol{y}) \in \mathbb{T}^{d} \times B,$$
(1.7)

with $\boldsymbol{\omega} \in \mathbb{R}^d$ and a real symmetric $d \times d$ matrix Q. It is said to be non-degenerate if $\det Q \neq 0$.

Theorem 1.2. Suppose H^0 is non-degenerate and $\boldsymbol{\omega}$ is Diophantine. If H is a real analytic Hamiltonian on $\mathbb{T}^d \times B$ sufficiently close to H^0 , then the Hamiltonian flow of H leaves invariant a Lagrangian d-dimensional torus where it is analytically conjugated to the linear flow $\phi_t(\boldsymbol{x}) = \boldsymbol{x} + t\boldsymbol{\omega}$ on \mathbb{T}^d , $t \geq 0$. The conjugacy depends analytically on H.

Outline of the proof of Theorem 1.1. Consider a vector field $X(\boldsymbol{x}, \boldsymbol{y}) = \boldsymbol{\omega} + f(\boldsymbol{x}) + \boldsymbol{y}$ where $\boldsymbol{x} \in \mathbb{T}^d$, and $\boldsymbol{y} \in \mathbb{R}^d$ is an auxiliary parameter. The vector field $f(\boldsymbol{x})$ is a sufficiently small analytic perturbation of a constant vector field. We may furthermore assume that $\boldsymbol{\omega} = \begin{pmatrix} \boldsymbol{\alpha} \\ 1 \end{pmatrix}$ for some Diophantine $\boldsymbol{\alpha} \in \mathbb{R}^{d-1}$; this achieved by a rescaling of time. The aim is to find a value of parameter $\boldsymbol{y} = \boldsymbol{y}_{\boldsymbol{\omega}}$ such that the vector field $X(\boldsymbol{x}, \boldsymbol{y}_{\boldsymbol{\omega}})$ is linearizable to a constant vector field identically equal to $\boldsymbol{\omega}$ by means of an analytic transformation of the coordinates on \mathbb{T}^d .

Renormalization is an iterative process, and we thus assume that after the (n-1)th renormalization step the vector field is of the slightly more general form

$$X_{n-1}(\boldsymbol{x}, \boldsymbol{y}) = \boldsymbol{\omega}^{(n-1)} + f_{n-1}(\boldsymbol{x}, \boldsymbol{y})$$
(1.8)

where $\boldsymbol{\omega}^{(n-1)} = (\boldsymbol{\alpha}_{1}^{(n-1)})$ and $\boldsymbol{\alpha}^{(n-1)}$ is given by the continued fraction algorithm, cf. (1.4). The Fourier modes of f_{n-1} are smaller than in the previous step, and decay exponentially as $\|\boldsymbol{k}\| \to +\infty$. We define a cone of resonant modes by a relation $I_{n-1}^{+} = \{\boldsymbol{k} \in \mathbb{Z}^d : |\boldsymbol{k} \cdot \boldsymbol{\omega}^{(n-1)}| \leq \sigma_{n-1} \|\boldsymbol{k}\|\}$. The *n*th step requires the following operations.

- (1) Eliminate all Fourier modes outside of the resonant cone I_{n-1}^+ .
- (2) Apply a linear operator corresponding to a coordinate transformation given by the inverse transfer matrix $T^{(n)^{-1}}$.
- (3) Rescale time to ensure that the frequency vector is of the form $\boldsymbol{\omega}^{(n)} = (\boldsymbol{\alpha}_1^{(n)})$.

The conjugate action on the Fourier modes is given by $\mathbf{k} \mapsto {}^{\mathsf{T}}T^{(n)^{-1}}\mathbf{k}$. It follows from the hyperbolicity of $T^{(n)}$ that this transformation contracts for $\mathbf{k} \in I_{n-1}^+$ if σ_{n-1} is small enough. This gives a significant improvement of the analyticity domain which results in the decrease of the estimates for the corresponding Fourier modes. As a result, all Fourier modes apart from the zero modes get smaller. To decrease the size of the latter, we choose a parameter $\mathbf{y} = \mathbf{y}_n$ in such a way that the corresponding zero modes vanish, and then consider a neighbourhood of \mathbf{y} -values centred at \mathbf{y}_n . That is, the auxiliary parameter \mathbf{y} is used to eliminate an instability in the direction of constant vector fields. To get enough control on the parameter dependence we perform an affine rescaling of this parameter on every renormalization step. One can then show that the corresponding sequence of parameter domains is nested and converges to a single point $\mathbf{y} = \mathbf{y}_{\boldsymbol{\omega}}$ for which the initial vector field is indeed linearizable.

In order for the scheme to be effective, the sequence of stopping times t_n and the sizes of the resonant cones defined by the sequence of σ_n must be chosen properly. Large intervals $\delta t_n = t_n - t_{n-1}$ improve hyperbolicity but, on the other hand, worsen estimates for the norms $||T^{(n)}||$, $||T^{(n)^{-1}}||$. Similarly, if σ_n is too small the elimination of non-resonant modes will give large contributions; on the other hand, for large values of σ_n the multiplication by ${}^{\mathsf{T}}T^{(n+1)^{-1}}$ will not yield a contraction for $\mathbf{k} \in I_n^+$. As we shall show below a right choice of sequences $\{(t_n, \sigma_n)\}$ can be made, depending on the Diophantine properties of the vector $\boldsymbol{\omega}$.

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The idea of renormalization was introduced to the theory of dynamical systems by Feigenbaum [6] in the late 1970's. In the case of Hamiltonian systems with two degrees of freedom MacKay proposed in the early 1980's a renormalization scheme for the construction of KAM invariant tori [22]. The scheme was realized for the construction of invariant curves for two-dimensional conservative maps of the cylinder. An important feature of MacKay's approach is the analysis of both smooth KAM invariant curves and so-called critical curves corresponding to critical values of a parameter above which invariant curves no longer exist. From the point of view of renormalization theory the KAM curves correspond to a trivial linear fixed point for the renormalization transformations, while critical curves give rise to very complicated fixed points with nontrivial critical behavior. MacKay's renormalization scheme was carried out only for a small class of Diophantine rotation numbers with periodic continued fraction expansion (such as the golden mean). Khanin and Sinai studied a different renormalization scheme for general Diophantine rotation numbers [14]. Both of the above early approaches were based on renormalization for maps or their generating functions. Essentially, the renormalization transformations are defined in the space of pairs of mappings which, being iterates of the same map, commute with each other. These commutativity conditions cause difficult technical problems, and led MacKay [23] to propose the development of alternative renormalization schemes acting directly on vector fields. The same idea was realized by Koch [16] who proves a KAM type result for analytic perturbations of linear Hamiltonians $H^0(\boldsymbol{x}, \boldsymbol{y}) = \boldsymbol{\omega} \cdot \boldsymbol{y}$, for frequencies $\boldsymbol{\omega}$ which are eigenvectors of hyperbolic matrices in $SL(d, \mathbb{Z})$ with only one unstable direction. Notice that the set of such frequencies has zero Lebesgue measure and in the case d = 2 corresponds to vectors with a quadratic irrational slope. Further improvements and applications of Koch's techniques appeared in [1, 17, 20, 21, 7], emphasizing the connection between KAM and renormalization theories. The results of this paper illustrate that such a programme can indeed be carried out in considerable generality. Another direction was followed in [18], presenting a computer-assisted proof of the existence of MacKay's golden mean critical renormalization fixed point in the context of Hamiltonian vector fields with two degrees of freedom.

Other renormalization ideas have appeared in the context of the stability of invariant tori for nearly integrable Hamiltonian systems (see e.g. [3, 8, 9] and references therein). Inspired by quantum field theory and an analogy with KAM theory, this approach uses an iterative resummation of Poincaré's Lindstedt series to prove its convergence.

For the sake of transparency we have restricted our attention to Diophantine vectors $\boldsymbol{\omega}$. A more detailed analysis under weaker Diophantine conditions is, in principle, possible within the present framework. It is however a fundamental open problem to state a sharp (i.e. the weakest possible) Diophantine condition under which the above conjugacy can be established. The answer to this question is known only in the classical case d = 2 [26] where the Diophantine condition is of Brjuno type.

It would also be interesting to see whether the multidimensional continued fraction algorithm presented here will allow generalizations of other one-dimensional renormalization constructions. A concrete challenge is for instance the extension of the recent results on the reducibility of cocycles over irrational rotations by Ávila and Krikorian [2].

In the next section we introduce the multidimensional continued fraction algorithm, and include a discussion of its hyperbolicity properties required in the renormalization schemes. The remainder of the paper provides a detailed account of two exemplary cases, the renormalization of vector fields (Section 3) and Hamiltonian flows (Section 4).

2. Multidimensional continued fractions and flows on homogeneous spaces

2.1. Flows on homogeneous spaces. Let us set $G = SL(d, \mathbb{R})$ and $\Gamma = SL(d, \mathbb{Z})$, and define the diagonal subgroup $\{E^t : t \in \mathbb{R}\}$ in G, where

$$E^t = \operatorname{diag}(e^{r_1 t}, \dots, e^{r_d t}) \tag{2.1}$$

with constants satisfying the conditions

$$r_1, \dots, r_{d-1} < 0 < r_d, \qquad \sum_{j=1}^a r_j = 0.$$
 (2.2)

The right action of E^t on the homogeneous space $\Gamma \setminus G \Gamma \setminus G$ generates the flow

$$\Phi^t: \Gamma \backslash G \to \Gamma \backslash G, \qquad \Gamma M \mapsto \Gamma M E^t.$$
(2.3)

Since G is a simple Lie group with finite center, Φ^t is ergodic and mixing [24].

Let $\mathcal{F} \subset G$ be a fundamental domain of the left action of Γ on G. Recall that, by definition of the fundamental domain of a free group action,

$$\bigcup_{P \in \Gamma} P\mathcal{F} = G, \qquad \mathcal{F} \cap P\mathcal{F} = \emptyset \text{ for all } P \in \Gamma - \{1\},$$
(2.4)

and hence, for any given $M \in G$, there is a unique family of $P(t) \in \Gamma$ such that

$$M(t) := P(t)ME^t \in \mathcal{F}$$
(2.5)

holds for all $t \in \mathbb{R}$.

2.2. A convenient parametrization. Let us consider those $M \in G$ which can be written as

$$M = \begin{pmatrix} 1 & \boldsymbol{\alpha} \\ 0 & 1 \end{pmatrix} \begin{pmatrix} A & \mathbf{0} \\ \top \boldsymbol{\beta} & \boldsymbol{\gamma} \end{pmatrix}$$
(2.6)

where $A \in \operatorname{Mat}_{d-1}(\mathbb{R})$ (the space of real $(d-1) \times (d-1)$ matrices), $\alpha, \beta \in \mathbb{R}^{d-1}$ are column vectors, $\gamma \in \mathbb{R}$ with $\gamma > 0$. This yields a local parametrization of G for the set

$$G_{+} := \left\{ \begin{pmatrix} 1 & \boldsymbol{\alpha} \\ 0 & 1 \end{pmatrix} \begin{pmatrix} A & \boldsymbol{0} \\ {}^{\mathsf{T}}\boldsymbol{\beta} & \gamma \end{pmatrix} \in G : A \in \operatorname{Mat}_{d-1}(\mathbb{R}), \, \boldsymbol{\alpha}, \boldsymbol{\beta} \in \mathbb{R}^{d-1}, \, \gamma \in \mathbb{R}_{>0} \right\}, \quad (2.7)$$

which is particularly convenient for our purposes. All other matrices are either of the above form with $\gamma < 0$ instead, or may be written as

$$M = S \begin{pmatrix} 1 & \boldsymbol{\alpha} \\ 0 & 1 \end{pmatrix} \begin{pmatrix} A & \mathbf{0} \\ \top \boldsymbol{\beta} & \gamma \end{pmatrix}$$
(2.8)

where $S \in \Gamma$ is a suitably chosen "signed permutation matrix", i.e., every row and every column contains one and only one non-zero coefficient, which is either 1 or -1. In the following we will stay clear of the parameter singularity at $\gamma = 0$, and thus may assume without loss of generality S = 1.

To work out the action of a general element $T \in G$ in the above parametrization, consider

$$T: M \mapsto \tilde{M} := TM \tag{2.9}$$

where

$$T = \begin{pmatrix} T_{11} & \boldsymbol{t}_{12} \\ {}^{\mathsf{T}}\boldsymbol{t}_{21} & \boldsymbol{t}_{22} \end{pmatrix}, \qquad (2.10)$$

M is as above and

$$\tilde{M} = \begin{pmatrix} 1 & \tilde{\boldsymbol{\alpha}} \\ 0 & 1 \end{pmatrix} \begin{pmatrix} \tilde{A} & \boldsymbol{0} \\ \top \boldsymbol{\beta} & \boldsymbol{\tilde{\gamma}} \end{pmatrix}.$$
(2.11)

A short calculation yields the fractional linear action

$$\boldsymbol{\alpha} \mapsto \tilde{\boldsymbol{\alpha}} = \frac{T_{11}\boldsymbol{\alpha} + \boldsymbol{t}_{12}}{\mathsf{T}\boldsymbol{t}_{21}\boldsymbol{\alpha} + \boldsymbol{t}_{22}},\tag{2.12}$$

and

$$\gamma \mapsto \tilde{\gamma} = ({}^{\mathsf{T}} \boldsymbol{t}_{21} \boldsymbol{\alpha} + t_{22}) \gamma, \qquad (2.13)$$

and more complicated expressions for $\tilde{A}, \tilde{\beta}$ which will not be needed in the following.

2.3. Multidimensional continued fractions. Let $t_0 = 0 < t_1 < t_2 < \ldots \rightarrow \infty$ be sequence of times, with gaps

$$\delta t_n := t_n - t_{n-1} \tag{2.14}$$

chosen large enough so that $P(t_n) \neq P(t_{n-1})$, where P(t) is defined by (2.5). The sequence $P^{(n)} := P(t_n)$ of matrices in Γ may be viewed as the continued fraction approximants of the vector $\boldsymbol{\alpha}$, which are the best possible for suitable choices of a fundamental domain \mathcal{F} and times t_n , see [19]. Let us furthermore put $M^{(n)} := M(t_n)$ with M(t) as in (2.5), and define $\boldsymbol{\alpha}^{(n)}, \gamma^{(n)}$ by the decomposition (2.6), i.e., by

$$M^{(n)} = \begin{pmatrix} 1 & \boldsymbol{\alpha}^{(n)} \\ 0 & 1 \end{pmatrix} \begin{pmatrix} A^{(n)} & \mathbf{0} \\ {}^{\mathsf{T}}\boldsymbol{\beta}^{(n)} & \gamma^{(n)} \end{pmatrix}.$$
 (2.15)

From $M^{(n)} = P^{(n)}ME^{t_n}$ and (2.12), (2.13) we deduce

$$\boldsymbol{\alpha}^{(n)} = \frac{P_{11}^{(n)} \boldsymbol{\alpha} + \boldsymbol{p}_{12}^{(n)}}{{}^{\mathsf{T}} \boldsymbol{p}_{21}^{(n)} \boldsymbol{\alpha} + p_{22}^{(n)}},$$
(2.16)

and

$$\gamma^{(n)} = \left({}^{\mathsf{T}}\boldsymbol{p}_{21}^{(n)}\boldsymbol{\alpha} + p_{22}^{(n)} \right) \mathrm{e}^{r_d t_n} \gamma \tag{2.17}$$

where

$$P^{(n)} = \begin{pmatrix} P_{11}^{(n)} & \boldsymbol{p}_{12}^{(n)} \\ \top \boldsymbol{p}_{21}^{(n)} & p_{22}^{(n)} \end{pmatrix}.$$
 (2.18)

It is evident that if the components of $({}^{\top}\alpha, 1)$ are linearly independent over \mathbb{Q} , then $\gamma \neq 0$ implies $\gamma^{(n)} \neq 0$ for all $n \geq 0$.

We shall later employ the transfer matrices $T^{(n)}$ defined by $P^{(n)} = T^{(n)}P^{(n-1)}$. Here, $M^{(n)} = T^{(n)}M^{(n-1)}E^{\delta t_n}$ implies

$$\boldsymbol{\alpha}^{(n)} = \frac{T_{11}^{(n)} \boldsymbol{\alpha}^{(n-1)} + \boldsymbol{t}_{12}^{(n)}}{{}^{\mathsf{T}} \boldsymbol{t}_{21}^{(n)} \boldsymbol{\alpha}^{(n-1)} + \boldsymbol{t}_{22}^{(n)}},\tag{2.19}$$

and

$$\gamma^{(n)} = \left({}^{\mathsf{T}} \boldsymbol{t}_{21}^{(n)} \boldsymbol{\alpha}^{(n-1)} + t_{22}^{(n)} \right) \mathrm{e}^{r_d \delta t_n} \gamma^{(n-1)}, \qquad (2.20)$$

where

$$T^{(n)} = \begin{pmatrix} T_{11}^{(n)} & \boldsymbol{t}_{12}^{(n)} \\ \top \boldsymbol{t}_{21}^{(n)} & \boldsymbol{t}_{22}^{(n)} \end{pmatrix}.$$
 (2.21)

2.4. Siegel sets. In dimensions d > 2 it is difficult to describe the geometry of a fundamental domain \mathcal{F} . To overcome this problem, C. Siegel introduced simply connected sets $\mathcal{S}_d \subset G$ which have the property that they contain \mathcal{F} and are contained in a finite number of translates $P\mathcal{F}$, $P \in \Gamma$. Consider the Iwasawa decomposition

$$M = nak \tag{2.22}$$

where

$$n = \begin{pmatrix} 1 & u_{12} & \dots & u_{1d} \\ & \ddots & \ddots & \vdots \\ & & \ddots & u_{d-1,d} \\ & & & 1 \end{pmatrix}, \qquad a = \begin{pmatrix} a_1 & & & \\ & \ddots & & \\ & & \ddots & \\ & & & a_d \end{pmatrix}$$
(2.23)

and $k \in SO(d)$, with $u_{ij}, a_j \in \mathbb{R}, a_j > 0, a_1 \cdots a_d = 1$. Then

$$\mathcal{S}_d = \{ nak : \ n \in \mathcal{F}_N, \ a_j \ge \frac{\sqrt{3}}{2} a_{j+1} > 0 \ (j = 1, \dots, d-1), \ k \in \mathrm{SO}(d) \}$$
(2.24)

is an example of a Siegel set [25]; here \mathcal{F}_N denotes a compact fundamental region of $(\Gamma \cap N) \setminus N$, where N is the upper triangular group of elements of the form n as above.

2.5. Dani's correspondence. We assume from now on that $r_1, \ldots, r_{d-1} = -1$, $r_d = d - 1$, i.e.,

$$E^{t} = \operatorname{diag}(e^{-t}, \dots, e^{-t}, e^{(d-1)t}).$$
 (2.25)

Let us denote by $|\cdot|$ the maximum norm in \mathbb{R}^{d-1} . A vector $\boldsymbol{\alpha} \in \mathbb{R}^{d-1}$ is called *badly* approximable or of *bounded type*, if one of the following equivalent conditions is satisfied.

(i) There exists a constant C > 0 such that

$$|k\boldsymbol{\alpha} + \boldsymbol{m}|^{d-1}|k| > C \tag{2.26}$$

for all $\boldsymbol{m} \in \mathbb{Z}^{d-1}$, $k \in \mathbb{Z} - \{0\}$.

(ii) There exists a constant C > 0 such that

$$|\boldsymbol{m}|^{d-1}|\boldsymbol{m}\cdot\boldsymbol{\alpha}+k| > C \tag{2.27}$$

for all $\boldsymbol{m} \in \mathbb{Z}^{d-1} - \{\boldsymbol{0}\}, k \in \mathbb{Z}$.

The statements (i) and (ii) are equivalent in view of Khintchine's transference principle ([4] Chapter V).

We recall *Dani's correspondence* in the following proposition (cf. [5], Theorem 2.20).

Proposition 2.1. The orbit $\{\Gamma M E^t : t \ge 0\}$, with M as in (2.6), is bounded in $\Gamma \setminus G$ if and only if the vector $\boldsymbol{\alpha}$ is of bounded type.

The reason why the parameters A, β, γ are irrelevant in the statement is that the family of matrices

$$W(t) = E^{-t} \begin{pmatrix} A & \mathbf{0} \\ {}^{\mathsf{T}}\boldsymbol{\beta} & \boldsymbol{\gamma} \end{pmatrix} E^{t}$$
(2.28)

is bounded in G for all $t \ge 0$.

The boundedness of the orbit { $\Gamma M E^t : t \ge 0$ } implies of course that there is a compact set $\mathcal{C} \in G$ such that $M(t) \in \mathcal{C}$ for all $t \ge 0$, with M(t) as in (2.5).

2.6. Diophantine conditions. A vector $\boldsymbol{\alpha} \in \mathbb{R}^{d-1}$ is called *Diophantine*, if there exist constants $\epsilon > 0, C > 0$ such that

$$|\boldsymbol{m}|^{(d-1)(1+\epsilon)}|\boldsymbol{m}\cdot\boldsymbol{\alpha}+k| > C$$
(2.29)

for all $m \in \mathbb{Z}^{d-1} - \{0\}$, $k \in \mathbb{Z}$. It is well known that Diophantine vectors form a set of full Lebesgue measure [4].

Let us show that (2.29) implies the inequality

$$\|\boldsymbol{k}\|^{(d-1)(1+\epsilon)}|\boldsymbol{k}\cdot\boldsymbol{\omega}| > C, \qquad (2.30)$$

for all $\mathbf{k} \in \mathbb{Z}^d - \{\mathbf{0}\}$, where $\boldsymbol{\omega} = \begin{pmatrix} \boldsymbol{\alpha} \\ 1 \end{pmatrix}$, cf. (1.6). With $\mathbf{k} = \begin{pmatrix} \boldsymbol{m} \\ k \end{pmatrix}$, (2.29) yields

$$|\boldsymbol{k} \cdot \boldsymbol{\omega}| > C |\boldsymbol{m}|^{-(d-1)(1+\epsilon)} \ge C |\boldsymbol{k}|^{-(d-1)(1+\epsilon)} \ge C ||\boldsymbol{k}||^{-(d-1)(1+\epsilon)}$$
(2.31)

for all $\mathbf{m} \in \mathbb{Z}^{d-1} - \{\mathbf{0}\}$, $k \in \mathbb{Z}$. In the case when $\mathbf{m} = \mathbf{0}$, we have $k \neq 0$ (since $\mathbf{k} \neq \mathbf{0}$) and thus (2.30) holds trivially.² Note also that (2.30) evidently implies (2.29), however with different choices for C in both inequalities.

Following [15] we define the following function on G,

$$\delta(M) = \inf_{\boldsymbol{k} \in \mathbb{Z}^d - \{\boldsymbol{0}\}} |^{\mathsf{T}} \boldsymbol{k} M|.$$
(2.32)

It is easily checked that $\delta(M)$ is invariant under left action of Γ , and may thus be viewed as a function on $\Gamma \backslash G$. In terms of the Iwasawa parametrization (2.22) and the Siegel set \mathcal{S}_d defined in (2.24) we have the following estimate.

Lemma 2.2. For $M = nak \in S_d$ as in (2.22), (2.24), there are constants $0 < C_1 \leq C_2$ such that for all $0 < a_d \leq 1$

$$C_1 a_d \le \delta(M) \le C_2 a_d. \tag{2.33}$$

Proof. Since $\|\boldsymbol{x}\| \ll |\boldsymbol{x}| \ll \|\boldsymbol{x}\|$ for all $\boldsymbol{x} \in \mathbb{R}^d$, we may prove the statement of the lemma for the function

$$\tilde{\delta}_d(M) = \inf_{\boldsymbol{k} \in \mathbb{Z}^d - \{\boldsymbol{0}\}} \|^{\mathsf{T}} \boldsymbol{k} M \|$$
(2.34)

instead.³ Due to the rotational invariance of the Euclidean distance we may assume that $k \in SO(d)$ is the identity.

Proof by induction. The statement trivially holds for d = 1. Therefore let us assume the assertion is true for dimension d - 1. The *j*th coefficient of the vector ${}^{\top}\mathbf{k}M$ is

$$({}^{\mathsf{T}}\boldsymbol{k}M)_j = \left(k_j + \sum_{i=1}^{j-1} k_i u_{ij}\right) a_j.$$
(2.35)

Since $a_1 \to \infty$ when $a_d \to 0$, this implies that when taking the infimum in (2.32) we must take $k_1 = 0$ for all sufficiently small a_d . Thus we now need to estimate

$$\inf_{\tilde{\boldsymbol{k}}\in\mathbb{Z}^{d-1}-\{\boldsymbol{0}\}}\max_{2\leq j\leq d}\left|\left(k_{j}+\sum_{i=1}^{j-1}k_{i}\,u_{ij}\right)a_{j}\right|=a_{1}^{-1/(d-1)}\inf_{\tilde{\boldsymbol{k}}\in\mathbb{Z}^{d-1}-\{\boldsymbol{0}\}}\max_{2\leq j\leq d}\left|\left(k_{j}+\sum_{i=1}^{j-1}k_{i}\,u_{ij}\right)\tilde{a}_{j}\right|$$
(2.36)

²Note that every admissible constant in (2.29) needs to satisfy C < 1/2; to see this, choose m = (1, 0, ..., 0), and $k \in \mathbb{Z}$ such that $|\alpha_1 + k| \leq 1/2$.

³In the following, $A \ll B$ means 'there is a constant C > 0 such that $A \leq CB$ '. If $A \ll B \ll A$ we will also use the notation $A \simeq B$.

where ${}^{\top}\tilde{\boldsymbol{k}} = (k_2, \ldots, k_d), \, \tilde{a}_j = a_1^{1/(d-1)} a_j$ so that $\tilde{a}_2 \cdots \tilde{a}_d = 1$. Now

$$\inf_{\tilde{\boldsymbol{k}}\in\mathbb{Z}^{d-1}-\{\boldsymbol{0}\}}\max_{2\leq j\leq d}\left|\left(k_{j}+\sum_{i=1}^{j-1}k_{i}\,u_{ij}\right)\tilde{a}_{j}\right|=:\delta_{d-1}(\tilde{M})\asymp\tilde{\delta}_{d-1}(\tilde{M})$$
(2.37)

where $\tilde{M} = \tilde{n}\tilde{a}$ with

$$\tilde{n} = \begin{pmatrix} 1 & u_{23} & \dots & u_{2d} \\ & \ddots & \ddots & \vdots \\ & & \ddots & u_{d-1,d} \\ & & & 1 \end{pmatrix}, \qquad \tilde{a} = \begin{pmatrix} \tilde{a}_2 & & & \\ & \ddots & & \\ & & \ddots & \\ & & & \tilde{a}_d \end{pmatrix}.$$
(2.38)

It is easily checked that $\tilde{M} \in \mathcal{S}_{d-1}$, so by the induction hypothesis, for suitable constants $0 < C_{1,d-1} \leq C_{2,d-1}$, we have

$$C_{1,d-1} \ \tilde{a}_d \le \tilde{\delta}_{d-1}(\tilde{M}) \le C_{2,d-1} \ \tilde{a}_d,$$
 (2.39)

provided $\tilde{a}_d = a_1^{1/(d-1)} a_d \leq 1$. So for a_d sufficiently small and $a_1^{1/(d-1)} a_d \leq 1$, we have

$$C_{1,d-1} \ a_d \le \delta_d(M) \le C_{2,d-1} \ a_d.$$
 (2.40)

In the remaining case $\tilde{a}_d > 1$, all \tilde{a}_j are bounded from above and below by positive constants, and hence $\tilde{\delta}_{d-1}(\tilde{M})$ is bounded from above and below by positive constants. Furthermore $\tilde{a}_d > 1$ implies $a_1^{-1/(d-1)} < a_d$, and, in view of our choice of the Siegel set, $a_1^{-1} = a_2 \cdots a_d \gg a_d^{d-1}$. So

$$a_d \ll a_1^{-1/(d-1)} < a_d$$
 (2.41)

and the required bound follows from (2.36) also for the case $\tilde{a}_d > 1$.

Lemma 2.3. Choose M as in (2.6) and suppose α satisfies condition (2.29). Then there exists a constant C' > 0 such that for all $t \ge 0$

$$\delta(ME^t) > C' \mathrm{e}^{-\theta t} \tag{2.42}$$

where

$$\theta = \frac{(d-1)\epsilon}{d+(d-1)\epsilon}.$$
(2.43)

Proof. Let us put ${}^{\mathsf{T}}\boldsymbol{k} = ({}^{\mathsf{T}}\boldsymbol{m}, k)$ with $\boldsymbol{m} \in \mathbb{Z}^{d-1}$ and $k \in \mathbb{Z}$. Then

$$\delta(ME^{t}) = \inf_{(\boldsymbol{m},k)\in\mathbb{Z}^{d}-\{\boldsymbol{0}\}} \left| \left({}^{\mathsf{T}}\boldsymbol{m}e^{-t}, \left({}^{\mathsf{T}}\boldsymbol{m}\boldsymbol{\alpha}+k\right)e^{(d-1)t} \right) W(t) \right| \\ \gg \inf_{(\boldsymbol{m},k)\in\mathbb{Z}^{d}-\{\boldsymbol{0}\}} \left| \left({}^{\mathsf{T}}\boldsymbol{m}e^{-t}, \left({}^{\mathsf{T}}\boldsymbol{m}\boldsymbol{\alpha}+k\right)e^{(d-1)t} \right) \right|$$
(2.44)

since W(t), as defined in (2.28), is bounded in G for all $t \ge 0$. Furthermore for t sufficiently large

$$\inf_{(\boldsymbol{m},k)\in\mathbb{Z}^{d}-\{\boldsymbol{0}\}}|({}^{\mathsf{T}}\boldsymbol{m}\mathrm{e}^{-t},({}^{\mathsf{T}}\boldsymbol{m}\boldsymbol{\alpha}+k)\mathrm{e}^{(d-1)t})|=\inf_{\boldsymbol{m}\in\mathbb{Z}^{d-1}-\{\boldsymbol{0}\},\ k\in\mathbb{Z}}|({}^{\mathsf{T}}\boldsymbol{m}\mathrm{e}^{-t},({}^{\mathsf{T}}\boldsymbol{m}\boldsymbol{\alpha}+k)\mathrm{e}^{(d-1)t})|$$
(2.45)

which, in view of the Diophantine condition (2.29), is bounded from below by

$$\geq \inf_{\boldsymbol{m} \in \mathbb{Z}^{d-1} - \{\boldsymbol{0}\}} |({}^{\top}\boldsymbol{m} e^{-t}, C |\boldsymbol{m}|^{-(d-1)(1+\epsilon)} e^{(d-1)t})| = e^{-\theta t} \inf_{\boldsymbol{m} \in \mathbb{Z}^{d-1} - \{\boldsymbol{0}\}} |({}^{\top}\boldsymbol{x}, C |\boldsymbol{x}|^{-(d-1)(1+\epsilon)})|$$

$$(2.46)$$

where $\boldsymbol{x} = e^{(\theta-1)t} \boldsymbol{m}$. We conclude the proof by noting that

$$\inf_{\boldsymbol{m}\in\mathbb{Z}^{d-1}-\{\boldsymbol{0}\}} |({}^{\top}\boldsymbol{x}, C|\boldsymbol{x}|^{-(d-1)(1+\epsilon)})| \ge \inf_{\boldsymbol{y}\in\mathbb{R}^{d-1}-\{\boldsymbol{0}\}} |({}^{\top}\boldsymbol{y}, C|\boldsymbol{y}|^{-(d-1)(1+\epsilon)})| > 0.$$
(2.47)

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The fact that $\epsilon = 0$ implies $\theta = 0$ is consistent with Dani's correspondence. On the other hand, $\theta < 1$ for any $\epsilon < \infty$.

2.7. Norm estimates. Let $\|\cdot\|$ denote the usual matrix norm

$$||M|| := \sup_{\boldsymbol{x}\neq \boldsymbol{0}} \frac{||M\boldsymbol{x}||}{||\boldsymbol{x}||}.$$
(2.48)

Proposition 2.4. Choose $M = M^{(0)}$ as in (2.6), and suppose α satisfies condition (2.29). Then there are constants $c_1, c_2, c_3, c_4, c_5, c_6 > 0$ such that for all $n \in \mathbb{N} \cup \{0\}$

$$||M^{(n)}|| \leq c_1 \exp[(d-1)\theta t_n],$$
 (2.49)

$$||M^{(n)^{-1}}|| \leq c_2 \exp(\theta t_n),$$
 (2.50)

$$||P^{(n)}|| \leq c_3 \exp[(d\theta + 1 - \theta)t_n],$$
 (2.51)

$$||P^{(n)^{-1}}|| \leq c_4 \exp[(d-1+\theta)t_n],$$
 (2.52)

$$||T^{(n)}|| \leq c_5 \exp[(1-\theta)\delta t_n + d\theta t_n],$$
 (2.53)

$$||T^{(n)^{-1}}|| \leq c_6 \exp[(d-1)(1-\theta)\delta t_n + d\theta t_n].$$
(2.54)

Proof. For any $M \in \mathcal{S}_d$ as in (2.22) we have, for all $0 < a_d \leq 1$,

$$|M|| \ll a_1 = (a_2 \cdots a_d)^{-1} \ll a_d^{-(d-1)},$$
 (2.55)

and

$$\|M^{-1}\| \ll a_d^{-1}.$$
(2.56)

Combine this with Lemmas 2.2 and 2.3 to obtain the bounds

$$\|M^{(n)^{-1}}\| \ll C_2 \,\delta(M^{(n)})^{-1} = C_2 \,\delta(M^{(0)}E^{t_n})^{-1} < C_2 C'^{-1} \exp(\theta t_n)$$
(2.57)

and

$$\|M^{(n)}\| \ll C_2^{d-1} \,\delta(M^{(n)})^{-(d-1)} = C_2^{d-1} \,\delta(M^{(0)}E^{t_n})^{-(d-1)} < C_2^{d-1}C'^{-(d-1)} \exp[(d-1)\theta t_n]. \quad (2.58)$$

The remaining estimates follow immediately from (2.49), (2.50) and the equations

$$P^{(n)} = M^{(n)} E^{-t_n} M^{(0)^{-1}}, \qquad T^{(n)} = M^{(n)} E^{-\delta t_n} M^{(n-1)^{-1}}.$$
(2.59)

Proposition 2.5. Choose $M = M^{(0)}$ as in (2.6), and suppose α satisfies condition (2.29). Then there is a constant $c_7 > 0$ such that for all $n \in \mathbb{N} \cup \{0\}$,

$$c_{7} \exp\left[-\theta\left(\frac{d^{2}}{1-\theta} - (d-1)\right)t_{n}\right] \le |\gamma^{(n)}| \le c_{1} \exp[(d-1)\theta t_{n}]$$
(2.60)

with c_1 as in (2.49).

Proof. The upper bound for $|\gamma^{(n)}|$ follows from (2.49), since $\gamma^{(n)} = (M^{(n)})_{dd}$ and hence $|\gamma^{(n)}| \leq ||M^{(n)}||$.

From (2.17) and the Diophantine condition (2.29) we have

$$|\gamma^{(n)}| = \gamma \exp[(d-1)t_n] |^{\mathsf{T}} \boldsymbol{p}_{21}^{(n)} \boldsymbol{\alpha} + p_{22}^{(n)})| > C \exp[(d-1)t_n] |\boldsymbol{p}_{21}^{(n)}|^{-(d-1)(1+\epsilon)}.$$
 (2.61)

Since

$$1 + \epsilon = \frac{d - 1 + \theta}{(d - 1)(1 - \theta)} \tag{2.62}$$

and

$$\boldsymbol{p}_{21}^{(n)} \leq \|P^{(n)}\| \tag{2.63}$$

the proposition follows from the estimate (2.51).

2.8. Hyperbolicity of the transfer matrices. Let

$$\boldsymbol{\omega}_{\perp}^{(n)} = \{ \boldsymbol{\xi} \in \mathbb{R}^d : \, \boldsymbol{\xi} \cdot \boldsymbol{\omega}^{(n)} = 0 \}$$
(2.64)

be the orthogonal complement of the vector

$$\boldsymbol{\omega}^{(n)} = \begin{pmatrix} \boldsymbol{\alpha}^{(n)} \\ 1 \end{pmatrix} \in \mathbb{R}^d.$$
(2.65)

Lemma 2.6. For all $\boldsymbol{\xi} \in \boldsymbol{\omega}_{\perp}^{(n-1)}, n \in \mathbb{N}$,

$$^{\mathsf{T}}T^{(n)^{-1}}\boldsymbol{\xi} = \exp(-\delta t_n)^{\mathsf{T}} (M^{(n-1)}M^{(n)^{-1}})\boldsymbol{\xi}$$
 (2.66)

Proof. This follows directly from the relation

$$E^{\delta t_n \, \mathsf{T}} M^{(n-1)} \boldsymbol{\xi} = \exp(-\delta t_n) \begin{pmatrix} \mathsf{T} A^{(n-1)} \boldsymbol{\xi}' \\ 0 \end{pmatrix} = \exp(-\delta t_n) \, \mathsf{T} M^{(n-1)} \boldsymbol{\xi}$$
(2.67)

where $\boldsymbol{\xi}' \in \mathbb{R}^{d-1}$ comprises the first d-1 components of $\boldsymbol{\xi}$.

Proposition 2.7. Choose $M = M^{(0)}$ as in (2.6), and suppose $\boldsymbol{\alpha}$ satisfies condition (2.29). Then there is a constant $\Lambda > 0$ such that for all $\boldsymbol{\xi} \in \boldsymbol{\omega}_{\perp}^{(n-1)}$, $n \in \mathbb{N}$,

$$\| {}^{\mathsf{T}}T^{(n)^{-1}}\boldsymbol{\xi} \| \leq \frac{1}{2}\Lambda \exp(-\varphi_n) \|\boldsymbol{\xi}\|$$
(2.68)

with

$$\varphi_n = (1 - \theta)\delta t_n - d\,\theta\,t_{n-1}.\tag{2.69}$$

Proof. From Lemma 2.6,

$$\| T T^{(n)^{-1}} \boldsymbol{\xi} \| \le \exp(-\delta t_n) \| M^{(n-1)} \| \| M^{(n)^{-1}} \| \| \boldsymbol{\xi} \|,$$
(2.70)

and the proposition follows from the bounds (2.49), (2.50).

Given any positive sequence $\varphi_0, \varphi_1, \ldots$, the values t_n that solve eq. (2.69) with $t_0 = 0$ are

$$t_n = \frac{1}{1 - \theta} \sum_{j=1}^n (1 + \beta)^{n-j} \varphi_j.$$
 (2.71)

where

$$\beta = \frac{d\theta}{1 - \theta}.\tag{2.72}$$

E.g., for constant $\varphi_n = \varphi > 0$, we have

$$t_n = \begin{cases} n\varphi & (\theta = 0)\\ \frac{\varphi}{d\theta} [(1+\beta)^n - 1] & (0 < \theta < 1). \end{cases}$$
(2.73)

2.9. The resonance cone. As we shall see, a crucial step in our renormalization scheme is to eliminate all far-from-resonance modes in the Fourier series, i.e., all modes labeled by integer vectors outside the cone

$$K^{(n)} = \{ \boldsymbol{\xi} \in \mathbb{R}^d : | \boldsymbol{\xi} \cdot \boldsymbol{\omega}^{(n)} | \le \sigma_n \| \boldsymbol{\xi} \| \}$$
(2.74)

for a given $\sigma_n > 0$.

Lemma 2.8. Choose $M = M^{(0)}$ as in (2.6), and suppose α satisfies condition (2.29). Then

$$\sup_{\boldsymbol{\xi} \in K^{(n-1)} - \{0\}} \frac{\|{}^{\top} T^{(n)^{-1}} \boldsymbol{\xi}\|}{\|\boldsymbol{\xi}\|} \le \left[\frac{\Lambda}{2} + c_6 \sigma_{n-1} \mathrm{e}^{d\,\delta t_n}\right] \exp\left[-(1-\theta)\delta t_n + d\,\theta\,t_{n-1}\right], \quad (2.75)$$

for all $n \in \mathbb{N}$.

Proof. We write $\boldsymbol{\xi} = \boldsymbol{\xi}_1 + \boldsymbol{\xi}_2$, where

$$\boldsymbol{\xi}_1 = \frac{\boldsymbol{\xi} \cdot \boldsymbol{\omega}^{(n-1)}}{\|\boldsymbol{\omega}^{(n-1)}\|^2} \, \boldsymbol{\omega}^{(n-1)}, \qquad \boldsymbol{\xi}_2 \in \boldsymbol{\omega}_{\perp}^{(n-1)}.$$
(2.76)

Firstly,

$$\| {}^{\mathsf{T}}T^{(n)^{-1}}\boldsymbol{\xi}_{1} \| \leq \| {}^{\mathsf{T}}T^{(n)^{-1}} \| \| \boldsymbol{\xi}_{1} \| = \| T^{(n)^{-1}} \| \frac{|\boldsymbol{\xi} \cdot \boldsymbol{\omega}^{(n-1)}|}{\| \boldsymbol{\omega}^{(n-1)} \|} \leq \sigma_{n-1} \| T^{(n)^{-1}} \| \| \boldsymbol{\xi} \|$$
(2.77)

since $\boldsymbol{\xi} \in K^{(n-1)}$ and $\|\boldsymbol{\omega}^{(n-1)}\| = \|({}^{\top}\boldsymbol{\alpha}^{(n-1)}, 1)\| \ge 1$. Hence in view of (2.54)

$$\| {}^{\mathsf{T}}T^{(n)^{-1}} \boldsymbol{\xi}_1 \| \le c_6 \sigma_{n-1} \exp[(d-1)(1-\theta)\delta t_n + d\,\theta\,t_n] \| \boldsymbol{\xi} \|.$$
(2.78)

Secondly, from Proposition 2.7 we infer

$$\|^{\top} T^{(n)^{-1}} \boldsymbol{\xi}_{2} \| \leq \frac{1}{2} \Lambda \exp\left[-(1-\theta)\delta t_{n} + d\,\theta\,t_{n-1}\right] \|\boldsymbol{\xi}\|.$$
(2.79)

This proves (2.75).

Remark 2.1. Note that if the t_n are chosen as in (2.73), and

$$\sigma_{n-1} \le \frac{1}{2} c_6^{-1} \Lambda \exp(-d\,\delta t_n),$$
(2.80)

then

$$\| {}^{\mathsf{T}}T^{(n)^{-1}}\boldsymbol{\xi} \| \le \Lambda \exp(-\varphi) \| \boldsymbol{\xi} \|$$
(2.81)

for all $\boldsymbol{\xi} \in K^{(n-1)}$, $n \in \mathbb{N}$ and $\varphi > 0$.

3. Renormalization of vector fields

3.1. **Definitions.** The transformation of a vector field X on a manifold M by a diffeomorphism $\psi: M \to M$ is given by the *pull-back* of X under ψ :

$$\psi^* X = (D\psi)^{-1} X \circ \psi.$$

As the tangent bundle of the *d*-torus is trivial, $T\mathbb{T}^d \simeq \mathbb{T}^d \times \mathbb{R}^d$, we identify the set of vector fields on \mathbb{T}^d with the set of functions from \mathbb{T}^d to \mathbb{R}^d , that can be regarded as maps of \mathbb{R}^d by lifting to the universal cover. We will make use of the analyticity to extend to the complex domain, so we will deal with complex analytic functions. We will also be considering an extra variable related to a parameter.

Remark 3.1. We will be using maps between Banach spaces over \mathbb{C} with a notion of analyticity stated as follows (cf. e.g. [13]): a map F defined on a domain is analytic if it is locally bounded and Gâteux differentiable. If it is analytic on a domain, it is continuous and Fréchet differentiable. Moreover, we have a convergence theorem which is going to be used later on. Let $\{F_k\}$ be a sequence of functions analytic and uniformly locally bounded on a domain D. If $\lim_{k\to+\infty} F_k = F$ on D, then F is analytic on D.

Let $\rho, a, b > 0, r = (a, b)$ and consider the domain

$$D_{\rho} \times B_r, \tag{3.1}$$

where $D_{\rho} = \{ \boldsymbol{x} \in \mathbb{C}^d : \| \operatorname{Im} \boldsymbol{x} \| < \rho/2\pi \}$ for the norm $\| \boldsymbol{u} \| = \sum_i |u_i|$ on \mathbb{C}^d , and

$$B_r = \left\{ \boldsymbol{y} = (y_1, \dots, y_d) \in \mathbb{C}^d \colon \sum_{i=1}^{d-1} |y_i| < a \text{ and } |y_d| < b \right\}.$$
 (3.2)

Take complex analytic functions $f: D_{\rho} \times B_r \to \mathbb{C}^d$ that are \mathbb{Z}^d -periodic on the first coordinate and on the form of the Fourier series

$$f(\boldsymbol{x}, \boldsymbol{y}) = \sum_{\boldsymbol{k} \in \mathbb{Z}^d} f_{\boldsymbol{k}}(\boldsymbol{y}) e^{2\pi i \boldsymbol{k} \cdot \boldsymbol{x}}.$$
(3.3)

Its coefficients are analytic functions $f_k \colon B_r \to \mathbb{C}^d$ with a continuous extension to the closure $\overline{B_r}$, endowed with the sup-norm:

$$\|f_{\boldsymbol{k}}\|_r = \sup_{\boldsymbol{y}\in B_r} \|f_{\boldsymbol{k}}(\boldsymbol{y})\|.$$

The Banach spaces $\mathcal{A}_{\rho,r}$ and $\mathcal{A}'_{\rho,r}$ are the subspaces of such functions such that the respective norms

$$\|f\|_{\rho,r} = \sum_{\boldsymbol{k}\in\mathbb{Z}^d} \|f_{\boldsymbol{k}}\|_r e^{\rho\|\boldsymbol{k}\|},$$

$$\|f\|'_{\rho,r} = \sum_{\boldsymbol{k}\in\mathbb{Z}^d} (1+2\pi\|\boldsymbol{k}\|) \|f_{\boldsymbol{k}}\|_r e^{\rho\|\boldsymbol{k}\|}$$

are finite. Also, write the constant Fourier mode of $f \in \mathcal{A}_{\rho,r}$ through the projection

$$\mathbb{E}f(\boldsymbol{y}) = \int_{\mathbb{T}^d} f(\boldsymbol{x}, \boldsymbol{y}) d\boldsymbol{x} = f_0(\boldsymbol{y})$$
(3.4)

into the projected space denoted by $\mathbb{E}A_r$. The norm of its derivative Df_0 is given by the operator norm $\|Df_0\|_r = \sup_{\|g\|_r=1} \|Df_0 g\|_r$.

Some of the properties of the above spaces are of easy verification. For instance, given any $f, g \in \mathcal{A}'_{\rho,r}$ we have:

- $||f(\boldsymbol{x}, \boldsymbol{y})|| \leq ||f||_{\rho, r} \leq ||f||'_{\rho, r}$ where $(\boldsymbol{x}, \boldsymbol{y}) \in D_{\rho} \times B_{r}$,
- $||f||_{\rho-\delta,r} \leq ||f||_{\rho,r}$ with $\delta < \rho$.

In order to setup notations write, according to section 2, $\boldsymbol{\omega}^{(0)} = \boldsymbol{\omega} \in \mathbb{R}^d - \{0\}, \lambda_0 = 1$ and, for $n \in \mathbb{N}$,

$$\boldsymbol{\omega}^{(n)} = \gamma^{(n)^{-1}} M^{(n)} \begin{pmatrix} 0\\ \vdots\\ 0\\ 1 \end{pmatrix} = \lambda_n P^{(n)} \boldsymbol{\omega} = \eta_n T^{(n)} \boldsymbol{\omega}^{(n-1)}, \qquad (3.5)$$

where

$$\lambda_n = \frac{\gamma}{\gamma^{(n)}} e^{(d-1)t_n} \quad \text{and} \quad \eta_n = \frac{\lambda_n}{\lambda_{n-1}}.$$
(3.6)

In the following, we will be interested in equilibria-free vector fields with a "twist" along the parameter direction. By rescaling this direction we will find the right parameter which guarantees the conjugacy to a linear flow. For a fixed $n \in \mathbb{N} \cup \{0\}$, we will be studying vector fields of the form

$$X(\boldsymbol{x},\boldsymbol{y}) = X_n^0(\boldsymbol{y}) + f(\boldsymbol{x},\boldsymbol{y}), \qquad (\boldsymbol{x},\boldsymbol{y}) \in D_\rho \times B_r, \qquad (3.7)$$

where $f \in \mathcal{A}_{\rho,r}$ and

$$X_n^0(\boldsymbol{y}) = \boldsymbol{\omega}^{(n)} + \gamma^{(n)^{-1}} M^{(n)} \boldsymbol{y}.$$
(3.8)

(We drop the second coordinate of the vector field because it will always be equal to zero – there is no dynamics along the parameter direction.) The linear transformation on y deforms the set B_r along the directions of the columns of $M^{(n)}$ (see (2.15)). In particular, its *d*th column corresponds to $\boldsymbol{\omega}^{(n)}$.

For the space of the above vector fields we use the same notation $\mathcal{A}_{\rho,r}$ and the same norm $\|\cdot\|_{\rho,r}$ without ambiguity.

3.2. Resonance modes. Given $\sigma_n > 0$ we define the far from resonance Fourier modes f_{k} as in (3.3) with respect to $\omega^{(n)}$ to be the ones whose indices k are in the cone

$$I_n^- = \{ \boldsymbol{k} \in \mathbb{Z}^d \colon |\boldsymbol{k} \cdot \boldsymbol{\omega}^{(n)}| > \sigma_n \|\boldsymbol{k}\| \}.$$
(3.9)

Similarly, the *resonant* modes correspond to the cone

$$I_n^+ = \mathbb{Z}^d - I_n^-. (3.10)$$

It is also useful to define the projections \mathbb{I}_n^+ and \mathbb{I}_n^- on $\mathcal{A}_{\rho,r}$ and $\mathcal{A}'_{\rho,r}$ by restricting the Fourier modes to I_n^+ and I_n^- , respectively. The identity operator is $\mathbb{I} = \mathbb{I}_n^+ + \mathbb{I}_n^-$. Moreover, take

$$A_{n} = \sup_{\boldsymbol{k} \in I_{n}^{+} - \{0\}} \frac{\|{}^{\top} T^{(n+1)^{-1}} \boldsymbol{k}\|}{\|\boldsymbol{k}\|}.$$
(3.11)

A useful property of the above cones is included in the Lemma below.

Lemma 3.2. If $\mathbf{k} \in I_n^-$ and $\mathbf{y} \in B_{r_n}$ with $r_n = (a_n, b_n)$,

$$a_n \le \sigma_n \left(\frac{1}{2} - b_n\right) |\gamma^{(n)}| \|M^{(n)}\|^{-1} \quad and \quad b_n < \frac{1}{2},$$
 (3.12)

then

$$\left| \boldsymbol{k} \cdot X_n^0(\boldsymbol{y}) \right| > \frac{\sigma_n}{2} \| \boldsymbol{k} \|.$$
 (3.13)

Proof. For every $\boldsymbol{y} \in B_{r_n}$ and $\boldsymbol{k} \in I_n^-$,

$$|\boldsymbol{k} \cdot (\omega^{(n)} + \gamma^{(n)^{-1}} M^{(n)} \boldsymbol{y})| = |(1 + y_d) \boldsymbol{k} \cdot \boldsymbol{\omega}^{(n)} + \gamma^{(n)^{-1}} \boldsymbol{k} \cdot M^{(n)} (y_1, \dots, y_{d-1}, 0)| > (1 - b_n) \sigma_n ||\boldsymbol{k}|| - a_n |\gamma^{(n)}|^{-1} ||M^{(n)}|| ||\boldsymbol{k}||.$$
(3.14)

Our choice of a_n yields (3.13).

3.3. Basis change, time rescaling and reparametrization. The fundamental step of the renormalization is a transformation of the domain of definition of our vector fields. This is done by a linear change of basis (coming essentially from the multidimensional continued fraction expansion of $\boldsymbol{\omega}$ – see section 2), a linear rescaling of time because the orbits take longer to cross the new torus, and a change of variables for the parameter \boldsymbol{y} in order to deal with the zero mode of the perturbation.

Let $\rho_{n-1}, a_{n-1}, b_{n-1} > 0, r_{n-1} = (a_{n-1}, b_{n-1})$ and consider a vector field

$$X(\boldsymbol{x},\boldsymbol{y}) = X_{n-1}^{0}(\boldsymbol{y}) + f(\boldsymbol{x},\boldsymbol{y}), \qquad (\boldsymbol{x},\boldsymbol{y}) \in D_{\rho_{n-1}} \times B_{r_{n-1}}, \qquad (3.15)$$

with $f \in \mathcal{A}_{\rho_{n-1},r_{n-1}}$. We are interested in the following coordinate and time linear changes:

$$\boldsymbol{x} \mapsto T^{(n)^{-1}} \boldsymbol{x}, \qquad t \mapsto \eta_n t.$$
 (3.16)

Notice that negative time rescalings are possible, meaning that we are inverting the direction of time. In addition to (3.16) we will use a transformation on \boldsymbol{y} , a map $\boldsymbol{y} \mapsto \Phi_n(X)(\boldsymbol{y})$ depending on X in a way to be defined later.

Therefore, consider the transformation

$$L_n(\boldsymbol{x}, \boldsymbol{y}) = (T^{(n)^{-1}} \boldsymbol{x}, \Phi_n(X)(\boldsymbol{y})), \qquad (\boldsymbol{x}, \boldsymbol{y}) \in \mathbb{C}^{2d},$$
(3.17)

that determines a vector field in the new coordinates as the image of the map

$$X \mapsto \mathcal{L}_n(X) = \eta_n L_n^* X.$$

That is, for $(\boldsymbol{x}, \boldsymbol{y}) \in L_n^{-1} D_{\rho_{n-1}} \times B_{r_{n-1}}$,

$$\mathcal{L}_{n}(X)(\boldsymbol{x}, \boldsymbol{y}) = \eta_{n} T^{(n)} [\boldsymbol{\omega}^{(n-1)} + \gamma^{(n-1)^{-1}} M^{(n-1)} \Phi_{n}(X)(\boldsymbol{y}) + f_{0} \circ \Phi_{n}(X)(\boldsymbol{y})] + \eta_{n} T^{(n)} (f - f_{0}) \circ L_{n}(\boldsymbol{x}, \boldsymbol{y}).$$
(3.18)

In order to eliminate the $\mathbf{k} = 0$ mode of the perturbation of X in the new coordinates and to normalise the linear term in \mathbf{y} to $\gamma^{(n)^{-1}} M^{(n)} \mathbf{y}$, using the definitions of $T^{(n)}$ and η_n we choose

$$\Phi_n(X): \boldsymbol{y} \mapsto \left(\mathrm{Id} + \gamma^{(n-1)} M^{(n-1)^{-1}} f_0 \right)^{-1} (\mathrm{e}^{-d\delta t_n} y_1, \dots, \mathrm{e}^{-d\delta t_n} y_{d-1}, y_d),$$
(3.19)

if possible. Hence,

$$\mathcal{L}_n(X)(\boldsymbol{x}, \boldsymbol{y}) = X_n^0(\boldsymbol{x}, \boldsymbol{y}) + \widetilde{\mathcal{L}}_n(f - f_0)(\boldsymbol{x}, \boldsymbol{y}), \qquad (3.20)$$

where

$$\widetilde{\mathcal{L}}_n \colon f \mapsto \eta_n T^{(n)} f \circ L_n.$$
(3.21)

Denote by Δ_{μ} the set of $X \in \mathcal{A}_{\rho_{n-1},r_{n-1}}$ such that $||f_0||_{r_{n-1}} < \mu$.

Lemma 3.3. Let $r_n = (a_n, b_n)$ and $\mu_{n-1} > 0$ such that

$$a_{n} \leq e^{d\delta t_{n}} \left[a_{n-1} - \left(1 + |\gamma^{(n-1)}| \| M^{(n-1)^{-1}} \| \right) \mu_{n-1} \right]$$

$$b_{n} \leq b_{n-1} - \left(1 + |\gamma^{(n-1)}| \| M^{(n-1)^{-1}} \| \right) \mu_{n-1}.$$
(3.22)

There exist an analytic map $\Phi_n: \Delta_{\mu_{n-1}} \to \text{Diff}(B_{r_n}, \mathbb{C}^d)$ such that, for each $X \in \Delta_{\mu_{n-1}}$, $\Phi_n(X)$ is given by (3.19) and

$$\Phi_n(X)(B_{r_n}) \subset B_{r_{n-1}}.$$
(3.23)

In case f_0 is real-analytic, $\Phi_n(X)|\mathbb{R}^d$ is also real-valued.

Proof. For $X \in \mathcal{A}_{\rho_{n-1},r_{n-1}}$ with $||f_0||_{r_{n-1}} < \mu_{n-1}$ and $\widehat{\delta}_{n-1} = (\delta_{n-1},\delta_{n-1})$ with $\delta_{n-1} = \mu_{n-1} |\gamma^{(n-1)}| \, \|M^{(n-1)^{-1}}\|,$ (3.24)

we have by the Cauchy estimate

$$\|Df_0\|_{r_{n-1}-\widehat{\delta}_{n-1}} \le \frac{\|f_0\|_{r_{n-1}}}{\delta_{n-1}} < \frac{1}{|\gamma^{(n-1)}| \|M^{(n-1)^{-1}}\|}$$

So, $F = \operatorname{Id} + \gamma^{(n-1)} M^{(n-1)^{-1}} f_0$ is a diffeomorphism on $B_{r_{n-1} - \hat{\delta}_{n-1}}$. Now, if $R_1 < a_{n-1} - \hat{\delta}_{n-1}$. $\delta_{n-1} - \mu_{n-1}, R_2 < b_{n-1} - \delta_{n-1} - \mu_{n-1} \text{ and } R = (R_1, R_2), \text{ we have } B_R \subset F(B_{r_{n-1}} - \hat{\delta}_{n-1})$ and $F^{-1}(B_R) \subset B_{r_{n-1}-\widehat{\delta}_{n-1}}$. Therefore, $\Phi_n(X)$ as given by (3.19) is a diffeomorphism on B_{r_n} by choosing $R = (e^{-d\delta t_n} a_n, b_n)$, and thus we get (3.23). In addition, $X \mapsto \Phi_n(X)$ is analytic from its dependence on f_0 . When restricted to a real domain for a real-analytic $f_0, \Phi_n(X)$ is also real-analytic.

Let the translation $R_{\boldsymbol{z}}$ on \mathbb{C}^{2d} be defined for $\boldsymbol{z} \in \mathbb{C}^d$ and given by

$$R_{\boldsymbol{z}}: (\boldsymbol{x}, \boldsymbol{y}) \mapsto (\boldsymbol{x} + \boldsymbol{z}, \boldsymbol{y}). \tag{3.25}$$

Notice that we have the following "commutative" relation:

$$L_n^* R_{\boldsymbol{z}}^* = R_{T^{(n)} \boldsymbol{z}}^* L_n^*, \qquad \boldsymbol{z} \in \mathbb{C}^d.$$
(3.26)

This also follows from the fact that Φ_n is unchanged by the introduction of the translation $R_{\boldsymbol{z}}$.

3.4. Analyticity improvement.

Lemma 3.4. If $\delta > 0$ and

$$\rho_n' \le \frac{\rho_{n-1}}{A_{n-1}} - \delta, \tag{3.27}$$

then $\widetilde{\mathcal{L}}_n$ as a map from $(\mathbb{I}_{n-1}^+ - \mathbb{E})\mathcal{A}_{\rho_{n-1},r_{n-1}} \cap \Delta_{\mu_{n-1}}$ into $(\mathbb{I} - \mathbb{E})\mathcal{A}'_{\rho'_n,r_n}$ is continuous and compact with

$$\|\widetilde{\mathcal{L}}_n\| \le |\eta_n| \|T^{(n)}\| \left(1 + \frac{2\pi}{\delta}\right).$$
(3.28)

Remark 3.5. This result means that every vector field in $\mathbb{I}_{n-1}^+ \mathcal{A}_{\rho_{n-1},r_{n-1}} \cap \Delta_{\mu_{n-1}}$, i.e. a function on $D_{\rho_{n-1}} \times B_{r_{n-1}}$ into \mathbb{C}^d , has an analytic extension to $T^{(n)^{-1}} D_{\rho'_n} \times B_{r_{n-1}}$.

Proof. Let $f \in (\mathbb{I}_{n-1}^+ - \mathbb{E})\mathcal{A}_{\rho_{n-1},r_{n-1}} \cap \Delta_{\mu_{n-1}}$. Then,

$$\|f \circ L_n\|'_{\rho'_n, r_n} \le \sum_{\mathbf{k} \in I_{n-1}^+ -\{0\}} \left(1 + 2\pi \|^{\mathsf{T}} T^{(n)^{-1}} \mathbf{k}\| \right) \|f_{\mathbf{k}} \circ M_n\|_{r_n} \mathrm{e}^{(\rho'_n - \delta + \delta)\|^{\mathsf{T}} T^{(n)^{-1}} \mathbf{k}\|}.$$
 (3.29)

By using the relation $\xi e^{-\delta \xi} \leq \delta^{-1}$ with $\xi \geq 0$, (3.11) and (3.23), we get

$$\|f \circ L_n\|'_{\rho'_n, r_n} \le (1 + 2\pi/\delta) \sum_{\substack{I_{n-1}^+ - \{0\}}} \|f_k\|_{r_{n-1}} e^{A_{n-1}(\rho'_n + \delta)} \|k\|$$

$$\le (1 + 2\pi/\delta) \|f\|_{\rho_{n-1}, r_{n-1}}.$$
(3.30)

Finally, $\|\widetilde{\mathcal{L}}_n f\|'_{\rho'_n, r_n} \leq |\eta_n| \|T^{(n)}\| \|f \circ L_n\|'_{\rho'_n, r_n}$. The above for $D_{\rho'_n} \times B_{r_n}$ is also valid for $D_{\zeta} \times B_{r_n}$, $\zeta > \rho'_n$ but satisfying a similar inequality to (3.27). Therefore, $\widetilde{\mathcal{L}}_n = \mathcal{I} \circ \mathcal{J}$, where $\mathcal{J} \colon (\mathbb{I}_{n-1}^+ - \mathbb{E})\mathcal{A}_{\rho_{n-1}, r_{n-1}} \to \mathcal{A}'_{\zeta, r_n}$ is bounded as $\widetilde{\mathcal{L}}_n$, and the inclusion map $\mathcal{I} \colon \mathcal{A}'_{\zeta,r_n} \to \mathcal{A}'_{\rho'_n,r_n}$ is compact.

For $0 < \rho''_n \le \rho'_n$, consider the inclusion

$$\mathcal{I}_n \colon \mathcal{A}'_{\rho'_n, r_n} \to \mathcal{A}'_{\rho''_n, r_n} \tag{3.31}$$

by restricting $X \in \mathcal{A}'_{\rho'_n, r_n}$ to the smaller domain $D_{\rho''_n} \times B_{r_n}$. When restricted to nonconstant modes, its norm can be estimated as follows.

Lemma 3.6. If $\phi_n \geq 1$ and

$$0 < \rho_n'' \le \rho_n' - \log(\phi_n), \tag{3.32}$$

then

$$|\mathcal{I}_n(\mathbb{I} - \mathbb{E})|| \le \phi_n^{-1}.$$
(3.33)

Proof. For $f \in (\mathbb{I} - \mathbb{E})\mathcal{A}'_{\rho'_n, r_n}$, we have

$$\|\mathcal{I}_{n}(f)\|_{\rho_{n}'',r_{n}}' \leq \sum_{\boldsymbol{k}\neq 0} (1+2\pi\|\boldsymbol{k}\|) \|f_{\boldsymbol{k}}\|_{r_{n}} \mathrm{e}^{\rho_{n}'\|\boldsymbol{k}\|} \phi_{n}^{-\|\boldsymbol{k}\|} \leq \phi_{n}^{-1} \|f\|_{\rho_{n}',r_{n}}'.$$
(3.34)

3.5. Elimination of far from resonance modes. The theorem below (to be proven in Section A.1) states the existence of a nonlinear change of coordinates U, isotopic to the identity, that cancels the I_n^- modes of any X as in (3.7) with sufficiently small f. We are eliminating only the far from resonance modes, this way avoiding the complications usually related to small divisors. We remark that the "parameter" direction \boldsymbol{y} is not affected by this change of coordinates.

For given $\rho_n, r_n, \varepsilon, \nu > 0$, denote by $\mathcal{V}_{\varepsilon}$ the open ball in $\mathcal{A}'_{\rho_n+\nu,r_n}$ centred at X^0_n with radius ε .

Theorem 3.7. Let r_n be as in (3.12), $\sigma_n < \|\boldsymbol{\omega}^{(n)}\|$ and

$$\varepsilon_n = \frac{\sigma_n}{42} \min\left\{\frac{\nu}{4\pi}, \frac{\sigma_n}{72\|\boldsymbol{\omega}^{(n)}\|}\right\}.$$
(3.35)

For all $X \in \mathcal{V}_{\varepsilon_n}$ there exists an isotopy

$$U_t: D_{\rho_n} \times B_{r_n} \to D_{\rho_n + \nu} \times B_{r_n}, (\boldsymbol{x}, \boldsymbol{y}) \mapsto (\boldsymbol{x} + u_t(\boldsymbol{x}, \boldsymbol{y}), \boldsymbol{y}),$$
(3.36)

of analytic diffeomorphisms with u_t in $\mathcal{A}'_{\rho_n,r_n}$, $t \in [0,1]$, satisfying

$$\mathbb{I}_{n}^{-}U_{t}^{*}X = (1-t)\,\mathbb{I}_{n}^{-}X, \qquad U_{0} = \mathrm{Id}\,.$$
(3.37)

This defines the maps

$$\begin{aligned} \mathfrak{U}_t \colon \mathcal{V}_{\varepsilon_n} \to \mathcal{A}'_{\rho_n, r_n} \\ X &\mapsto \mathrm{Id} + u_t \end{aligned} \tag{3.38}$$

and

$$\begin{aligned}
\mathcal{U}_t \colon \mathcal{V}_{\varepsilon_n} \to \mathbb{I}^+ \mathcal{A}_{\rho_n, r_n} \oplus (1-t) \mathbb{I}_n^- \mathcal{A}'_{\rho_n + \nu, r_n} \\
X \mapsto U_t^* X
\end{aligned} (3.39)$$

which are analytic, and satisfy the inequalities

$$\|\mathfrak{U}_{t}(X) - \operatorname{Id}\|_{\rho_{n},r_{n}}^{\prime} \leq \frac{42t}{\sigma_{n}} \|\mathbb{I}_{n}^{-}f\|_{\rho_{n},r_{n}}$$

$$\|\mathcal{U}_{t}(X) - X_{n}^{0}\|_{\rho_{n},r_{n}} \leq (3-t)\|f\|_{\rho_{n}+\nu,r_{n}}^{\prime}.$$
(3.40)

If X is real-analytic, then $\mathfrak{U}_t(X)(\mathbb{R}^{2d}) \subset \mathbb{R}^{2d}$.

Remark 3.8. Further on we will be using the above result for t = 1. So that all far from resonance modes are eliminated.

Recall the definition of the translation R_z in (3.25).

Lemma 3.9. In the conditions of Theorem 3.7, if $\boldsymbol{x} \in \mathbb{R}^d$ and $X \in \mathcal{V}_{\varepsilon_n}$, then

$$\mathfrak{U}_t(X \circ R_{\boldsymbol{x}}) = R_{\boldsymbol{x}}^{-1} \circ \mathfrak{U}_t(X) \circ R_{\boldsymbol{x}}$$
(3.41)

on \mathcal{D}_{ρ_n,r_n} .

Proof. Notice that $R_{\boldsymbol{x}}(D_{\rho_n} \times B_{r_n}) = D_{\rho_n} \times B_{r_n}$. If $U_t = \mathfrak{U}_t(X)$ is a solution of the homotopy equation (3.37) on $D_{\rho_n} \times B_{r_n}$, then $\tilde{U}_t = R_{\boldsymbol{x}}^{-1} \circ \mathfrak{U}_t(X) \circ R_{\boldsymbol{x}}$ solves the same equation for $\tilde{X} = X \circ R_{\boldsymbol{x}}$, i.e. $\mathbb{I}_n^- \tilde{X} \circ \tilde{U}_t = (1-t)\mathbb{I}_n^- \tilde{X}$, on $D_{\rho_n} \times B_{r_n}$.

3.6. Trivial limit of renormalization. Let a sequence of "widths" $0 < \sigma_n < 1$ of the resonance cones I_n^+ be given. The *n*th step renormalization operator is thus

$$\mathcal{R}_n = \mathcal{U}_n \circ \mathcal{I}_n \circ \mathcal{L}_n \circ \mathcal{R}_{n-1} \quad ext{and} \quad \mathcal{R}_0 = \mathcal{U}_0$$

where \mathcal{U}_n is the full elimination of the modes in I_n^- as in Theorem 3.7 (for t = 1). Notice that $\mathcal{R}_n(X^0 + \boldsymbol{v}) = X_n^0$, for every $\boldsymbol{v} \in \mathbb{C}^d$. From the previous sections the map \mathcal{R}_n on its domain is analytic. Also, in case a vector field X is real-analytic, the same is true for $\mathcal{R}_n(X)$.

Fix the constants ν and δ as in Theorem 3.7 and Lemma 3.4, respectively, and choose $0 < \lambda < 1$. Take

$$\Theta_{n} = \min\left\{\varepsilon_{n}, \frac{\lambda^{n}\sigma_{n}^{2}}{\prod_{i=1}^{n} \|T^{(i)}\|^{-1}\|^{2}}, \lambda^{n}\frac{\frac{\sigma_{n}|\gamma^{(n)}|}{\|M^{(n)}\|} - \frac{\sigma_{n+1}|\gamma^{(n+1)}|}{e^{d\delta t_{n+1}}\|M^{(n+1)}\|}}{1 + |\gamma^{(n)}|\|M^{(n)}\|^{-1}\|}\right\}$$
(3.42)

by assuming that the sequence of times t_n guarantees that $\Theta_n > 0$. Now, write

$$B_n = \prod_{i=0}^n A_i.$$
 (3.43)

with A_i given by (3.11). By recalling the inequalities (3.27) and (3.32) we choose, for a given $\rho_0 > 0$,

$$\rho_n = \frac{1}{B_{n-1}} \left[\rho_0 - \sum_{i=0}^{n-1} B_i \log \left(\phi_{i+1}\right) - \left(\delta + \nu\right) \sum_{i=0}^{n-1} B_i \right], \qquad (3.44)$$

where

$$\phi_n = \max\left\{2|\eta_n| \, \|T^{(n)}\| (1+2\pi\delta^{-1})\frac{\Theta_{n-1}}{\Theta_n}, 1\right\} \ge 1 \tag{3.45}$$

is to be used in Lemma 3.6.

Define the following function for every $\boldsymbol{\omega} \in \mathbb{R}^d$ associated to the choice of σ_n :

$$\mathcal{B}(\boldsymbol{\omega}) = \sum_{i=0}^{+\infty} B_i \log \left(\phi_{i+1}\right) + \left(\delta + \nu\right) \sum_{i=0}^{+\infty} B_i.$$
(3.46)

The convergence of the renormalization scheme now follows directly from our construction.

Theorem 3.10. Suppose that

$$\mathcal{B}(\boldsymbol{\omega}) < +\infty \tag{3.47}$$

and $\rho > \mathcal{B}(\boldsymbol{\omega}) + \nu$. There is K, b > 0 and $r_n = (a_n, b_n)$ with $a_n > 0$ and $b_n > b > 0$, such that if X is in a sufficiently small open ball around X^0 in \mathcal{A}_{ρ,r_0} , then (i) X is in the domain of \mathcal{R}_n and

$$\|\mathcal{R}_{n}(X) - \mathcal{R}_{n}(X^{0})\|_{\rho_{n}, r_{n}} \leq K\Theta_{n}\|X - X^{0}\|_{\rho, r_{0}}, \quad n \in \mathbb{N} \cup \{0\},$$
(3.48)

(ii) for each |s| < b there exists in $B_{r_{n-1}} \subset \mathbb{C}^d$ the limits

$$p_n^s(X) = \lim_{m \to +\infty} \Phi_n(\mathcal{R}_{n-1}(X)) \dots \Phi_m(\mathcal{R}_{m-1}(X))(0, \dots, 0, s)$$
(3.49)

and

$$\lim_{n \to +\infty} \|p_n^s(X) - (0, \dots, 0, s)\| = 0,$$
(3.50)

(iii) the map $X \mapsto p_n(X)$ is analytic and takes any real-analytic X into an analytic curve $s \mapsto p_n^s(X)$ in \mathbb{R}^d .

Proof. Let $\xi > 0$ and $\rho_0 = \rho - \nu - \xi > 0$ such that $\rho_0 > \mathcal{B}(\boldsymbol{\omega})$. Hence, by (3.44), we have R > 0 satisfying $\rho_n > RB_{n-1}^{-1}$ for all $n \in \mathbb{N}$.

Denote by c the radius of an open ball in \mathcal{A}_{ρ,r_0} centred at X^0 and containing X. If $c \leq \varepsilon_0$ we can use Theorem 3.7 to obtain $\mathcal{R}_0(X) \in \mathbb{I}_0^+ \mathcal{A}_{\rho_0,r_0}$ with $r_0 = (a_0, b_0)$ satisfying (3.12) and

$$\|\mathcal{R}_0(X) - \mathcal{R}_0(X^0)\|_{\rho_0, r_0} \le 2\|X - X^0\|_{\rho+\xi, r_0}' \le 2\xi^{-1}\|X - X^0\|_{\rho, r_0}.$$

Let $K = 2(\xi \Theta_0)^{-1}$ and assume that $c \leq K^{-1} \min\{b_0(1-\lambda), \frac{1}{2} - b_0\}$. So, (3.48) holds for n = 0.

Now, with $n \in \mathbb{N}$ we choose the following r_n :

$$a_n = \sigma_n \left(\frac{1}{2} - b_0\right) \frac{|\gamma^{(n)}|}{\|M^{(n)}\|} \quad \text{and} \quad b_n = b_0 - cK \sum_{i=0}^{n-1} \lambda^i,$$
(3.51)

so that $1/2 > b_n > b = b_0 - cK(1-\lambda)^{-1}$. The inequalities in (3.12) follow immediately. Moreover, (3.22) is also satisfied with $\mu_{n-1} = cK\Theta_{n-1}$ because

$$a_{n-1} - e^{-d\delta t_n} a_n \ge (\frac{1}{2} - b_0)(1 + |\gamma^{(n-1)}| \| M^{(n-1)^{-1}} \|) \Theta_{n-1}$$

$$\ge (1 + |\gamma^{(n-1)}| \| M^{(n-1)^{-1}} \|) c K \Theta_{n-1},$$

$$b_{n-1} - b_n = c K \lambda^{n-1}$$

$$\ge (1 + |\gamma^{(n-1)}| \| M^{(n-1)^{-1}} \|) c K \Theta_{n-1}.$$
(3.52)

Suppose that $X_{n-1} = \mathcal{R}_{n-1}(X) \in \mathbb{I}_{n-1}^+ \mathcal{A}_{\rho_{n-1}, r_{n-1}}$ and

$$||X_{n-1} - X_{n-1}^{0}||_{\rho_{n-1}, r_{n-1}} \le K\Theta_{n-1}||X - X^{0}||_{\rho, r}.$$

Since (3.22) holds, Lemmas 3.3 and 3.4 are valid and, together with (3.20) and Lemma 3.6, can be used to estimate $\mathcal{I}_n \circ \mathcal{L}_n(X_{n-1})$:

$$\begin{aligned} \|\mathcal{I}_{n} \circ \mathcal{L}_{n}(X_{n-1}) - X_{n}^{0}\|_{\rho_{n}^{\prime\prime},r_{n}}^{\prime} &\leq |\eta_{n}| \, \|T^{(n)}\|(1 + 2\pi\delta^{-1})\phi_{n}^{-1}K\Theta_{n-1}\|X - X^{0}\|_{\rho,r_{0}} \\ &= \frac{1}{2}K\Theta_{n}\|X - X^{0}\|_{\rho,r_{0}}. \end{aligned}$$
(3.53)

This vector field is inside the domain of \mathcal{U}_n as (3.12) and $\frac{1}{2}c K\Theta_n < \varepsilon_n$ are satisfied. Thus (3.48) follows from (3.40).

Denote by $f_0^{(n)}$ the constant mode of the perturbation term of X_n . By Lemma 3.3, $\Phi_n(X_{n-1}): B_{r_n} \to B_{r_{n-1}}$ is given by

$$\boldsymbol{y} \mapsto (\mathrm{Id} + g_n) \operatorname{diag}(\mathrm{e}^{-d\delta t_n}, \dots, \mathrm{e}^{-d\delta t_n}, 1) \boldsymbol{y},$$

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where

$$g_n = \left(\mathrm{Id} + \gamma^{(n-1)} M^{(n-1)^{-1}} f_0^{(n-1)} \right)^{-1} - \mathrm{Id}$$
 (3.54)

is defined on $B_{r'_n}$ with $r'_n = (e^{-d\delta t_n}a_n, b_n)$. So, for $\boldsymbol{z} \in B_{r'_n}$ there is $\boldsymbol{\xi} \in B_{r'_n}$ such that

$$g_{n}(\boldsymbol{z}) = [I + \gamma^{(n-1)} M^{(n-1)^{-1}} Df_{0}^{(n-1)}(\boldsymbol{\xi})]^{-1} [\boldsymbol{z} - \gamma^{(n-1)} M^{(n-1)^{-1}} f_{0}^{(n-1)}(0)] - \boldsymbol{z}$$

= $-[I + \gamma^{(n-1)} M^{(n-1)^{-1}} Df_{0}^{(n-1)}(\boldsymbol{\xi})]^{-1} \gamma^{(n-1)} M^{(n-1)^{-1}} [Df_{0}^{(n-1)}(\boldsymbol{\xi}) \boldsymbol{z} + f_{0}^{(n-1)}(0)]$
(3.55)

and

$$\|g_n\|_{r'_n} \le \frac{|\gamma^{(n-1)}| \|M^{(n-1)^{-1}}\|}{1 - |\gamma^{(n-1)}| \|M^{(n-1)^{-1}}\| \|Df_0^{(n-1)}\|_{r'_n}} \left(\|r'_n\| \|Df_0^{(n-1)}\|_{r'_n} + \|f_0^{(n-1)}\|_{r'_n}\right).$$
(3.56)

The choice of r_n means that

$$\min\{a_{n-1} - e^{-d\delta t_n} a_n, b_{n-1} - b_n\} \gg \min\left\{\frac{\sigma_{n-1}|\gamma^{(n-1)}|}{\|M^{(n-1)}\|}, \lambda^{n-1}\right\}.$$
(3.57)

By using (3.42) and the Cauchy estimate,

$$\|Df_0^{(n-1)}\|_{r'_n} \le \frac{\|f_0^{(n-1)}\|_{r_{n-1}}}{\min\{a_{n-1} - e^{-d\delta t_n}a_n, b_{n-1} - b_n\}} \ll \frac{\lambda^{n-1}}{|\gamma^{(n-1)}| \|M^{(n-1)^{-1}}\|}$$
(3.58)

Thus,

$$\|g_n\|_{r'_n} \ll \lambda^{n-1}.\tag{3.59}$$

Writing $\boldsymbol{y}_s = (0, \dots, 0, s)$, by induction we have

$$\Phi_n(X_{n-1})\dots\Phi_m(X_{m-1})(\boldsymbol{y}_s) = \boldsymbol{y}_s + \sum_{i=n}^m \operatorname{diag}(\mathrm{e}^{-d(t_{i-1}-t_{n-1})},\dots,\mathrm{e}^{-d(t_{i-1}-t_{n-1})},1)g_i(\boldsymbol{\xi}_i),$$
(3.60)

for some $\boldsymbol{\xi}_k \in B_{r'_k}$. Therefore, from (3.59), there exists $p_n^s(X) \in \mathbb{C}^d$ unless X is real which clearly gives $p_n^s(X) \in \mathbb{R}^d$. In addition,

$$\|p_n^s(X) - \boldsymbol{y}_s\| \le \sum_{i=n}^{+\infty} \|g_i\|_{r_i'} \ll \frac{\lambda^{n-1}}{1-\lambda}.$$
(3.61)

The maps $X \mapsto p_n^s(X)$ are analytic since the convergence is uniform. Lemma 3.3 gives us the nested sequence $\Phi_n(X_{n-1})(B_{r_n}) \subset B_{r_{n-1}}$. So, as $\boldsymbol{y}_s \in \bigcap_{i \in \mathbb{N}} B_{r_i}$, it follows that $p_n^s(X) \in B_{r_{n-1}}$.

Remark 3.11. The above can be generalised for a small analyticity radius ρ by considering a sufficiently large N and applying the above theorem to $\widetilde{X} = \mathcal{U}_N \mathcal{L}_N \dots \mathcal{U}_1 \mathcal{L}_1 \mathcal{U}_0(X)$, where X is close enough to X^0 . We recover the large strip case since ρ_N is of the order of B_{N-1}^{-1} . It remains to check that $\rho_N > \mathcal{B}(\boldsymbol{\omega}^{(N)}) + \nu$. This follows from the fact that $\mathcal{B}(\boldsymbol{\omega}^{(N)}) = B_{N-1}^{-1}[\mathcal{B}(\boldsymbol{\omega}) - \mathcal{B}_N(\boldsymbol{\omega})]$ where $\mathcal{B}_N(\boldsymbol{\omega})$ is the sum of the first N terms of $\mathcal{B}(\boldsymbol{\omega})$ so that $\mathcal{B}_N(\boldsymbol{\omega}) \to \mathcal{B}(\boldsymbol{\omega})$ as $N \to +\infty$.

Lemma 3.12. If $\boldsymbol{\omega} = \begin{pmatrix} \boldsymbol{\alpha} \\ 1 \end{pmatrix}$ in \mathbb{R}^d is diophantine, i.e. $\boldsymbol{\alpha}$ satisfies (2.29) with exponent $\boldsymbol{\epsilon}$ (related to $\boldsymbol{\theta}$ by (2.43) and to $\boldsymbol{\beta}$ by (2.72)), then (3.47) is verified.

Proof. Let us set $\delta t_n = \xi t_{n-1}$, $\sigma_n = \exp(-c\delta t_n)$, $n \ge 1$, where positive constants ξ, c will be chosen later. Obviously, $t_n = (1+\xi)t_{n-1} = [(1+\xi)/\xi]\delta t_n$ and $\delta t_n = (1+\xi)\delta t_{n-1}$. We shall assume that

$$c < d(1+\xi),$$
 (3.62)

so that $\sigma_{n-1} \exp(d\delta t_n) = \exp(-c\delta t_{n-1} + d\delta t_n) = \exp[(d - c/(1 + \xi))\delta t_n]$ is much larger than Λ given by Proposition 2.7. Hence, using (2.75) we have

$$A_{n-1} \ll \exp\left[\left(-\frac{c}{1+\xi} + d - (1-\theta) + \frac{d\theta}{\xi}\right)\delta t_n\right].$$
(3.63)

We next estimate $\|\boldsymbol{\omega}^{(n)}\|$ and ε_n . It follows from (3.5) that $\|\boldsymbol{\omega}^{(n)}\| \ll \|M^{(n)}\| |\gamma^{(n)^{-1}}|$. Thus, using (2.49), (2.60) we have

$$\|\boldsymbol{\omega}^{(n)}\| \ll \exp\left(\frac{\theta}{1-\theta}d^2t_n\right) = \exp\left(\frac{\theta}{1-\theta}d^2\frac{1+\xi}{\xi}\delta t_n\right).$$
(3.64)

Since $\|\boldsymbol{\omega}^{(n)}\| \geq 1$ one gets from (3.35) that $\varepsilon_n \sim \sigma_n^2 / \|\boldsymbol{\omega}^{(n)}\|$ which together with (3.64) implies

$$\exp\left[\left(-2c - \frac{\theta}{1-\theta}d^2\frac{1+\xi}{\xi}\right)\delta t_n\right] \ll \varepsilon_n \ll \exp(-2c\delta t_n). \tag{3.65}$$

Here $X \sim Y$ means that there exist two positive constants $C_1, C_2 > 0$ such that $C_1 Y < X < C_2 Y$. Using again (2.49), (2.60) we get

$$\frac{\sigma_n |\gamma^{(n)}|}{\|M^{(n)}\|} \gg \exp\left[\left(-c - \frac{\theta}{1-\theta} d^2 \frac{1+\xi}{\xi}\right) \delta t_n\right].$$
(3.66)

Also, since $||M^{(n+1)}|| \ge \gamma^{(n+1)}$,

$$\frac{\sigma_{n+1}|\gamma^{(n+1)}|}{\|M^{(n+1)}\|\exp(d\delta t_{n+1})} \ll \exp[-(c+d)(1+\xi)\delta t_n].$$
(3.67)

We shall assume that c and ξ are chosen in such a way that

$$-c - \frac{\theta}{1-\theta} d^2 \frac{1+\xi}{\xi} > -(c+d)(1+\xi), \qquad (3.68)$$

so that

$$\frac{\sigma_n |\gamma^{(n)}|}{\|M^{(n)}\|} - \frac{\sigma_{n+1} |\gamma^{(n+1)}|}{\|M^{(n+1)}\| \exp(dt_{n+1})} \gg \frac{\sigma_n |\gamma^{(n)}|}{\|M^{(n)}\|}.$$
(3.69)

Inequality (3.68) is equivalent to the following condition

$$c > \frac{\theta}{1-\theta} \frac{1+\xi}{\xi^2} d^2 - \frac{1+\xi}{\xi} d.$$
 (3.70)

Finally, we want A_n to be small and, hence, require the exponent in (3.63) to be negative

$$-\frac{c}{1+\xi} + d - (1-\theta) + \frac{d\theta}{\xi} < 0.$$
(3.71)

Suppose that conditions (3.62), (3.70), (3.71) are satisfied. It follows immediately from the estimates above and (3.6), (2.49), (2.50), (2.54), (2.60) that

$$|\log \Theta_n|, |\log \Theta_{n-1}|, \log ||T^{(n)}||, |\log |\eta_n|| \ll \delta t_n.$$
 (3.72)

At the same time

$$B_n = \prod_{i=0}^n A_i \ll C^n \exp(-\alpha t_{n+1}), \qquad (3.73)$$

where

$$\alpha = \frac{c}{1+\xi} - d + (1-\theta) - \frac{d\theta}{\xi} > 0.$$
(3.74)

Since B_n decays exponentially with t_n and $\log \phi_n$ grows at most linearly the series (3.47) converges. To finish the proof it is enough to show that conditions (3.62), (3.70), (3.71) can be satisfied. Indeed, since $0 < \theta < 1$ we can choose ξ so large that $1 - \theta - d\theta/\xi > 0$ and

$$\frac{\theta}{1-\theta}\frac{1+\xi}{\xi^2}d^2 - \frac{1+\xi}{\xi}d < 0.$$
(3.75)

It is easy to see that all three inequalities (3.62), (3.70), (3.71) are satisfied if $(1+\xi)(d-\beta) < c < (1+\xi)d$, where $\beta = 1 - \theta - d\theta/\xi > 0$.

3.7. Analytic conjugacy to linear flow. As a consequence of Theorem 3.10, we obtain an analytic conjugacy between a vector field and the linear flow, thus proving Theorem 1.1. In the following we always assume to be in the conditions of Section 3.6.

Let $r = r_0$ and

$$\Delta = \{ X \in \mathcal{A}_{\rho,r} \colon \| X - X^0 \|_{\rho,r} < c \}$$
(3.76)

inside the domain of \mathcal{R}_n for all $n \in \mathbb{N} \cup \{0\}$. By taking $X \in \Delta$, we denote $X_n = \mathcal{R}_n(X) \in \mathbb{I}_n^+ \mathcal{A}_{\rho_n, r_n}$ so that

$$X_n = \lambda_n \left(U_0 \circ L_1 \circ U_1 \cdots L_n \circ U_n \right)^* (X), \tag{3.77}$$

where $U_k = \mathfrak{U}_k(\mathcal{I}_k\mathcal{L}_k(X_{k-1}))$ is given by Theorem 3.7 for t = 1 at the kth step and L_k is the linear rescaling as in (3.17) for X_{k-1} .

Denote by V_n the coordinate change

$$V_n: (\boldsymbol{x}, \boldsymbol{y}) \mapsto (P^{(n)^{-1}} \boldsymbol{x}, \Phi_1(X_0) \dots \Phi_n(X_{n-1})(\boldsymbol{y}))$$
(3.78)

and set $V_0 = \text{Id.}$ Thus, $L_n = V_{n-1}^{-1} \circ V_n$ and

$$X_n = \lambda_n \left(V_n \circ U_n \right)^* \left(V_{n-1} \circ U_{n-1} \circ V_{n-1}^{-1} \right)^* \cdots \left(V_1 \circ U_1 \circ V_1^{-1} \right)^* U_0^*(X).$$
(3.79)

In particular, the y-coordinate is only transformed by the second component of V_n .

Notice that if $X_n = X_n^0$ for some $n \in \mathbb{N}$,

$$\boldsymbol{y} = \Phi_1(X_0) \dots \Phi_n(X_{n-1})(0, \dots, 0, s) \in \mathbb{C}^d,$$

with |s| < b, corresponds to the parameter for which X is conjugated to $(1 + s) \omega^{(n)}$. The parameter value for the general case $X_n - X_n^0 \to 0$ as $n \to +\infty$ is $p^s(X) = p_1^s(X)$.

Lemma 3.13. There is an open ball B about X^0 in Δ such that we can find a sequence $R_n > 0$ satisfying $R_{-1} = \rho$,

$$R_n + 2\pi 42K\Theta_n^{1/2} \|X - X^0\|_{\rho,r} \le R_{n-1} \le \frac{\rho_{n-1}}{\|P^{(n-1)}\|}, \quad X \in B,$$
(3.80)

and

$$\lim_{n \to +\infty} R_n^{-1} \Theta_n^{1/2} = 0.$$
(3.81)

Proof. Let $\rho_* = \min \rho_n$. It is enough to check that $\Theta_n^{1/2} \ll \lambda^n \rho_* \prod_{i=1}^n ||T^{(i)}||^{-1}$ with $0 < \lambda < 1$ and taking $R_n = c\lambda^{-n}\Theta_n^{1/2}$ for some positive constant c. This immediately implies (3.81) and (3.80) by considering a small enough upper bound for $||H-H^0||_{\rho,r}$. \Box

Let $\operatorname{Diff}_{per}(D_{\zeta}, \mathbb{C}^d)$, $\zeta > 0$, be the Banach space of \mathbb{Z}^d -periodic diffeomorphisms $g: D_{\zeta} \to \mathbb{C}^d$ with finite norm $\|g\|_{\zeta} = \sum_{\boldsymbol{k}} \|g_{\boldsymbol{k}}\| \mathrm{e}^{\zeta \|\boldsymbol{k}\|}$, where $g_{\boldsymbol{k}} \in \mathbb{C}^d$ are the coefficients of the Fourier representation. It is simple to check that $\|g \circ P^{(n)}\|_{R_n} \leq \|g\|_{\rho_n}$.

Denote by u_n the analytic function

$$u_n \colon \Delta \to \operatorname{Diff}_{per}(D_{\rho_n}, \mathbb{C}^d) X \mapsto \mathfrak{U}_n(\mathcal{I}_n \mathcal{L}_n(X_{n-1}))(\cdot, p_{n+1}^s(X)).$$
(3.82)

As $p_{n+1}^s(X) \in B_{r_n}$, $D_{\rho_n} \times \{ \boldsymbol{y} = p_{n+1}^s(X) \}$ is inside the domain \mathcal{D}_{ρ_n,r_n} of $\mathfrak{U}_n(\mathcal{I}_n\mathcal{L}_n(X_{n-1}))$ given in Theorem 3.7. Now, for each X, define the isotopic to the identity diffeomorphism

$$W_n(X) = P^{(n)^{-1}} \circ u_n(X) \circ P^{(n)}, \qquad (3.83)$$

on $P^{(n)^{-1}}D_{\rho_n}$. If X is real-analytic, then $W_n(X)(\mathbb{R}^d) \subset \mathbb{R}^d$, since this property holds for $u_n(X)$. We also have $W_n(X^0) = \text{Id}$.

Lemma 3.14. For all $n \in \mathbb{N} \cup \{0\}$, $W_n \colon B \to \text{Diff}_{per}(D_{R_n}, \mathbb{C}^d)$ is analytic satisfying $W_n(X) \colon D_{R_n} \to D_{R_{n-1}}$ and

$$||W_n(X) - \operatorname{Id}||_{R_n} \le 42K\Theta_n^{1/2} ||X - X^0||_{\rho,r}, \quad X \in B.$$
(3.84)

Proof. For any $X \in \Delta$, in view of (3.40) we get

$$||W_n(X) - \mathrm{Id}||_{R_n} = ||P^{(n)^{-1}} \circ [u_n(X) - \mathrm{Id}] \circ P^{(n)}||_{R_n}$$

$$\leq \frac{42}{\sigma_n} ||P^{(n)^{-1}}|| ||\mathcal{I}_n \mathcal{L}_n(X_{n-1}) - X_n^0||_{\rho_n, r_n}.$$

We can bound the above by (3.84).

Now, for $\boldsymbol{x} \in D_{R_n}$ and $X \in B \subset \Delta$,

$$\|\operatorname{Im} W_n(X)(\boldsymbol{x})\| \le \|\operatorname{Im}(W_n(X)(\boldsymbol{x}) - \boldsymbol{x})\| + \|\operatorname{Im} \boldsymbol{x}\| < \|W_n(X) - \operatorname{Id}\|_{R_n} + R_n/2\pi \le R_{n-1}/2\pi.$$

So we have $W_n(X): D_{R_n} \to D_{R_{n-1}}$ and $W_n(X) \in \text{Diff}_{per}(D_{R_n}, \mathbb{C}^d)$. From the properties of \mathfrak{U}_n, W_n is analytic as a map from B into $\text{Diff}_{per}(D_{R_n}, \mathbb{C}^d)$.

Consider the analytic map $H_n: B \to \text{Diff}_{per}(D_{R_n}, \mathbb{C}^d)$ defined by the coordinate transformation $H_n(X): D_{R_n} \to D_{\rho_0}$ as

$$H_n(X) = W_0(X) \circ \dots \circ W_n(X).$$
(3.85)

Lemma 3.15. There exists c > 0 such that for $X \in B$ and $n \in \mathbb{N}$,

$$||H_n(X) - H_{n-1}(X)||_{R_n} \le c\Theta_n^{1/2} ||X - X^0||_{\rho,r}.$$
(3.86)

Proof. For each k = 0, ..., n - 1, consider the transformations

$$G_k(z, X) = (W_k(X) - \mathrm{Id}) \circ (\mathrm{Id} + G_{k+1}(z, X)) + G_{k+1}(z, X)$$

$$G_n(z, X) = z(W_n(X) - \mathrm{Id}),$$

with $(z, X) \in \{z \in \mathbb{C} : |z| < 1 + d_n\} \times B$, where we have c' > 0 such that

$$d_n = \frac{c'}{\Theta_n^{1/2} \|X - X^0\|_{\rho,r}} - 1 > 0.$$

If the image of D_{R_n} under Id $+G_{k+1}(z, X)$ is inside the domain of $W_k(X)$, or simply

$$||G_{k+1}(z,X)||_{R_n} \le (R_k - R_n)/2\pi,$$

then G_k is well-defined as an analytic map into $\text{Diff}_{per}(D_{R_n}, \mathbb{C}^d)$, and

$$||G_k(z,X)||_{R_n} \le ||W_k(X) - \mathrm{Id}||_{R_k} + ||G_{k+1}(z,X)||_{R_n}.$$

An inductive scheme shows that

$$\begin{aligned} \|G_n(z,X)\|_{R_n} &\leq (R_{n-1} - R_n)/2\pi, \\ \|G_k(z,X)\|_{R_n} &\leq \sum_{i=k}^{n-1} \|W_i(X) - \operatorname{Id}\|_{R_i} + |z| \|W_n(X) - \operatorname{Id}\|_{R_n} \\ &\leq (R_{k-1} - R_n)/2\pi. \end{aligned}$$

By Cauchy's formula

$$||H_n(X) - H_{n-1}(X)||_{R_n} = ||G_0(1, X) - G_0(0, X)||_{R_n}$$
$$= \left\| \frac{1}{2\pi i} \oint_{|z|=1+d_n/2} \frac{G_0(z, X)}{z(z-1)} dz \right\|_{R_n},$$

and

$$\|H_n(X) - H_{n-1}(X)\|_{R_n} \le \frac{2}{d_n} \sup_{|z|=1+d_n/2} \|G_0(z,X)\|_{R_n}$$
$$\ll \Theta_n^{1/2} \|X - X^0\|_{\rho,r}.$$

Consider $C^1_{per}(\mathbb{R}^d, \mathbb{C}^d)$ to be the Banach space of the \mathbb{Z}^d -periodic C^1 functions between \mathbb{R}^d and \mathbb{C}^d with norm

$$||f||_{C^1} = \max_{k \le 1} \max_{\boldsymbol{x} \in \mathbb{R}^d} ||D^k f(\boldsymbol{x})||.$$
(3.87)

Lemma 3.16. There exists C > 0, an open ball $B' \subset B$ about X^0 and an analytic map $H: B' \to \text{Diff}_{per}(\mathbb{R}^d, \mathbb{C}^d)$ such that for $X \in B'$, $H(X) = \lim_{n \to +\infty} H_n(X)$ and

$$||H(X) - \mathrm{Id}||_{C^1} \le C ||X - X^0||_{\rho, r}.$$
(3.88)

If $X \in B'$ is real-analytic, then $H(X) \in \text{Diff}_{per}(\mathbb{R}^d, \mathbb{R}^d)$.

Proof. As the domains D_{R_n} are shrinking, we consider the restrictions of $W_n(X)$ and $H_n(X)$ to \mathbb{R}^d , and estimate their C^1 norms from the respective norms in $\text{Diff}_{per}(D_{R_n}, \mathbb{C}^d)$. More precisely, for any $X \in B$, making use of Lemma 3.15,

$$\|H_{n}(X) - H_{n-1}(X)\|_{C^{1}} \leq \max_{k \leq 1} \sup_{\boldsymbol{x} \in D_{R_{n}/2}} \|D^{k}[H_{n}(X)(\boldsymbol{x}) - H_{n-1}(X)(\boldsymbol{x})]\|$$

$$\leq \frac{4}{R_{n}} \|H_{n}(X) - H_{n-1}(X)\|_{R_{n}},$$
(3.89)

which goes to zero by (3.81). Notice that here we have used Cauchy's estimate $||D^1g||_{\zeta} \le (2\pi/e\delta)||g||_{\zeta+\delta}$ with $\zeta, \delta > 0$.

Therefore, it is shown the existence of the limit $H_n(X) \to H(X)$ as $n \to +\infty$, in the Banach space $C_{per}^1(\mathbb{R}^d, \mathbb{C}^d)$. Moreover, $||H(X) - \operatorname{Id}||_{C^1} \ll ||X - X^0||_{\rho,r}$. The convergence of H_n is uniform in B so H is analytic. As the space of close to identity diffeomorphisms is closed for the C^1 norm, H(X) is a diffeomorphism for any X sufficiently close to X^0 , i.e. $X \in B'$. The fact that, for real-analytic X, H(X) takes real values for real arguments, follows from the same property of each $W_n(X)$.

To simplify notation, write $\pi_{\boldsymbol{y}} X = X(\cdot, \boldsymbol{y})$.

Lemma 3.17. For every real-analytic $X \in B'$ and |s| < b, $[H(X)]^*(\pi_{p^s(X)}X) = (1+s) \omega$ on \mathbb{R}^d .

Proof. For each $n \in \mathbb{N}$ the definition of $H_n(X)$ and (3.79) imply that

$$H_n(X)^*(\pi_{p^s(X)}X) = \lambda_n^{-1}\pi_{p^s(X)}V_n^{-1*}(X_n).$$
(3.90)

The r.h.s. can be written as

$$\lambda_n^{-1} P^{(n)^{-1}} [\boldsymbol{\omega}^{(n)} + \gamma^{(n)^{-1}} M^{(n)} \Phi_n(X_{n-1})^{-1} \cdots \Phi_1(X_0)^{-1} p^s(X)] + \lambda_n^{-1} \pi_{p^s(X)} V_n^{-1*} (X_n - X_n^0) =$$

= $(1+s) \boldsymbol{\omega} + \lambda_n^{-1} P^{(n)^{-1}} \gamma^{(n)^{-1}} M^{(n)} p_{n+1}^s(X) - s \boldsymbol{\omega} + \lambda_n^{-1} \pi_{p^s(X)} V_n^{-1*} (X_n - X_n^0).$ (3.91)

Its terms can be estimated, for $\boldsymbol{x} \in \mathbb{R}^d$, by

$$\|\lambda_n^{-1}\pi_{p^s(X)}V_n^{-1*}(X_n - X_n^0)(\boldsymbol{x})\| \le |\lambda_n^{-1}| \|P^{(n)^{-1}}\| \|X_n - X_n^0\|_{\rho_n, r_n} \ll \Theta_n^{1/2}, \qquad (3.92)$$

and

$$\frac{1}{|\lambda_n \gamma^{(n)}|} \|P^{(n)^{-1}} M^{(n)}[p_{n+1}^s(X) - (0, \dots, 0, s)]\| = \frac{\|M^{(0)}\|}{|\gamma|e^{(d-1)t_n}} \|E^{t_n}[p_{n+1}^s(X) - (0, \dots, 0, s)]\| \\ \ll \|p_{n+1}^s(X) - (0, \dots, 0, s)\|$$
(3.93)

which is controlled by (3.50). Consequently, the limit of (3.90) as $n \to +\infty$ is $(1+s) \omega$. Using the convergence of H_n we complete the proof.

Lemma 3.18. If $X \in B'$ and $\boldsymbol{x} \in \mathbb{R}^d$, then

$$H(X \circ R_{\boldsymbol{x}}) = \widehat{R}_{\boldsymbol{x}}^{-1} \circ H(X) \circ \widehat{R}_{\boldsymbol{x}}, \qquad (3.94)$$

where $\widehat{R}_{\boldsymbol{x}} \colon \boldsymbol{z} \mapsto \boldsymbol{z} + \boldsymbol{x}$ is a translation on \mathbb{C}^d .

Proof. For each $n \in \mathbb{N}$, (3.41) and (3.26) yield that $\mathcal{U}_n(X \circ R_{\boldsymbol{x}}) = \mathcal{U}_n(X) \circ R_{\boldsymbol{x}}$ and $\mathcal{L}_n(X \circ R_{\boldsymbol{x}}) = \mathcal{L}_n(X) \circ R_{T^{(n)}\boldsymbol{x}}$. This implies immediately that

$$\mathcal{R}_n(X \circ R_{\boldsymbol{x}}) = \mathcal{R}_n(X) \circ R_{P^{(n)}\boldsymbol{x}}.$$
(3.95)

Next, from a simple adaptation of (3.41) and the formula $\widehat{R}_{P^{(n)}\boldsymbol{z}} = P^{(n)}\widehat{R}_{\boldsymbol{z}}P^{(n)^{-1}}$ for $\boldsymbol{z} \in \mathbb{C}^d$, we get

$$W_n(X \circ R_{\boldsymbol{x}}) = P^{(n)^{-1}} \circ \mathfrak{U}_{a_n}(\mathcal{L}_n \mathcal{R}_{n-1}(X \circ R_{\boldsymbol{x}})) \circ P^{(n)}$$

= $\widehat{R}_{\boldsymbol{x}}^{-1} \circ W_n(X) \circ \widehat{R}_{\boldsymbol{x}}.$ (3.96)

Thus, $H_n(X \circ R_x) = \widehat{R}_x^{-1} \circ H_n(X) \circ \widehat{R}_x$. The convergence of H_n implies (3.94).

Theorem 3.19. If $v \in \text{Vect}^{\boldsymbol{\omega}}(\mathbb{T}^d)$ is sufficiently close to $\boldsymbol{\omega}$, then there exists an analytic curve $s \mapsto \boldsymbol{p}^s \in \mathbb{R}^d$ for |s| < b, and $h \in \text{Diff}^{\boldsymbol{\omega}}(\mathbb{T}^d)$ homotopic to the identity such that

$$h^*(\boldsymbol{v} + \boldsymbol{p}^s) = (1+s)\boldsymbol{\omega}. \tag{3.97}$$

The maps $\boldsymbol{v} \mapsto \boldsymbol{p}$ and $\boldsymbol{v} \mapsto h$ are analytic.

Proof. The lift $\tilde{\boldsymbol{v}}$ to \mathbb{R}^d of \boldsymbol{v} is assumed to have an analytic extension in D_{ρ} . Consider the real-analytic vector field

$$Y(\boldsymbol{x}, \boldsymbol{y}) = \widetilde{\boldsymbol{v}}(\boldsymbol{x}) + \boldsymbol{y}$$

in $\mathcal{A}_{\rho,r}$. Suppose that \boldsymbol{v} is close enough to $\boldsymbol{\omega}$ such that $Y \in B'$ and $Y \circ R_{\boldsymbol{z}} \in B'$ for some $\eta > 0$ and $\boldsymbol{z} \in D_{\eta}$. Then, the parameter $\boldsymbol{p}^s = p^s(Y) \in \mathbb{R}^d$ and the C^1 -diffeomorphism $h = H(Y) \mod 1$ verify (3.97).

We now want to extend h analytically to a complex neighbourhood of its domain. Take $\tilde{h}(\boldsymbol{z}) = \boldsymbol{z} + H(Y \circ R_{\boldsymbol{z}})(0), \, \boldsymbol{z} \in D_{\eta}$. The maps $\boldsymbol{z} \mapsto Y \circ R_{\boldsymbol{z}}$ and $X \mapsto H(X)$ are analytic and $C_{per}^1(\mathbb{R}^d, \mathbb{C}^d) \ni g \mapsto g(0)$ is bounded. As \tilde{h} involves their composition, it is analytic on the domain D_{η} and \mathbb{Z}^d -periodic. From (3.94), for any $\boldsymbol{x} \in \mathbb{R}^d$, we have

$$\widetilde{h}(\boldsymbol{x}) \mod 1 = (\boldsymbol{x} + \widehat{R}_{\boldsymbol{x}}^{-1} \circ H(Y) \circ \widehat{R}_{\boldsymbol{x}}(0)) \mod 1$$
$$= (\boldsymbol{x} + H(Y)(\boldsymbol{x}) - \boldsymbol{x}) \mod 1$$
$$= h(\boldsymbol{x}).$$
(3.98)

The extension of h is a complex analytic diffeomorphism, thus h is a real-analytic diffeomorphism.

4. Renormalization of Hamiltonian flows

4.1. **Preliminaries.** Consider the symplectic manifold $T^*\mathbb{T}^d$ with respect to the canonical symplectic form $\sum_{i=1}^d dy_i \wedge dx_i$. As the cotangent bundle of \mathbb{T}^d is trivial, $T^*\mathbb{T}^d \simeq \mathbb{T}^d \times \mathbb{R}^d$, we identify functions on $T^*\mathbb{T}^d$ with functions on $\mathbb{T}^d \times \mathbb{R}^d$. By lifting to the universal cover, we consider functions from \mathbb{R}^{2d} into \mathbb{R} and extend them to the complex domain.

Let Ω be a neighbourhood of $\mathbb{R}^d \times \{0\}$ in \mathbb{C}^{2d} . A Hamiltonian is a complex analytic function $H: \Omega \to \mathbb{C}$, \mathbb{Z}^d -periodic on the first coordinate, written on the form of a Taylor-Fourier series

$$H(\boldsymbol{x}, \boldsymbol{y}) = \sum_{(\boldsymbol{k}, \boldsymbol{\nu}) \in I} H_{\boldsymbol{k}, \boldsymbol{\nu}} \boldsymbol{y}^{\boldsymbol{\nu}} e^{2\pi i \boldsymbol{k} \cdot \boldsymbol{x}}, \qquad (\boldsymbol{x}, \boldsymbol{y}) \in \Omega,$$
(4.1)

where $I = \mathbb{Z}^d \times (\mathbb{N} \cup \{0\})^d$, $H_{\boldsymbol{k},\boldsymbol{\nu}} \in \mathbb{C}$ and $\boldsymbol{y}^{\boldsymbol{\nu}} = y_1^{\nu_1} \dots y_d^{\nu_d}$.

Let the positive real numbers ρ and r be given in order to determine the domain

$$\mathcal{D}_{\rho,r} = D_{\rho} \times B_r,\tag{4.2}$$

where

$$D_{\rho} = \{ \boldsymbol{x} \in \mathbb{C}^{d} \colon \|\operatorname{Im} \boldsymbol{x}\| < \rho/2\pi \} \text{ and}$$

$$B_{r} = \{ \boldsymbol{y} \in \mathbb{C}^{d} \colon \|\boldsymbol{y}\| < r \},$$

$$(4.3)$$

for the norm $\|\boldsymbol{u}\| = \sum_{i=1}^{d} |u_i|$ on \mathbb{C}^d . Consider the Banach space $\mathcal{A}_{\rho,r}$ of Hamiltonians defined on $\Omega = \mathcal{D}_{\rho,r}$, which extend continuously to the boundary and with finite norm

$$||H||_{\rho,r} = \sum_{(\mathbf{k},\mathbf{\nu})\in I} |H_{\mathbf{k},\mathbf{\nu}}| r^{||\mathbf{\nu}||} e^{\rho ||\mathbf{k}||}.$$
(4.4)

Similarly, take a norm on the product space $\mathcal{A}_{\rho,r}^{2d} = \mathcal{A}_{\rho,r} \times \cdots \times \mathcal{A}_{\rho,r}$ given by $\|(H_1, \ldots, H_{2d})\|_{\rho,r} = \sum_{i=1}^{2d} \|H_i\|_{\rho,r}$. Using this we have the Banach space $\mathcal{A}'_{\rho,r}$ of the Hamiltonians $H \in \mathcal{A}_{\rho,r}$ with finite norm

$$||H||'_{\rho,r} = ||H||_{\rho,r} + ||\nabla H||_{\rho,r}$$

A useful property is the Cauchy estimate: for any $\delta > 0$ we have

$$\begin{aligned} \|\partial_i H\|_{\rho,r} &\leq \frac{2\pi}{\delta} \|H\|_{\rho+\delta,r}, \quad H \in \mathcal{A}_{\rho+\delta,r}, \quad 1 \leq i \leq d, \\ \|\partial_j H\|_{\rho,r} &\leq \frac{1}{\delta} \|H\|_{\rho,r+\delta}, \quad H \in \mathcal{A}_{\rho,r+\delta}, \quad d+1 \leq j \leq 2d, \end{aligned}$$
(4.5)

where ∂_k denotes the partial derivative with respect to the kth argument. In particular

$$\|H\|'_{\rho,r} \le \left(1 + \frac{2\pi + 1}{\delta}\right) \|H\|_{\rho+\delta,r+\delta}.$$
(4.6)

The constant Fourier modes are written by the projections

$$\mathbb{E}_{0}F = \int_{\mathbb{T}^{d}} F(\boldsymbol{x}, 0) d\boldsymbol{x} = F_{0,0} \quad \text{and} \quad \mathbb{E}F(\boldsymbol{y}) = \int_{\mathbb{T}^{d}} F(\boldsymbol{x}, \boldsymbol{y}) d\boldsymbol{x} = \sum_{\nu} F_{0,\nu} \boldsymbol{y}^{\nu}.$$
(4.7)

The space where $\mathbb{E}F$ lies is denoted by $\mathbb{E}A_r$ and the natural induced norm is $\|\cdot\|_r$. Similarly, we define $\mathbb{E}A'_r$ with norm $\|\cdot\|'_r$. Recall Remark 3.1 about the notion of analyticity of maps over the above spaces.

In the next sections we will be studying Hamiltonians close enough to H^0 as given by (1.7) for a given choice of $\boldsymbol{\omega} \in \mathbb{R}^d - \{0\}$ and an invertible symmetric $d \times d$ matrix Q. The norm of matrices is given by $\|Q\| = \max_{j=1...d} \sum_{i=1}^d |Q_{i,j}|$, where $Q_{i,j}$ are the entries of the $d \times d$ matrix Q.

4.2. Change of basis and rescaling. The following transformations leave invariant the dynamics of the flow generated by the Hamiltonian, producing an equivalent system. They consist of

• an affine linear symplectic transformation of the phase space,

$$L_n: (\boldsymbol{x}, \boldsymbol{y}) \mapsto (T^{(n)^{-1}} \boldsymbol{x}, {}^{\mathsf{T}} T^{(n)} \boldsymbol{y} + \boldsymbol{b}_n), \qquad (\boldsymbol{x}, \boldsymbol{y}) \in \mathbb{C}^{2d},$$
(4.8)

for some $\boldsymbol{b}_n \in \mathbb{C}^d$,

• a linear time (energy) change,

$$H \mapsto \eta_n H,\tag{4.9}$$

where η_n is as in (3.6),

• a linear action rescaling,

$$H \mapsto \frac{1}{\mu_n} H(\cdot, \mu_n \cdot) \tag{4.10}$$

with a choice of $\mu_n > 0$ to be specified later on,

• and the (trivial) elimination of the constant term

$$H \mapsto (\mathbb{I} - \mathbb{E}_0)H. \tag{4.11}$$

Recall the definition of the sequence $\boldsymbol{\omega}^{(n)}$ given by (3.5). For $n \in \mathbb{N}$ and $\rho_{n-1}, r > 0$, consider Hamiltonians of the form

$$H(\boldsymbol{x},\boldsymbol{y}) = H_{n-1}^{0}(\boldsymbol{y}) + F(\boldsymbol{x},\boldsymbol{y}), \qquad (\boldsymbol{x},\boldsymbol{y}) \in \mathcal{D}_{\rho_{n-1},r},$$
(4.12)

where

$$H_{n-1}^{0}(\boldsymbol{y}) = \boldsymbol{\omega}^{(n-1)} \cdot \boldsymbol{y} + \frac{1}{2} {}^{\mathsf{T}} \boldsymbol{y} Q_{n-1} \boldsymbol{y}$$
(4.13)

and $H_0^0 = H^0$, with Q_{n-1} being a $d \times d$ symmetric matrix. By performing the transformations above, we get a new Hamiltonian that is the image of the map

$$H \mapsto \mathcal{L}_n(H) = (\mathbb{I} - \mathbb{E}_0) \frac{\eta_n}{\mu_n} H \circ L_n(\cdot, \mu_n \cdot).$$

In order to simplify notation, we write

$$\Phi_n(\boldsymbol{y}) = \mu_n \,^{\mathsf{T}} T^{(n)} \boldsymbol{y} + \boldsymbol{b}_n.$$
(4.14)

So, for any $(\boldsymbol{x}, \boldsymbol{y}) \in L_n^{-1} \mathcal{D}_{\rho_{n-1}, r}$,

$$\mathcal{L}_{n}(H)(\boldsymbol{x},\boldsymbol{y}) = \\ = (\mathbb{I} - \mathbb{E}_{0})\frac{\eta_{n}}{\mu_{n}} \left[\boldsymbol{\omega}^{(n-1)} \cdot \Phi_{n}(\boldsymbol{y}) + \frac{1}{2} {}^{\mathsf{T}} \Phi_{n}(\boldsymbol{y}) Q^{(n-1)} \Phi_{n}(\boldsymbol{y}) + F \circ L_{n}(\boldsymbol{x},\mu_{n}\boldsymbol{y}) \right]. \quad (4.15)$$

By decomposing $F = (\mathbb{I} - \mathbb{E})F + F_0$ and using the Taylor expansion of F_0 :

$$F_0 \circ \Phi_n(\boldsymbol{y}) = F_0(\boldsymbol{b}_n) + \mu_n \,^{\mathsf{T}} \nabla F_0(\boldsymbol{b}_n) \,^{\mathsf{T}} T^{(n)} \boldsymbol{y} + \frac{\mu_n^2}{2} \,^{\mathsf{T}} \boldsymbol{y} T^{(n)} D^2 F_0(\boldsymbol{b}_n) \,^{\mathsf{T}} T^{(n)} \boldsymbol{y} + \Upsilon_n(\boldsymbol{y}), \quad (4.16)$$

with $\Upsilon_n(\boldsymbol{y}) = \mathcal{O}(\|\boldsymbol{y}\|^3)$, we get

$$\mathcal{L}_{n}(H)(\boldsymbol{x},\boldsymbol{y}) = \boldsymbol{\omega}^{(n)} \cdot \boldsymbol{y} + \eta_{n} \left[{}^{\mathsf{T}}\boldsymbol{b}_{n}Q_{n-1} + {}^{\mathsf{T}}\nabla F_{0}(\boldsymbol{b}_{n}) \right] {}^{\mathsf{T}}T^{(n)}\boldsymbol{y} + \frac{\eta_{n}\mu_{n}}{2} {}^{\mathsf{T}}\boldsymbol{y}T^{(n)} \left[Q_{n-1} + D^{2}F_{0}(\boldsymbol{b}_{n}) \right] {}^{\mathsf{T}}T^{(n)}\boldsymbol{y} + \frac{\eta_{n}}{\mu_{n}}\Upsilon_{n}(\boldsymbol{y}) + \frac{\eta_{n}}{\mu_{n}}(\mathbb{I} - \mathbb{E})F \circ L_{n}(\boldsymbol{x},\mu_{n}\boldsymbol{y}).$$

$$(4.17)$$

To "normalize" the linear terms in \boldsymbol{y} of $\mathbb{E}\mathcal{L}_n(H)$ equal to $\omega^{(n)} \cdot \boldsymbol{y}$, we choose \boldsymbol{b}_n inside the domain of ∇F_0 such that

$$Q_{n-1}\boldsymbol{b}_n + \nabla F_0(\boldsymbol{b}_n) = 0.$$
(4.18)

Hence,

$$\mathcal{L}_n(H)(\boldsymbol{x}, \boldsymbol{y}) = H_n^0(\boldsymbol{x}, \boldsymbol{y}) + \widehat{\mathcal{L}}_n(F_0)(\boldsymbol{y}) + \widetilde{\mathcal{L}}_n(F - F_0)(\boldsymbol{x}, \boldsymbol{y}).$$
(4.19)

Here we have introduced the operators

$$\widehat{\mathcal{L}}_n \colon F_0 \mapsto \frac{\eta_n}{\mu_n} \Upsilon_n \tag{4.20}$$

and

$$\widetilde{\mathcal{L}}_n \colon (\mathbb{I} - \mathbb{E})F \mapsto \frac{\eta_n}{\mu_n} (\mathbb{I} - \mathbb{E})F \circ L_n(\cdot, \mu_n \cdot)$$
(4.21)

defined in $\mathbb{E}\mathcal{A}_r$ and $(\mathbb{I} - \mathbb{E})\mathcal{A}_{\rho_{n-1},r}$, respectively. Moreover, we have used the symmetric $d \times d$ matrix

$$Q_n = \eta_n \mu_n T^{(n)} \left[Q_{n-1} + D^2 F_0(\boldsymbol{b}_n) \right]^{\top} T^{(n)}.$$
(4.22)

Denote by Δ_{γ} the set of $H \in \mathcal{A}_{\rho_{n-1},r}$ such that $||F_0||_{\rho_{n-1},r} < \gamma$.

Lemma 4.1. If $\det(Q_{n-1}) \neq 0$ and

$$\gamma_n = \frac{r^2}{16\|Q_{n-1}^{-1}\|},\tag{4.23}$$

there is $b_n \in C^1(\Delta_{\gamma_n}, \mathbb{C}^d)$ such that, for all $H \in \Delta_{\gamma_n}$, $\boldsymbol{b}_n = b_n(H)$ satisfies (4.18) and

$$\|b_n(H)\| < (2/r) \|Q_{n-1}^{-1}\| \|F_0\|_r < \frac{r}{8}.$$
(4.24)

Moreover, $det(Q_n) \neq 0$ and

$$\|Q_n^{-1}\| \le \frac{\|T^{(n)^{-1}}\| \|^{\mathsf{T}} T^{(n)^{-1}}\|}{\mu_n |\eta_n| (\|Q_{n-1}^{-1}\|^{-1} - \frac{16}{r^2} \|F_0\|_r)}.$$
(4.25)

In the case F_0 is real-analytic and Q_{n-1} is real, then $b_n(H) \in \mathbb{R}^d$ and Q_n is also real.

Proof. Consider the differentiable function $\mathcal{F}(H, \mathbf{b}) = \mathbf{b} + Q_{n-1}^{-1} \nabla F_0(\mathbf{b})$ defined on $\Delta_{\gamma_n} \times B_{r/2}$. Notice that $\mathcal{F}(H_{n-1}^0, 0) = 0$ and the derivative with respect to the second argument,

$$D_2 \mathcal{F}(H, \boldsymbol{b}) = I + Q_{n-1}^{-1} D^2 F_0(\boldsymbol{b}), \qquad (H, \boldsymbol{b}) \in \Delta_{\gamma_n} \times B_{r/2},$$

admits a bounded inverse because

$$|D^{2}F_{0}||_{r/2} = \max_{d+1 \le j \le 2d} ||\partial_{j}\nabla F_{0}||_{r/2}$$

$$\le (4/r) ||\nabla F_{0}||_{3r/4}$$

$$\le (16/r^{2}) ||F_{0}||_{r}$$

$$< ||Q_{n-1}^{-1}||^{-1}$$
(4.26)

by the Cauchy estimate. The implicit function theorem implies the existence of a C^1 function $b: H \mapsto b(H)$ in a neighbourhood of H^0_{n-1} such that

$$\mathcal{F}(H, b(H)) = b(H) + Q_{n-1}^{-1} \nabla F_0(b(H)) = 0$$

i.e. a solution of (4.18). Notice that for any $H \in \Delta_{\gamma_n}$ the operator $\mathrm{Id} - \mathcal{F}(H, \cdot)$ is a contraction with a unique fixed point **b**. Hence the domain of the C^1 function $H \mapsto b(H) = \mathbf{b}$ is extendable to Δ_{γ_n} and thus (4.24). Assuming F_0 to be real-analytic and Q_{n-1} with real entries, the same argument is still valid when considering $B_{r/2} \cap \mathbb{R}^d$. Thus, b(H) is real and Q_n is a real symmetric matrix.

From (4.26),

 $\|Q_{n-1}^{-1}D^2F_0(b_n(H))\| < 1, \quad H \in \Delta_{\gamma_n}.$ Hence, $A = Q_{n-1}[I + Q_{n-1}^{-1}D^2F_0(b_n(H))]$ is invertible. Moreover,

$$||A^{-1}|| \le 1/(||Q_{n-1}^{-1}||^{-1} - ||D^2F_0||_{r/2}).$$
(4.27)

Now,
$$Q_n^{-1} = (\eta_n \mu_n)^{-1} T^{(n)^{-1}} A^{-1} T^{(n)^{-1}}$$
, thus (4.25).

Lemma 4.2. If r < r' and

$$\mu_n < \frac{r}{4r' \| \,^{\top} T^{(n)} \|},\tag{4.28}$$

then $\widehat{\mathcal{L}}_n \colon \mathbb{E}\mathcal{A}_r \cap \Delta_{\gamma_n} \to \mathbb{E}\mathcal{A}'_{r'}$ and

$$\|\widehat{\mathcal{L}}_{n}\| \leq \mu_{n}^{2} |\eta_{n}| \left(1 + \frac{1}{2r'}\right) \frac{(4r' \|^{\top} T^{(n)}\|)^{3}}{r^{2}(r - 4r' \mu_{n} \|^{\top} T^{(n)}\|)}.$$
(4.29)

Proof. Let $H \in \Delta_{\gamma_n}$, $R = \frac{r}{4r'\mu_n \| {}^{\top}T^{(n)} \|} > 1$ and the map

$$f: \{z \in \mathbb{C} : |z| \le R\} \to \mathbb{C}^d$$

$$z \mapsto F_0(z\mu_n {}^{\mathsf{T}}T^{(n)}\boldsymbol{y} + b_n(H)).$$
(4.30)

Hence Υ_n as in (4.16) can be written as

$$f(1) - f(0) - Df(0) - \frac{1}{2}D^2f(0) = \frac{1}{2\pi i} \oint_{|z|=R} \frac{f(z)}{z^3(z-1)} dz.$$

Therefore,

$$\begin{aligned} \|\Upsilon_n\|'_{r'} &= \frac{1}{2\pi} \left\| \oint_{|z|=R} \frac{f(z)}{z^3(z-1)} dz \right\|'_{r'} \\ &\leq \frac{1}{R^2(R-1)} \sup_{|z|=R} \|F_0(z\mu_n \,^\top T^{(n)} \cdot + b_n(H))\|'_{r'}. \end{aligned}$$

For $\|y\| < r'$, in view of (4.24),

$$\sup_{|z|=R} \|z\mu_n \,^{\mathsf{T}} T^{(n)} \boldsymbol{y} + b_n(H)\| \le R\mu_n \|\,^{\mathsf{T}} T^{(n)} \|\, r' + \|b_n(H)\| < r/2,$$

and

$$\sup_{|z|=R} \|F_0(z\mu_n \,^{\top}T^{(n)} \cdot +b_n(H))\|'_{r'} \le \|F_0\|_{r/2} + R\mu_n \|\,^{\top}T^{(n)}\| \|\nabla F_0\|_{r/2} \\ \le \|F_0\|_{r/2} + \frac{1}{2r'}\|F_0\|_r \le \left(1 + \frac{1}{2r'}\right)\|F_0\|_r.$$

$$(4.31)$$

Thus, $\|\Upsilon_n\|'_{r'} \leq (1+1/2r')[R^2(R-1)]^{-1}\|F_0\|_r$ and

$$\|\widehat{\mathcal{L}}_{n}(F_{0})\|_{r'}' = \frac{|\eta_{n}|}{\mu_{n}} \|\Upsilon_{n}\|_{r'}' \leq \frac{|\eta_{n}|}{\mu_{n}} \left(1 + \frac{1}{2r'}\right) \frac{(4r'|\mu_{n}| \|^{\top}T^{(n)}\|)^{3}}{r^{2}(r - 4r'|\mu_{n}| \|^{\top}T^{(n)}\|)} \|F_{0}\|_{r}.$$

Finally, we have the relation for each $\boldsymbol{z} \in \mathbb{C}^d$:

$$R_{\boldsymbol{z}} \circ L_n = L_n \circ R_{T^{(n)}\boldsymbol{z}},\tag{4.32}$$

because $\mathbb{E}H \circ R_{\boldsymbol{z}} = \mathbb{E}H$.

4.3. Far from resonance modes. Given $\sigma_n, \tau_n > 0$, we call far from resonance modes with respect to $\boldsymbol{\omega}^{(n)}$ the Taylor-Fourier modes with indices in

$$I_n^- = \left\{ (\boldsymbol{k}, \boldsymbol{\nu}) \in I : |\boldsymbol{\omega}^{(n)} \cdot \boldsymbol{k}| > \sigma_n \|\boldsymbol{k}\|, \|\boldsymbol{\nu}\| < \tau_n \|\boldsymbol{k}\| \right\}.$$
(4.33)

The resonant modes are in $I_n^+ = I - I_n^-$. We also have the projections \mathbb{I}_n^+ and \mathbb{I}_n^- over the spaces of Hamiltonians by restricting the Taylor-Fourier modes to I_n^+ and I_n^- , respectively. The identity operator is $\mathbb{I} = \mathbb{I}_n^+ + \mathbb{I}_n^-$.

Furthermore, consider the same sequence of positive numbers A_n given in (3.11).

4.4. Analyticity improvement.

Lemma 4.3. If $\delta > 0$,

$$\rho'_{n} \leq \frac{\rho_{n-1}}{A_{n-1}} - \delta \quad and \quad \tau_{n} \geq \frac{2}{\log 2} (\rho'_{n} + \delta) \|^{\top} T^{(n)^{-1}} \|,$$
(4.34)

then $\widetilde{\mathcal{L}}_n$ as a map from $(\mathbb{I}_{n-1}^+ - \mathbb{E})\mathcal{A}_{\rho_{n-1},r} \cap \Delta_{\gamma_n}$ to $(\mathbb{I} - \mathbb{E})\mathcal{A}'_{\rho'_n,r'}$ is continuous and compact with

$$\|\widetilde{\mathcal{L}}_n\| \le \left(1 + \frac{2\pi}{\delta} + \frac{r}{2r'^2 \log 2}\right) \frac{|\eta_n|}{\mu_n}.$$
(4.35)

Proof. Let $F \in (\mathbb{I}_{n-1}^+ - \mathbb{E})\mathcal{A}_{\rho_{n-1},r} \cap \Delta_{\gamma_n}$,

$$E = \{ (0, \boldsymbol{\nu}) \colon \boldsymbol{\nu} \in (\mathbb{N} \times \{0\})^d \} \text{ and } J_n = \{ \boldsymbol{k} \in \mathbb{Z}^d \colon |\boldsymbol{k} \cdot \boldsymbol{\omega}^{(n)}| \le \sigma_n \|\boldsymbol{k}\| \}.$$
(4.36)

Using Lemma 4.1 and (4.28) we have

$$\psi_n = \mu_n \|^{\mathsf{T}} T^{(n)} \| r' + \| b_n(H) \| \le \frac{r}{4} + \frac{2}{r} \| Q_{n-1}^{-1} \| \| F_0 \|_r < \frac{r}{2}.$$
(4.37)

We want to find an upper bound on

$$\|F \circ L_{n}(\cdot,\mu_{n}\cdot)\|'_{\rho'_{n},r'} \leq \sum_{I_{n-1}^{+}-E} \left(1+2\pi \|^{\top}T^{(n)^{-1}}\boldsymbol{k}\|+\mu_{n}\|^{\top}T^{(n)}\|\|\boldsymbol{\nu}\|/r'\right)|F_{\boldsymbol{k},\boldsymbol{\nu}}|\psi_{n}^{\|\boldsymbol{\nu}\|}\mathrm{e}^{\rho'_{n}\|^{\top}T^{(n)^{-1}}\boldsymbol{k}\|} \\ \leq \sum_{I_{n-1}^{+}-E} \left(1+\frac{2\pi}{\delta}\mathrm{e}^{\delta\|^{\top}T^{(n)^{-1}}\boldsymbol{k}\|}+\frac{r}{4r'^{2}\xi_{n}}\,\mathrm{e}^{\xi_{n}\|\boldsymbol{\nu}\|}\right)|F_{\boldsymbol{k},\boldsymbol{\nu}}|\psi_{n}^{\|\boldsymbol{\nu}\|}\mathrm{e}^{\rho'_{n}\|^{\top}T^{(n)^{-1}}\boldsymbol{k}\|}, \quad (4.38)$$

where we have used the inequality $\zeta e^{-\delta\zeta} \leq \delta^{-1}$ with $\zeta \geq 0$ and again a choice of μ_n verifying (4.28). Here $\xi_n = \frac{1}{2} \log(r/\psi_n) > \frac{1}{2} \log 2$.

Consider separately the two cases corresponding to the definition of the resonance cone I_{n-1}^+ .

We deal first with the modes corresponding to $\mathbf{k} \in J_{n-1} - \{0\}$. By (3.11) and (4.34) each of these modes in (4.38) is bounded from above by

$$\left(1 + \frac{2\pi}{\delta} + \frac{r}{2r'^2 \log 2}\right) r^{\|\boldsymbol{\nu}\|} \mathrm{e}^{\rho_{n-1}\|\boldsymbol{k}\|}.$$
(4.39)

Now, consider $\|\boldsymbol{\nu}\| \geq \tau_n \|\boldsymbol{k}\|$ with $\boldsymbol{k} \neq 0$, so that

$$\| {}^{\mathsf{T}}T^{(n)^{-1}}\boldsymbol{k} \| \le \tau_n^{-1} \| {}^{\mathsf{T}}T^{(n)^{-1}} \| \| \boldsymbol{\nu} \|.$$
(4.40)

These modes in (4.38) are estimated by

$$\left(1 + \frac{2\pi}{\delta} + \frac{r}{4r'^{2}\xi_{n}} e^{\xi_{n} \|\boldsymbol{\nu}\|}\right) \left(r e^{-2\xi_{n} + (\rho_{n}' + \delta)\|^{\top} T^{(n)^{-1}} \|/\tau_{n}}\right)^{\|\boldsymbol{\nu}\|} \le \left(1 + \frac{2\pi}{\delta} + \frac{r}{2r'^{2}\log 2}\right) r^{\|\boldsymbol{\nu}\|},$$
(4.41)

where we have used (4.34).

Finally, we get

$$||F \circ L_n(\cdot, \mu_n \cdot)||'_{\rho'_n, r'} \le \left(1 + \frac{2\pi}{\delta} + \frac{r}{2r'^2 \log 2}\right) ||F||_{\rho_{n-1}, r},$$

and (4.35) follows from (4.21).

The compacticity of the operator follows in the same way as in the proof of Lemma 3.4. $\hfill \Box$

Let $0 < \rho''_n \le \rho'_n$ and the inclusion

$$\mathcal{I}_n \colon \mathcal{A}'_{\rho'_n, r'} \to \mathcal{A}'_{\rho''_n, r'}, \quad H \mapsto H | \mathcal{D}_{\rho''_n, r'}.$$
(4.42)

The norm of the $k \neq 0$ modes can be improved by the application of such inclusion. That is done in a similar way as in Lemma 3.6, therefore we omit here the proof and just repeat the result for convenience.

Lemma 4.4. If $\phi_n \geq 1$ and

$$0 < \rho_n'' \le \rho_n' - \log(\phi_n), \tag{4.43}$$

then

$$\|\mathcal{I}_n(\mathbb{I} - \mathbb{E})\| \le \phi_n^{-1}.$$
(4.44)

4.5. Elimination of far from resonance modes. The theorem below states the existence of a symplectomorphism isotopic to the identity that cancels the far from resonance modes of a Hamiltonian close to H_n^0 .

For given $\rho_n, r', \varepsilon, \nu > 0$, denote by $\mathcal{V}_{\varepsilon}$ the open ball in $\mathcal{A}'_{\rho_n+\nu,r'}$ centred at H^0_n with radius ε .

Theorem 4.5. Let 0 < r < r', $\sigma_n > 2r' ||Q_n||$ and

$$\varepsilon_n = \frac{\sigma_n^2 \min\left\{1, \frac{\nu}{2\pi}, r' - r\right\}^2}{48 \|H_n^0\|_{r'} (2\pi + 1)^2 (1 + 2\pi + \frac{\tau_n + 1}{r'})^2}.$$
(4.45)

For every $H \in \mathcal{V}_{\varepsilon_n}$ there exist an analytic symplectomorphism $g: \mathcal{D}_{\rho_n, r} \to \mathcal{D}_{\rho_n + \nu, r'}$ in $\mathcal{A}^{2d}_{\rho_n, r}$ satisfying

$$\mathbb{I}_n^- H \circ g = 0. \tag{4.46}$$

This defines the maps

$$\mathfrak{G}: \mathcal{V}_{\varepsilon_n} \to \mathcal{A}^{2d}_{\rho_n, r} \\
H \mapsto g$$
(4.47)

and

$$\begin{aligned} \mathcal{U} \colon \mathcal{V}_{\varepsilon_n} &\to \mathbb{I}_n^+ \mathcal{A}_{\rho_n, r} \\ H &\mapsto H \circ g, \end{aligned}$$

$$(4.48)$$

which are analytic, and verify

$$\|\mathfrak{G}(H) - \mathrm{Id} \|_{\rho_{n},r}^{\prime} \leq \frac{1}{\varepsilon_{n}} \|\mathbb{I}_{n}^{-}H\|_{\rho_{n},r}$$

$$\|\mathcal{U}(H) - H_{n}^{0}\|_{\rho_{n},r} \leq \left(1 + \sqrt{\frac{12\|H_{n}^{0}\|_{r'}}{\varepsilon_{n}}}\right) \|H - H_{n}^{0}\|_{\rho_{n}+\nu,r'}^{\prime}.$$

$$(4.49)$$

If H is real-analytic, then $\mathfrak{G}(H)(\mathbb{R}^{2d}) \subset \mathbb{R}^{2d}$.

We prove this theorem in Section A.2.

Recall the definition of the translation $R_{\mathbf{z}}$ on \mathbb{C}^{2d} given by (3.25).

Lemma 4.6. In the conditions of Theorem 4.5, if $\boldsymbol{x} \in \mathbb{R}^d$ and $H \in \mathcal{V}_{\epsilon_n}$, then

$$\mathfrak{G}(H \circ R_{\boldsymbol{x}}) = R_{\boldsymbol{x}}^{-1} \circ \mathfrak{G}(H) \circ R_{\boldsymbol{x}}$$
(4.50)

on $\mathcal{D}_{\rho_n,r}$.

Proof. If $g = \mathfrak{G}(H)$ is a solution of (4.46) in $\mathcal{D}_{\rho_n,r}$, then $\tilde{g} = R_{\boldsymbol{x}}^{-1} \circ \mathfrak{G}(H) \circ R_{\boldsymbol{x}}$ solves the same equation for $\tilde{H} = H \circ R_{\boldsymbol{x}}$, i.e. $\mathbb{I}^- \tilde{H} \circ \tilde{g} = 0$ in $\mathcal{D}_{\rho_n,r}$.

4.6. Convergence of renormalization. For a resonance cone I_n^+ and $\mu_n > 0$, the *n*th step renormalization operator is defined to be

$$\mathcal{R}_n = \mathcal{U}_n \circ \mathcal{I}_n \circ \mathcal{L}_n \circ \mathcal{R}_{n-1}$$
 and $\mathcal{R}_0 = \mathcal{U}_0$

where \mathcal{U}_n is as in Theorem 4.5, cancelling all the I_n^- modes. Notice that if $H^+(\boldsymbol{y}) = H^0(\boldsymbol{y}) + \boldsymbol{v} \cdot \boldsymbol{y}$ with H^0 non-degenerate, $\mathcal{R}_n(H^+) = H_n^0$ for every $\boldsymbol{v} \in \mathbb{C}^d$. This means that the renormalizations eliminate the direction corresponding to linear terms in \boldsymbol{y} . From the previous sections the map \mathcal{R}_n on its domain of validity is analytic by construction. In addition, whenever a Hamiltonian H is real-analytic, the same is true for $\mathcal{R}_n(H)$.

Let r' > r > 0, $\rho_0 > 0$ and fix a sequence σ_n such that $\sigma_0 > 2r' ||Q||$. To complete the specification of the resonant modes and of ε_n in Theorem 4.5, take

$$\tau_n = \frac{2\rho_0 \| {}^{\top} T^{(n)^{-1}} \|}{B_{n-1} \log 2}$$
(4.51)

according to Lemma 4.3. Consider also the constants ν and δ as they appear in Theorem 4.5 and Lemma 4.3, respectively. For $n \in \mathbb{N}$, we define the non-increasing sequence

$$\Theta_{n} = \min\left\{\Theta_{n-1}, \frac{\sigma_{n}^{2}}{(4r'\|Q\|)^{2}} \prod_{i=1}^{n} \frac{2^{6}\zeta_{i}^{2}}{\|T^{(i)}\|^{2}\|^{\top}T^{(i)}\|^{2}}, \frac{\varepsilon_{n}^{3}}{\prod_{i=1}^{n} \|T^{(i)^{-1}}\|^{3}}, \prod_{i=1}^{n} \frac{\min\{|\eta_{i}|^{-3}, |\eta_{i}|\}}{2^{10}\zeta_{i}^{3}\|T^{(i)^{-1}}\|^{2}\|^{\top}T^{(i)^{-1}}\|^{6}}\right\} \leq 1, \quad (4.52)$$

with $\Theta_0 = 1$. The sequence ρ_n is as defined in (3.44) with

$$\phi_{n} = \max\left\{1, 2\left(1 + \frac{2\pi}{\delta} + \frac{r}{2r'^{2}\log 2}\right)\left(1 + \sqrt{\frac{12\|H_{n}^{0}\|_{r'}}{\varepsilon_{n}}}\right)\frac{|\eta_{n}|\Theta_{n-1}}{\mu_{n}\Theta_{n}}\right\} \ge 1, \quad (4.53)$$
$$\mu_{n} = \left(\frac{\Theta_{n}}{2^{8}\zeta_{n}\max\{1, |\eta_{n}|\}\Theta_{n-1}}\right)^{1/2} \le 1.$$

Here,

$$\zeta_{i} = \left(1 + \sqrt{\frac{12\|H_{n}^{0}\|_{r'}}{\varepsilon_{n}}}\right) \left(1 + \frac{1}{2r'}\right) \left(\frac{r'}{r}\right)^{3} \|^{\top} T^{(i)}\|^{3} > 1$$

Notice that ϕ_n is our choice for Lemma 4.4.

Define the function $\boldsymbol{\omega} \mapsto \bar{\mathcal{B}}(\boldsymbol{\omega})$ as in (3.46) but using the present choice of ϕ_n .

Theorem 4.7. Suppose that $det(Q) \neq 0$,

$$\bar{\mathcal{B}}(\boldsymbol{\omega}) < +\infty, \tag{4.54}$$

and $\rho > \hat{\mathcal{B}}(\boldsymbol{\omega}) + \nu$. There exists c, K > 0 such that if $H \in \mathcal{A}_{\rho,r'}$ and $\|H - H^0\|_{\rho,r'} < c$, then H is in the domain of \mathcal{R}_n and

$$\|\mathcal{R}_{n}(H) - \mathcal{R}_{n}(H^{0})\|_{\rho_{n},r} \le K\Theta_{n}\|H - H^{0}\|_{\rho,r'}, \quad n \in \mathbb{N} \cup \{0\}.$$
(4.55)

Proof. Let $\xi > 0$ and $\rho_0 = \rho - \nu - \xi > 0$ such that $\rho_0 > \overline{\mathcal{B}}(\boldsymbol{\omega})$. Hence, by the definition (3.44) of ρ_n , there is R > 0 satisfying $\rho_n > RB_{n-1}^{-1}$ for all $n \in \mathbb{N}$.

If $c \leq \varepsilon_0$ we use Theorem 4.5 to get $\mathcal{R}_0(H) \in \mathbb{I}_0^+ \mathcal{A}_{\rho_0, r}$ with

$$\|\mathcal{R}_0(H) - \mathcal{R}_0(H^0)\|_{\rho_0, r} \le K\Theta_0 \|H - H^0\|_{\rho, r'}$$

for some K > 0.

Now, suppose that $H_{n-1} = \mathcal{R}_{n-1}(H) \in \mathbb{I}_{n-1}^+ \mathcal{A}_{\rho_{n-1},r}, n \in \mathbb{N}$, and

$$\|H_{n-1} - H_{n-1}^{0}\|_{\rho_{n-1},r} \leq K\Theta_{n-1} \|H - H^{0}\|_{\rho,r'},$$

$$\|Q_{n-1}\| \leq \|Q\| \prod_{i=1}^{n-1} \frac{3}{2} \mu_{i} |\eta_{i}| \|T^{(i)}\| \|^{\top} T^{(i)}\|,$$

$$\|Q_{n-1}^{-1}\| \leq \|Q^{-1}\| \prod_{i=1}^{n-1} 2\mu_{i}^{-1} |\eta_{i}|^{-1} \|T^{(i)-1}\| \|^{\top} T^{(i)-1}\|.$$

$$(4.56)$$

So, for c small enough, we get

$$\|Q_{n-1}^{-1}\| \ll \frac{\Theta_{n-1}^{1/2}}{\Theta_{n-1}} \prod_{i=1}^{n-1} \frac{2^5 \zeta_i^{1/2} \|T^{(i)^{-1}}\| \|^{\top} T^{(i)^{-1}}\|}{|\eta_i|^{1/2}} \le \frac{r^2}{32cK\Theta_{n-1}}.$$
 (4.57)

Thus, Lemma 4.1 is valid and as a consequence $||b_n(H_{n-1})|| < r/8$.

After performing the operators \mathcal{L}_n and \mathcal{I}_n , we want to estimate the norm of the resulting Hamiltonians. The constant and non-constant Fourier modes are dealt separately in

$$\mathcal{I}_n \mathcal{L}_n(H) = H_n^0 + \widehat{\mathcal{L}}_n(\mathbb{E}H_{n-1} - H_{n-1}^0) + \mathcal{I}_n \widetilde{\mathcal{L}}_n((\mathbb{I} - \mathbb{E})H_{n-1}).$$
(4.58)

For the former we use Lemma 4.2 and for the latter Lemmas 4.3 and 4.4. That is, the definition of μ_n implies that $\mu_n \leq \frac{r}{8r' || ||^T T^{(n)} ||}$ and

$$\begin{aligned} \|\widehat{\mathcal{L}}_{n}(\mathbb{E}H_{n-1} - H_{n-1}^{0})\|_{r'}^{\prime} &\leq 2^{7}K\left(1 + \frac{1}{2r'}\right)\left(\frac{r'}{r}\right)^{3}\mu_{n}^{2}|\eta_{n}| \|^{\top}T^{(n)}\|^{3}\Theta_{n-1}\|H - H^{0}\|_{\rho,r'} \\ &\leq \frac{K}{2(1 + \sqrt{12}\|H_{n}^{0}\|_{r'}/\varepsilon_{n})}\Theta_{n}\|H - H^{0}\|_{\rho,r'}. \end{aligned}$$

$$(4.59)$$

Furthermore, ϕ_n yields

$$\begin{aligned} \|\mathcal{I}_{n}\widetilde{\mathcal{L}}_{n}((\mathbb{I}-\mathbb{E})H_{n-1})\|_{\rho_{n}',r'}' &\leq K\left(1+\frac{2\pi}{\delta}+\frac{r}{2r'^{2}\log 2}\right)\mu_{n}^{-1}\phi_{n}^{-1}|\eta_{n}|\Theta_{n-1}\|H-H^{0}\|_{\rho,r'}\\ &\leq \frac{K}{2(1+\sqrt{12}\|H_{n}^{0}\|_{r'}/\varepsilon_{n})}\Theta_{n}\|H-H^{0}\|_{\rho,r'}. \end{aligned}$$

$$(4.60)$$

Moreover, assuming c to be small enough, we obtain from (4.22), $||Q_{n-1}||^{-1} \le ||Q_{n-1}^{-1}||$ and (4.57) that

$$\begin{aligned} \|Q_n\| &\leq \mu_n |\eta_n| \, \|T^{(n)}\| \, \|^{\mathsf{T}} T^{(n)}\| \|Q_{n-1}\| (1 + 16r^{-2}cK\Theta_{n-1}\|Q_{n-1}\|^{-1}) \\ &\leq \|Q\| \prod_{i=1}^n \frac{3}{2} \mu_i |\eta_i| \, \|T^{(i)}\| \, \|^{\mathsf{T}} T^{(i)}\| \leq \frac{\sigma_n}{4r'}. \end{aligned}$$

$$\tag{4.61}$$

By using (4.25) and again (4.57),

$$\begin{aligned} \|Q_{n}^{-1}\| &\leq \frac{\|T^{(n)^{-1}}\| \|^{\top} T^{(n)^{-1}}\| \|Q_{n-1}^{-1}\|}{\mu_{n} |\eta_{n}| (1 - 16r^{-2}cK\Theta_{n-1}\|Q_{n-1}^{-1}\|)} \\ &\leq \|Q^{-1}\| \prod_{i=1}^{n} 2\mu_{i}^{-1} |\eta_{i}|^{-1} \|T^{(i)^{-1}}\| \|^{\top} T^{(i)^{-1}}\|. \end{aligned}$$

$$(4.62)$$

The Hamiltonian $\mathcal{I}_n \mathcal{L}_n(H_{n-1})$ is inside the domain of \mathcal{U}_n since for c small enough $\frac{1}{2}c K\Theta_n < \varepsilon_n$ and $||Q_n|| < \sigma_n/(2r')$. The result follows from (4.49).

Cf. Remark 3.11 on how to generalise the above for a small analyticity radius ρ .

Lemma 4.8. If $\boldsymbol{\omega} = \begin{pmatrix} \boldsymbol{\alpha} \\ 1 \end{pmatrix} \in \mathbb{R}^d$ is diophantine, then (4.54) is verified.

Proof. The proof follows the same lines as for Lemma 3.12, using the same choices of σ_n and t_n . In fact, to show (4.54) it is only necessary to check that the series $\sum B_n |\log |\eta_{n+1}||$, $\sum B_n \log ||T^{(n+1)}||$, $\sum B_n \log ||\boldsymbol{\omega}^{(n+1)}||$ and $\sum B_n |\log \sigma_{n+1}|$ converge. This is already done in the proof of the lemma.

4.7. Construction of the invariant torus. In the following we will always assume to be in the conditions of section 4.6. We use Theorem 4.7 to determine the existence of an $\boldsymbol{\omega}$ -invariant torus for the flow of analytic Hamiltonians H close enough to H^0 (Theorem 1.2). This follows from the construction of an analytic conjugacy between the linear flow on \mathbb{T}^d of rotation vector $\boldsymbol{\omega}$ and an orbit of H. We will use the notations

$$\lambda_n = \prod_{i=1}^n \eta_i$$
 and $\chi_n = \prod_{i=1}^n \mu_i$

Let the set Δ be given by

$$\Delta = \{ H \in \mathcal{A}_{\rho, r'} \colon \| H - H^0 \|_{\rho, r'} < c \},$$
(4.63)

which is contained in the domain of \mathcal{R}_n for all $n \in \mathbb{N} \cup \{0\}$. Given $H \in \Delta$, denote by $H_n = \mathcal{R}_n(H) \in \mathbb{I}_n^+ \mathcal{A}_{\rho_n, r}$. It is simple to check that

$$H_{n} = \frac{\lambda_{n}}{\chi_{n}} [(\mathbb{I} - \mathbb{E}_{0})H \circ g_{0} \circ L_{1}^{\mu} \circ g_{1} \circ \cdots \circ L_{n}^{\mu}] \circ g_{n}$$

$$= \frac{\lambda_{n}}{\chi_{n}} ((\mathbb{I} - \mathbb{E}_{0})H \circ g_{0} \circ [\mathcal{P}_{1}(H) \circ g_{1} \circ \mathcal{P}_{1}(H)^{-1}] \circ$$

$$\cdots \circ [\mathcal{P}_{n-1}(H) \circ g_{n-1} \circ \mathcal{P}_{n-1}(H)^{-1}] \circ \mathcal{P}_{n}(H)) \circ g_{n}.$$

(4.64)

Here, $g_k = \mathfrak{G}_k(\mathcal{L}_k(H_{k-1}))$ is given by Theorem 4.5 at the kth step and

$$L_k^{\mu} \colon (\boldsymbol{x}, \boldsymbol{y}) \mapsto L_k(\boldsymbol{x}, \mu_k \boldsymbol{y}),$$
 (4.65)

with L_k as in (4.8) obtained for H_{k-1} . In addition, the affine symplectic map $\mathcal{P}_n(H) = L_1^{\mu} \cdots L_n^{\mu}$ is

$$\mathcal{P}_n(H)\colon (\boldsymbol{x}, \boldsymbol{y}) \mapsto \left(P^{(n)^{-1}}\boldsymbol{x}, \Phi_1(H) \dots \Phi_n(H_{n-1})(\boldsymbol{y})\right)$$
(4.66)

and $\mathcal{P}_0(H) = \mathrm{Id}$, where $\Phi_k(H_{k-1})(\boldsymbol{y}) = \mu_k \,^{\mathsf{T}} T^{(k)} \boldsymbol{y} + b_k(H_{k-1})$. Notice that

$$\Phi_1(H)\dots\Phi_n(H_{n-1})(\boldsymbol{y}) = \chi_n {}^{\top} P^{(n)} \boldsymbol{y} + v_n(H), \qquad (4.67)$$

with

$$v_n(H) = b_1(H) + \sum_{i=2}^n \chi_{i-1} T^{(i-1)} b_i(H_{i-1}).$$

Let

$$a_{n}(H) = \lim_{m \to +\infty} \Phi_{n}(H_{n-1}) \dots \Phi_{m}(H_{m-1})(0)$$

= $b_{n}(H_{n-1}) + \sum_{i=n+1}^{+\infty} \mu_{n} \dots \mu_{i-1} {}^{\mathsf{T}}T^{(n)} \dots {}^{\mathsf{T}}T^{(i-1)}b_{i}(H_{i-1})$ (4.68)

if the limit exists. If that is the case,

$$a(H) = a_1(H) = \lim_{n \to +\infty} v_n(H)$$
 (4.69)

and

$$a(H) - v_n(H) = \chi_n^{\top} P^{(n)} a_{n+1}(H).$$
(4.70)

Lemma 4.9. For $H \in \Delta$ there is $a_n(H) \in B_{r/2} \subset \mathbb{C}^d$ and the map $a_n \colon H \mapsto a_n(H)$ from Δ into $B_{r/2}$ is analytic taking any real-analytic H into \mathbb{R}^d . *Proof.* From Lemma 4.1 we obtain $||b_k(H_{k-1})|| \ll 1$ for any $k \in \mathbb{N}$. Thus, by the definition of μ_n , there is $0 < \lambda < 1$ such that

$$\mu_n \dots \mu_{i-1} \| {}^{\mathsf{T}} T^{(n)} \dots {}^{\mathsf{T}} T^{(i-1)} b_i(H_{i-1}) \| \ll \lambda^{i-n}, \tag{4.71}$$

where $1 \leq n \leq i-1$. Hence, (4.68) converges and each $a_n(H)$ is well-defined in \mathbb{C}^d , unless H is real which gives $a_n(H) \in \mathbb{R}^d$. The maps $H \mapsto a_n(H)$ are analytic since the convergence is uniform.

Lemma 4.10. There is an open ball B centred at H^0 in Δ such that we can find sequences $R_n, r_n > 0$ satisfying $R_{-1} = \rho, r_{-1} = r'$,

$$R_n + 2\pi K \Theta_n^{2/3} \| H - H^0 \|_{\rho, r'} \le R_{n-1} \le \frac{\rho_{n-1}}{\| P^{(n-1)} \|},$$
(4.72)

$$r_n + K\Theta_n^{2/3} \|H - H^0\|_{\rho, r'} \le r_{n-1} \le \frac{\chi_{n-1}r}{2\| {}^{\top}P^{(n-1)^{-1}}\|},$$
(4.73)

 $H \in B$, and

$$\lim_{n \to +\infty} R_n^{-1} \Theta_n^{2/3} = 0.$$
(4.74)

Proof. Let $\rho_* = \min \rho_n$. Since χ_n is decreasing, it is enough to check that

$$\Theta_n^{2/3} \ll \min\left\{\lambda^n \rho_* \prod_{i=1}^n \|T^{(i)}\|^{-1}, \chi_n \prod_{i=1}^n \|^\top T^{(i)^{-1}}\|^{-1}\right\}$$

for some $0 < \lambda < 1$ by taking $R_n = c_1 \lambda^{-n} \Theta_n^{2/3}$ and $r_n = c_2 \Theta_n^{2/3}$ with small constants $c_1, c_2 > 0$. Thus, the inequalities (4.72) and (4.73) hold whenever we take a sufficiently small bound on $||H - H^0||_{\rho,r}$. The limit (4.74) is now immediate.

For each $H \in \Delta$, consider the isotopic to the identity analytic symplectomorphism

$$W_n(H) = \mathcal{P}_n(H) \circ \mathfrak{G}_n(\mathcal{L}_n(H_{n-1})) \circ \mathcal{P}_n(H)^{-1}$$
(4.75)

on $\mathcal{P}_n(H)\mathcal{D}_{\rho_n,r}$. In particular, $W_n(H^0) = \text{Id.}$ Notice that for H real-analytic,

$$W_n(H)(\mathbb{R}^{2d}) \subset \mathbb{R}^{2d}.$$

For a given $H \in \Delta$, define the norm $||X||_n = ||X \circ V_{a(H)}||_{R_n,r_n}$, whenever $X \circ V_{a(H)} \in \mathcal{A}^{2d}_{R_n,r_n}$, where we have introduced the vertical translation

$$V_{\boldsymbol{z}}:(\boldsymbol{x},\boldsymbol{y})\mapsto(\boldsymbol{x},\boldsymbol{y}+\boldsymbol{z}),\tag{4.76}$$

for any $\boldsymbol{z} \in \mathbb{C}^d$.

Lemma 4.11. W_n is analytic on B, and satisfies

$$W_n(H): V_{a(H)}(\mathcal{D}_{\rho_n,r_n}) \to V_{a(H)}(\mathcal{D}_{\rho_{n-1},r_{n-1}})$$

and

$$|W_n(H) - \operatorname{Id}\|_n \le K' \Theta_n^{2/3} ||H - H^0||_{\rho,r}, \quad H \in B,$$
(4.77)

for some constant K' > 0.

Proof. For $H \in \Delta$ and $(\boldsymbol{x}, \boldsymbol{y}) \in \mathcal{D}_{R_n, r_n}$,

$$\|\operatorname{Im} P^{(n)}\boldsymbol{x}\| < \|P^{(n)}\|R_n/2\pi \le \rho_n/2\pi,$$

$$\|\Phi_n^{-1}(H_{n-1})\dots\Phi_1^{-1}(H)(\boldsymbol{y}+a(H))\| = \|\chi_n^{-1} {}^{\mathsf{T}}P^{(n)^{-1}}(\boldsymbol{y}+a(H)-v_n(H))\| \qquad (4.78)$$

$$\le \chi_n^{-1}\|{}^{\mathsf{T}}P^{(n)^{-1}}\|r_n+\|a_{n+1}(H)\| < r.$$

Therefore, $\mathcal{P}_n(H)^{-1} \circ V_{a(H)}(\mathcal{D}_{R_n,r_n}) \subset \mathcal{D}_{\rho_n,r}$. Moreover, using (4.49),

$$\|W_{n}(H) - \mathrm{Id}\|_{n} = \|\widehat{\mathcal{P}}_{n}(H) \circ [\mathfrak{G}_{n}(\mathcal{I}_{n}\mathcal{L}_{n}(H_{n-1})) - \mathrm{Id}] \circ \mathcal{P}_{n}(H)^{-1} \circ V_{a(H)}\|_{R_{n},r_{n}}$$

$$\leq \varepsilon_{n}^{-1} \|\widehat{\mathcal{P}}_{n}(H)\| \|\mathcal{I}_{n}\mathcal{L}_{n}(H_{n-1}) - H_{n}^{0}\|_{\rho_{n},r'} \qquad (4.79)$$

$$\ll \Theta_{n}^{2/3}, \|H - H^{0}\|_{\rho,r},$$

where $\widehat{\mathcal{P}}_n(H)$ corresponds to the linear part $(\boldsymbol{x}, \boldsymbol{y}) \mapsto (P^{(n)^{-1}}\boldsymbol{x}, \chi_n^{\top} P^{(n)}\boldsymbol{y})$ of $\mathcal{P}_n(H)$ which has norm bounded by $\|\widehat{\mathcal{P}}_n(H)\| \leq \|P^{(n)^{-1}}\| + \chi_n\|^{\top} P^{(n)}\|.$

Now, for $(\boldsymbol{x}, \boldsymbol{y}) \in \mathcal{D}_{R_n, r_n}$ and $H \in B$,

$$\begin{aligned} \|\pi_{1} \operatorname{Im} W_{n}(H) \circ V_{a(H)}(\boldsymbol{x}, \boldsymbol{y})\| &\leq \|\operatorname{Im}(\pi_{1} W_{n}(H) \circ V_{a(H)}(\boldsymbol{x}, \boldsymbol{y}) - \boldsymbol{x})\| + \|\operatorname{Im} \boldsymbol{x}\| \\ &< \|W_{n}(H) - \operatorname{Id}\|_{n} + R_{n}/2\pi < R_{n-1}/2\pi, \\ \|\pi_{2} W_{n}(H) \circ V_{a(H)}(\boldsymbol{x}, \boldsymbol{y}) - a(H)\| &\leq \|\pi_{2} W_{n}(H) \circ V_{a(H)}(\boldsymbol{x}, \boldsymbol{y}) - \boldsymbol{y} - a(H)\| + \|\boldsymbol{y}\| \end{aligned}$$

 $< ||W_n(H) - \mathrm{Id} ||_n + r_n < r_{n-1}.$ So, $W_n(H) : V_{a(H)}(\mathcal{D}_{R_n,r_n}) \to V_{a(H)}(\mathcal{D}_{R_{n-1},r_{n-1}}).$

Consider the analytic map Γ_n on B satisfying $\Gamma_n(H) \colon V_{a(H)}(\mathcal{D}_{R_n,r_n}) \to V_{a(H)}(\mathcal{D}_{\rho,r'}),$ $\Gamma_n(H) = W_0(H) \circ \cdots \circ W_n(H).$ (4.80)

We then rewrite (4.64) as

$$H \circ \Gamma_n(H) = \frac{\chi_n}{\lambda_n} H_n \circ \mathcal{P}_n(H)^{-1} + E(H), \qquad (4.81)$$

where E(H) represents a constant (irrelevant) term. Since each $W_n(H)$ is symplectic, thus $\Gamma_n(H)$ is symplectic and $H \circ \Gamma_n(H)$ is canonically equivalent to the Hamiltonian H_n . In particular, if $H_n = H_n^0$ for some n, there is an $\boldsymbol{\omega}$ -invariant torus in the phase space of H_n . We are interested in the general case, $H_n - H_n^0 \to 0$ as $n \to +\infty$.

By a simple adaptation of Lemma 3.15, we have the following result for Γ_n .

Lemma 4.12. There is c > 0 such that for $H \in B$

$$\|\Gamma_n(H) - \Gamma_{n-1}(H)\|_n \le c\Theta_n^{2/3} \|H - H^0\|_{\rho, r'}$$

Consider the Banach space $C^1_{per}(\mathbb{R}^d, \mathbb{C}^{2d})$ of C^1 functions \mathbb{Z}^d -periodic, endowed with the norm

$$\|f\|_{C^1} = \max_{k \leq 1} \max_{\boldsymbol{x} \in \mathbb{R}^d} \|D^k f(\boldsymbol{x})\|$$

Lemma 4.13. There exist C > 0, an open ball $B' \subset B$ centred at H^0 and an analytic map $\Gamma: B' \to \text{Diff}_{per}(\mathbb{R}^d, \mathbb{C}^{2d})$ such that, for every $H \in B'$, $\Gamma(H) = \lim_{n \to +\infty} \Gamma_n(H)(\cdot, a(H))$ and

$$\|\Gamma(H) - (\mathrm{Id}, a(H))\|_{C^1} \le C \|H - H^0\|_{\rho, r'}.$$
(4.82)

If $H \in B'$ is real-analytic, then $\Gamma(H) \in \text{Diff}_{per}(\mathbb{R}^d, \mathbb{R}^{2d})$.

Proof. For each
$$H \in B$$
, by the first inequality in (4.5),

$$\|[\Gamma_n(H) - \Gamma_{n-1}(H)](\cdot, a(H))\|_{C^1} \leq \max_{k \leq 1} \sup_{\boldsymbol{x} \in D_{\rho_n/2}} \|D^k[\Gamma_n(H)(\boldsymbol{x}, a(H)) - \Gamma_{n-1}(H)(\boldsymbol{x}, a(H))]\|_{C^1}$$

$$\leq \frac{4\pi}{R_n} \|\Gamma_n(H) - \Gamma_{n-1}(H)\|_n$$
(4.83)

which is estimated using (4.74). Hence, $\Gamma_n(H)(\cdot, a(H))$ converges in the Banach space $C^1_{per}(\mathbb{R}^d, \mathbb{C}^{2d})$, and (4.82) holds. The convergence of Γ_n is uniform in B, thus Γ is

analytic. If H is sufficiently close to H^0 , $\Gamma(H)$ is in fact a diffeomorphism as the space of close to identity diffeomorphisms is closed for the C^1 norm. Finally, for H realanalytic we have $\Gamma(H)(\mathbb{R}^d) \subset \mathbb{R}^{2d}$ in view of the similar property for each $W_n(H)$. \Box

The Hamiltonian vector field of H is denoted by $X_H = \mathbb{J}\nabla H$, where $\mathbb{J}: (x, y) \mapsto (y, -x)$.

Lemma 4.14. For $H \in B'$, we have on \mathbb{R}^d

$$X_H \circ \Gamma(H) = D(\Gamma(H)) \boldsymbol{\omega}. \tag{4.84}$$

Proof. Since $\Gamma_n(H)$ is a symplectomorphism, we have for $\boldsymbol{x} \in \mathbb{R}^d$,

$$Y_{n}(\boldsymbol{x}) = X_{H} \circ \Gamma_{n}(H) \circ V_{a(H)}(\boldsymbol{x}, 0) - D(\Gamma_{n}(H)) \circ V_{a(H)}(\boldsymbol{x}, 0) X_{H^{0}}(\boldsymbol{x}, 0)$$

= $[D(\Gamma_{n}(H)) \circ V_{a(H)} X_{H \circ \Gamma_{n}(H) \circ V_{a(H)} - H^{0}}](\boldsymbol{x}, 0).$ (4.85)

Hence,

$$Y_n(\boldsymbol{x}) \| \le \|D(\Gamma_n(H))(\boldsymbol{x}, a(H))\| \|\nabla[H \circ \Gamma_n(H) \circ V_{a(H)} - H^0](\boldsymbol{x}, 0)\|.$$

$$(4.86)$$

We can now estimate the above norm by recalling (4.81) in

$$\nabla [H \circ \Gamma_n(H) \circ V_{a(H)} - H^0] = \frac{\chi_n}{\lambda_n} \nabla [(H_n - H_n^0) \circ \mathcal{P}_n(H)^{-1} \circ V_{a(H)} + H_n^0 \circ \mathcal{P}_n(H)^{-1} \circ V_{a(H)} - H_n^0 \circ \mathcal{P}_n(H_0)^{-1}]$$
(4.87)

and

$$\frac{\chi_n}{\lambda_n} \nabla [H_n^0 \circ \mathcal{P}_n(H)^{-1} \circ V_{a(H)} - H_n^0 \circ \mathcal{P}_n(H_0)^{-1}] = \\ = {}^{\mathsf{T}} [a(H) - v_n(H)] \left[Q + \sum_{i=1}^{n-1} \frac{1}{\lambda_i \chi_i} P^{(i)^{-1}} D^2 F_0^{(i)}(b_{i+1}(H_i)) {}^{\mathsf{T}} P^{(i)^{-1}} \right]. \quad (4.88)$$

Notice that by induction we get

$$Q_{n} = \lambda_{n} \chi_{n} P^{(n)} Q^{\top} P^{(n)} + \sum_{i=0}^{n-1} \frac{\lambda_{n} \chi_{n}}{\lambda_{i} \chi_{i}} P^{(n)} P^{(i)^{-1}} D^{2} F_{0}^{(i)} (b_{i+1}(H_{i}))^{\top} P^{(i)^{-1} \top} P^{(n)}.$$
(4.89)

Since
$$\sum_{i=1}^{n-1} (\chi_i | \lambda_i |)^{-1} \| P^{(i)^{-1}} \| \|^{\top} P^{(i)^{-1}} \| \Theta_i \ll 1,$$

$$\frac{\chi_n}{|\lambda_n|} \| \nabla [(H_n - H_n^0) \circ \mathcal{P}_n(H)^{-1} \circ V_{a(H)}](\boldsymbol{x}, 0) \| \ll \frac{1}{|\lambda_n|} \|^{\top} P^{(n)^{-1}} \| \| H_n - H_n^0 \|_{\rho_{n,r}} \quad (4.90)$$

and

$$||a(H) - v_n(H)|| \le \chi_n ||^{\top} P^{(n)} || ||a_{n+1}(H)|| \ll \chi_n ||^{\top} P^{(n)} || \Theta_n$$
(4.91)

we have that

$$\|\nabla [H \circ \Gamma_n(H) \circ V_{a(H)} - H^0](\boldsymbol{x}, 0)\| \ll \Theta_n^{2/3}.$$
(4.92)

Finally, from the convergence of Γ_n and

$$||D\Gamma_n(H)(\boldsymbol{x}, a(H))|| \ll \frac{1}{R_n} ||\Gamma_n(H)||_n \ll \frac{1}{R_n},$$
 (4.93)

we find that $||Y_n(\boldsymbol{x})||$ converges uniformly to 0 as $n \to +\infty$. **Lemma 4.15.** If $H \in B'$ and $\boldsymbol{x} \in \mathbb{R}^d$, then

$$\Gamma(H \circ R_{\boldsymbol{x}}) = R_{\boldsymbol{x}}^{-1} \circ \Gamma(H) \circ \widehat{R}_{\boldsymbol{x}}$$
(4.94)

where $\widehat{R}_{\boldsymbol{x}} : \boldsymbol{z} \mapsto \boldsymbol{z} + \boldsymbol{x}$ is a translation on \mathbb{C}^d .

Proof. For each $n \in \mathbb{N}$, (4.32) implies that $\mathcal{P}_n(H \circ R_z) = \mathcal{P}_n(H)$ and we know that $\mathcal{P}_n(H) \circ R_{P^{(n)}\boldsymbol{z}}^{-1} = R_{\boldsymbol{z}}^{-1} \circ \mathcal{P}_n(H), \ \boldsymbol{z} \in \mathbb{C}^d.$ So, from Lemma 4.6,

$$W_n(H \circ R_{\boldsymbol{x}}) = \mathcal{P}_n(H) \circ \mathfrak{G}_n(\mathcal{L}_n \mathcal{R}_{n-1}(H \circ R_{\boldsymbol{x}})) \circ \mathcal{P}_n(H)^{-1}$$

= $R_{\boldsymbol{x}}^{-1} \circ W_n(H) \circ R_{\boldsymbol{x}}.$ (4.95)

Thus, $\Gamma_n(H \circ R_x) = R_x^{-1} \circ \Gamma_n(H) \circ R_x$ and (4.94) follows from the convergence of Γ_n . \Box

The flow generated by X_H is denoted by ϕ_H^t taken at time $t \ge 0$. Hence,

$$\phi_{H^0}^t|_{\mathbb{T}^d \times \{0\}} = \widehat{R}_{\omega t}.$$

We prove below the existence of an invariant torus \mathcal{T} for H, i.e. an analytic conjugacy between $\phi_H^t|_{\mathcal{T}}$ and $R_{\omega t}$.

Theorem 4.16. Let $D \subset \mathbb{R}^d$ be an open ball about the origin. If $H \in C^{\boldsymbol{\omega}}(\mathbb{T}^d \times D)$ is sufficiently close to H^0 , then there exist $\mathbf{a} \in \mathbb{R}^d$ and a C^{ω} -diffeomorphism $\gamma \colon \mathbb{T}^d \to \mathbb{T}^d$ $\mathbb{T}^d \times D$ such that

$$\phi_H^t \circ \gamma = \gamma \circ \widehat{R}_{\omega t}, \quad t \ge 0, \tag{4.96}$$

and $\mathcal{T} = \gamma(\mathbb{T}^d) \simeq \mathbb{T}^d$ is a submanifold homotopic to $\{y = a\}$. Furthermore, the maps $H \mapsto \boldsymbol{a}$ and $H \mapsto \gamma$ are analytic.

Proof. The lift \widetilde{H} to $\mathbb{R}^d \times D$ of H is assumed to have a unique analytic extension to $\mathcal{D}_{\rho,r'}$. Consider the real-analytic Hamiltonian $G = \widetilde{H} \in \mathcal{A}_{\rho,r'}$. Suppose that G is close enough to H^0 such that $G \in B'$ and $G \circ R_z \in B'$ for $\eta > 0$ and $z \in D_\eta$. Then, a = a(G) and $\gamma = \Gamma(G)|_{[0,1]^d}$, C^1 and homotopic to (Id, \boldsymbol{a}), verifies (4.96). This follows from (4.84) and the equivalent equation

$$\frac{d}{dt}\Big|_{t=0}(\phi_H^t \circ \gamma) = \frac{d}{dt}\Big|_{t=0}(\gamma \circ \widehat{R}_{\omega t}),$$

which we integrate for initial condition $\phi_H^0 = \hat{R}_0 = \text{Id.}$

We now want to extend analytically γ to a complex neighbourhood of its domain. Take $\widetilde{\gamma}(\boldsymbol{z}) = R_{\boldsymbol{z}} \circ \Gamma(G \circ R_{\boldsymbol{z}})(0), \ \boldsymbol{z} \in D_{\eta}$. The maps $\boldsymbol{z} \mapsto G \circ R_{\boldsymbol{z}}$ and $H \mapsto \Gamma(H)$ are analytic and $C^1_{per}(\mathbb{R}^d, \mathbb{C}^{2d}) \ni g \mapsto g(0)$ is bounded. As $\widetilde{\gamma} \colon D_\eta \to \mathbb{C}^{2d}$ involves their composition, it is analytic and \mathbb{Z}^d -periodic. From (4.94), for any $\boldsymbol{x} \in \mathbb{R}^d$, we have

$$\widetilde{\gamma}(\boldsymbol{x}) = \Gamma(G) \circ \widehat{R}_{\boldsymbol{x}}(0) = \Gamma(G)(\boldsymbol{x}) = \gamma(\boldsymbol{x}).$$

Finally, since Γ is analytic, the same is true for the map $H \mapsto \gamma$.

Remark 4.17. A quasiperiodic invariant torus \mathcal{T} is always Lagrangian. If in addition \mathcal{T} is homotopic to $\{\boldsymbol{y} = \boldsymbol{a}\}$, it is the graph of a function $\psi \in C^{\omega}(\mathbb{T}^d, D)$ (see [12]).

APPENDIX A. ELIMINATION OF MODES

A.1. Homotopy method for vector fields. In this section we prove Theorem 3.7 using the homotopy method (cf. [20]). As n is fixed, we will drop it from our notations. In addition we write $\rho' = \rho_n$ and $\rho = \rho_n + \nu$. We will be using the symbol D_x for the derivative with respect to \boldsymbol{x} .

Firstly, we include a technical lemma that will be used in the following.

Lemma A.1. Let $f \in \mathcal{A}'_{\rho,r}$. If $U = \mathrm{Id} + (u, 0)$ where $u: D_{\rho'} \times B_r \to D_{(\rho - \rho')/2}$ is in $\mathcal{A}_{\rho',r}$ and $||u||_{\rho',r} < (\rho - \rho')/4\pi$, then

- $||f \circ U||_{\rho',r} \le ||f||_{(\rho+\rho')/2,r},$ $||D_{\boldsymbol{x}}f \circ U|| \le ||f||'_{(\rho+\rho')/2,r},$

- $||f \circ U f||_{\rho',r} \le ||f||'_{(\rho+\rho')/2,r} ||u||_{\rho',r},$ $||D_{\boldsymbol{x}}f \circ U D_{\boldsymbol{x}}f|| \le \frac{4\pi}{\rho-\rho'} ||f||'_{\rho,r} ||u||_{\rho',r}.$

The proof of these inequalities is straightforward and thus will be omitted. Now, assume that

$$\delta = 42\varepsilon/\sigma < 1/2.$$

For vector fields in the form $X = \boldsymbol{\omega} + \pi_2 + f$, where $\pi_2 \colon (\boldsymbol{x}, \boldsymbol{y}) \mapsto \boldsymbol{y}$ is seen as a function in $\mathcal{A}'_{\rho,r}$, consider f to be in the open ball in $\mathcal{A}'_{\rho,r}$ centred at the origin with radius ε . The coordinate transformation U is written as $U = \mathrm{Id} + (u, 0)$, with u in

$$\mathcal{B} = \left\{ u \in \mathbb{I}^- \mathcal{A}'_{\rho',r} \colon u \colon D_{\rho'} \times B_r \to D_\rho, \|u\|'_{\rho',r} < \delta \right\}.$$

Notice that we have

$$\mathbb{I}^{-}U^{*}(X) = \mathbb{I}^{-}(DU)^{-1}(\boldsymbol{\omega} + \pi_{2} + f \circ U, 0)$$

= $(\mathbb{I}^{-}(I + D_{\boldsymbol{x}}u)^{-1}(\boldsymbol{\omega} + \pi_{2} + f \circ U), 0)$

From now on the parameter r is omitted whenever there is no ambiguity. Define the operator $F: \mathcal{B} \to \mathbb{I}^- \mathcal{A}_{\rho'},$

$$F(u) = \mathbb{I}^{-}(I + D_{\boldsymbol{x}}u)^{-1}(\boldsymbol{\omega} + \pi_2 + f \circ U).$$
(A.1)

F(u) takes real values for real arguments whenever u has that property. It is easy to see that the derivative of F at u is the linear map from $\mathbb{I}^-\mathcal{A}'_{o'}$ to $\mathbb{I}^-\mathcal{A}_{o'}$:

$$DF(u)h = \mathbb{I}^{-}(I + D_{\boldsymbol{x}}u)^{-1}[D_{\boldsymbol{x}}f \circ Uh$$

$$-D_{\boldsymbol{x}}h(I + D_{\boldsymbol{x}}u)^{-1}(\boldsymbol{\omega} + \pi_{2} + f \circ U)].$$
(A.2)

We want to find a solution of

$$F(u_t) = (1 - t)F(u_0), \tag{A.3}$$

with $0 \le t \le 1$ and "initial" condition $u_0 = 0$. Differentiating the above equation with respect to t, we get

$$DF(u_t)\frac{du_t}{dt} = -F(0).$$
(A.4)

Proposition A.2. If $u \in \mathcal{B}$, then $DF(u)^{-1}$ is a bounded linear operator from $\mathbb{I}^-\mathcal{A}_{\rho'}$ to $\mathbb{I}^-\mathcal{A}'_{\rho'}$ and

$$\|DF(u)^{-1}\| < \delta/\varepsilon.$$

From the above proposition (to be proved in Section A.1.1) we integrate (A.4) with respect to t, obtaining the integral equation:

$$u_t = -\int_0^t DF(u_s)^{-1} F(0) \, ds. \tag{A.5}$$

In order to check that $u_t \in \mathcal{B}$ for any $0 \leq t \leq 1$, we estimate its norm:

$$\|u_t\|'_{\rho'} \leq t \sup_{v \in \mathcal{B}} \|DF(v)^{-1}F(0)\|'_{\rho'} \\ \leq t \sup_{v \in \mathcal{B}} \|DF(v)^{-1}\| \|\mathbb{I}^-f\|_{\rho'} < t\delta \|f\|_{\rho'}/\varepsilon,$$

so, $||u_t||'_{\rho'} < \delta$. Therefore, the solution of (A.3) exists in \mathcal{B} and is given by (A.5). Moreover, if X is real-analytic, then u_t takes real values for real arguments.

It is now easy to see that

$$U_t^* X - X^0 = \mathbb{I}^+ \sum_{n \ge 2} (-D(U_t - \mathrm{Id}))^n X^0 + \mathbb{I}^+ U_t^* f + (1 - t) \mathbb{I}^- f.$$

So, using Lemma A.1,

$$\begin{split} \|U_t^*X - X^0\|_{\rho'} &\leq \frac{1}{1 - \|u_t\|_{\rho'}'} (\|\boldsymbol{\omega}\| \|u_t\|_{\rho'}'^2 + \|f\|_{\rho}) + (1 - t)\|f\|_{\rho'} \\ &< \frac{1}{1 - \delta} \left(\delta^2 \|\boldsymbol{\omega}\| \|f\|_{\rho'}/\varepsilon^2 + 1\right) \|f\|_{\rho} + (1 - t)\|f\|_{\rho'} \\ &< \left[\frac{1}{1 - \delta} \left(\frac{\delta^2 \|\boldsymbol{\omega}\|}{\varepsilon} + 1\right) + 1 - t\right] \|f\|_{\rho'}'. \end{split}$$

Moreover, $||U_t^*X - X^0 - \mathbb{I}^+ f - (1-t)\mathbb{I}^- f||_{\rho'} = \mathcal{O}(||f||_{\rho}^2)$, hence the derivative of $X \mapsto U_t^*X$ at X^0 is $\mathbb{I} - t\mathbb{I}^-$.

A.1.1. Proof of Proposition A.2.

Lemma A.3. If $||f||'_{\rho} < \varepsilon < \sigma/6$, then

$$DF(0)^{-1} \colon \mathbb{I}^- \mathcal{A}_{\rho'} \to \mathbb{I}^- \mathcal{A}'_{\rho'}$$

is continuous and

$$\|DF(0)^{-1}\| < \frac{3}{\sigma - 6\|f\|_{\rho}'}.$$

Proof. From (A.2) one has

$$DF(0) h = \mathbb{I}^{-}(\widehat{f} - D_{\omega}) h$$

= $-\left(\mathbb{I} - \mathbb{I}^{-}\widehat{f} D_{\omega}^{-1}\right) D_{\omega} h,$

where $\widehat{f}h = Dfh - Dhf$ and $D_{\omega}h = D_{x}h(\omega + \pi_2)$. Thus, the inverse of this operator, if it exists, is given by

$$DF(0)^{-1} = -D_{\omega}^{-1} \left(\mathbb{I} - \mathbb{I}^{-} \widehat{f} D_{\omega}^{-1} \right)^{-1}$$

The inverse of $D_{\boldsymbol{\omega}}$ is the linear map from $\mathbb{I}^- \mathcal{A}_{\rho'}$ to $\mathbb{I}^- \mathcal{A}'_{\rho'}$:

$$D_{\boldsymbol{\omega}}^{-1} g(\boldsymbol{x}, \boldsymbol{y}) = \sum_{\boldsymbol{k} \in I^{-}} \frac{g_{\boldsymbol{k}}(\boldsymbol{y})}{2\pi \mathrm{i} \boldsymbol{k} \cdot X^{0}(\boldsymbol{y})} \mathrm{e}^{2\pi \mathrm{i} \boldsymbol{k} \cdot \boldsymbol{x}},$$

and is well-defined since Lemma 3.2 implies that

$$|\boldsymbol{k} \cdot X^0(\boldsymbol{y})| > \sigma \|\boldsymbol{k}\|/2,$$

with $\boldsymbol{k} \in I^-$ and $\boldsymbol{y} \in B_r$. So,

$$\begin{split} \|D_{\boldsymbol{\omega}}^{-1}g\|_{\rho'}^{\prime} &= \sum_{\boldsymbol{k}\in I^{-}} \frac{1+2\pi\|\boldsymbol{k}\|}{2\pi} \sup_{\boldsymbol{y}\in B_{r}} \left\|\frac{g_{\boldsymbol{k}}(\boldsymbol{y})}{\boldsymbol{k}\cdot X^{0}(\boldsymbol{y})}\right\| e^{\rho'\|\boldsymbol{k}\|} \\ &< \sum_{\boldsymbol{k}\in I^{-}} \frac{1+2\pi\|\boldsymbol{k}\|}{\pi\sigma\|\boldsymbol{k}\|} \|g_{\boldsymbol{k}}\|_{r} e^{\rho'\|\boldsymbol{k}\|} \\ &\leq \frac{3}{\sigma} \|g\|_{\rho'}. \end{split}$$

Hence, $||D_{\boldsymbol{\omega}}^{-1}|| < 3/\sigma$. It is possible to bound from above the norm of $\widehat{f} \colon \mathbb{I}^- \mathcal{A}'_{\rho'} \to \mathcal{A}_{\rho'}$ by $||\widehat{f}|| \leq 2||f||'_{\rho'}$. Therefore,

$$\|\mathbb{I}^{-}\widehat{f} D_{\omega}^{-1}\| < \frac{6}{\sigma} \|f\|_{\rho'}' < 1,$$

and

$$\left\| \left(\mathbb{I} - \mathbb{I}^{-} \widehat{f} D_{\boldsymbol{\omega}}^{-1} \right)^{-1} \right\| < \frac{\sigma}{\sigma - 6 \|f\|_{\rho'}'}.$$

The statement of the lemma is now immediate.

Lemma A.4. Given $u \in \mathcal{B}$, the linear operator DF(u) - DF(0) mapping $\mathbb{I}^-\mathcal{A}'_{\rho'}$ into $\mathbb{I}^-\mathcal{A}_{\rho'}$, is bounded and

$$\|DF(u) - DF(0)\| < \frac{\|u\|'_{\rho'}}{1 - \|u\|'_{\rho'}} \left[\left(\frac{4\pi}{\rho - \rho'} + \frac{4 - 2\|u\|'_{\rho'}}{1 - \|u\|'_{\rho'}} \right) \|f\|'_{\rho} + \frac{2 - \|u\|'_{\rho'}}{1 - \|u\|'_{\rho'}} \|\boldsymbol{\omega} + \pi_2\| \right].$$

Proof. The formula (A.2) gives

$$\begin{aligned} [DF(u) - DF(0)] h &= \mathbb{I}^{-}(I + D_{x}u)^{-1} [D_{x}f \circ Uh - (I + D_{x}u)D_{x}f h] \\ &- D_{x}h (I + D_{x}u)^{-1} (\omega + \pi_{2} + f) \circ U \\ &+ (I + D_{x}u)D_{x}h (\omega + \pi_{2} + f)] \\ &= \mathbb{I}^{-}(I + D_{x}u)^{-1} \{A + B + C\}, \end{aligned}$$

where

$$A = [D_{\boldsymbol{x}} f \circ U - D_{\boldsymbol{x}} f - D_{\boldsymbol{x}} u D_{\boldsymbol{x}} f] h$$

$$B = D_{\boldsymbol{x}} u D_{\boldsymbol{x}} h (\boldsymbol{\omega} + \pi_2 + f)$$

$$C = -D_{\boldsymbol{x}} h (I + D_{\boldsymbol{x}} u)^{-1} [f \circ U - f - D_{\boldsymbol{x}} u (\boldsymbol{\omega} + \pi_2 + f)].$$

Using Lemma A.1,

$$\begin{aligned} \|A\|_{\rho'} &\leq \left(\frac{4\pi}{\rho-\rho'}\|f\|_{\rho}'\|u\|_{\rho'}+\|f\|_{\rho'}'\|u\|_{\rho'}'\right)\|h\|_{\rho'},\\ \|B\|_{\rho'} &\leq \left(\|\omega+\pi_2\|+\|f\|_{\rho'}\right)\|u\|_{\rho'}'\|h\|_{\rho'}',\\ \|C\|_{\rho'} &\leq \frac{1}{1-\|u\|_{\rho'}'}\left[\|f\|_{(\rho+\rho')/2}'\|u\|_{\rho'}+\|u\|_{\rho'}'\left(\|\omega+\pi_2\|_r+\|f\|_{\rho'}\right)\right]\|h\|_{\rho'}'.\end{aligned}$$

To conclude the proof of Proposition A.2, notice that

$$\begin{aligned} \|DF(u)^{-1}\| &\leq \left(\|DF(0)^{-1}\|^{-1} - \|DF(u) - DF(0)\|\right)^{-1} \\ &< \left\{\frac{\sigma}{3} - 2\varepsilon - \frac{\delta}{1 - \delta} \left[\left(\frac{4\pi}{\rho - \rho'} + \frac{4 - 2\delta}{1 - \delta}\right)\varepsilon + \frac{2 - \delta}{1 - \delta}\|\boldsymbol{\omega} + \pi_2\|_r\right]\right\}^{-1} \\ &< \frac{\delta}{\varepsilon}. \end{aligned}$$

The last inequality is true if

$$\varepsilon < \delta \left[\frac{\sigma}{3} - \frac{2\delta}{(1-\delta)^2} \| \boldsymbol{\omega} + \pi_2 \|_r \right] \left[1 + 2\delta + \frac{\delta^2}{1-\delta} \left(\frac{4\pi}{\rho - \rho'} + \frac{4-2\delta}{1-\delta} \right) \right]^{-1}$$

with a positive numerator N and denominator D in the r.h.s. This is so for our choices of ε and $\delta < \frac{1}{2}$, by observing that

$$\|\pi_2\|_r = \sup_{\boldsymbol{y}\in B_r} \|\boldsymbol{y}\| \le a \, |\gamma^{(n)}|^{-1} \|M^{(n)}\| + b\|\boldsymbol{\omega}\| < \frac{1}{2} \|\boldsymbol{\omega}\|, \tag{A.6}$$

thus $\|\boldsymbol{\omega} + \pi_2\|_r < \frac{3}{2} \|\boldsymbol{\omega}\|$ and

$$\frac{2\delta}{(1-\delta)^2} \|\boldsymbol{\omega} + \pi_2\|_r < 12\delta \|\boldsymbol{\omega}\| < \frac{\sigma}{6}$$

So, $N > \delta\sigma/6$, D < 7, and finally $\varepsilon \le \frac{\sigma^2}{42||\boldsymbol{\omega}||} < \frac{\sigma}{42} < N/D$.

A.2. Elimination of modes for Hamiltonians. Here we present a proof of Theorem 4.5. It is similar to related methods appearing in e.g. [16, 1]. As we have fixed n, we will not include it in our notations.

Let $R = (R_1, R_2)$ and $R' = (R'_1, R'_2)$ be such that R > R' > 0 componentwise. We will be interested on the set $\mathcal{G}_{R'}$ of analytic symplectomorphisms $g: \mathcal{D}_{R'} \to \mathcal{D}_R$ satisfying $g - \mathrm{Id} \in \mathcal{A}_{R'}^{2d}$ and

$$||g - \operatorname{Id}||_{R'} < \delta = \min\{(R_1 - R'_1)/2\pi, R_2 - R'_2\}$$

We use the notation $\{\cdot, \cdot\}$ for the usual Poisson bracket associated to $\mathbb{J}: (\boldsymbol{x}, \boldsymbol{y}) \mapsto (\boldsymbol{y}, -\boldsymbol{x})$. In the following $R - \delta$ stands for $R - \delta(1, 1)$ and $\pi_2: (\boldsymbol{x}, \boldsymbol{y}) \mapsto \boldsymbol{y}$ is the projection on the second component.

Lemma A.5. Let $0 < \xi \leq \frac{1}{2}$. If $G \in \mathcal{A}'_{R'}$ and $||G||'_{R'} < \xi \delta/(2\pi + 1)$, then there is a unique analytic symplectomorphism $g: \mathcal{D}_{R'-2\delta} \to \mathbb{C}^{2d}$ such that $||g - \operatorname{Id}||_{R'-2\delta} < \xi \delta$ and

$$g = \mathrm{Id} + \mathbb{J}\nabla G \circ \widehat{g},\tag{A.7}$$

where $\widehat{g}(\boldsymbol{x}, \boldsymbol{y}) = (\boldsymbol{x}, \pi_2 g(\boldsymbol{x}, \boldsymbol{y})), (\boldsymbol{x}, \boldsymbol{y}) \in \mathcal{D}_{R'-2\delta}$. Moreover, for any $H \in \mathcal{A}_{R'}$

$$\|H \circ g\|_{R'-2\delta} \le \|H\|_{R'}$$

$$\|H \circ g - H\|_{R'-2\delta} \le 2\xi \|H\|_{R'}$$

$$\|H \circ g - H - \{H, G\}\|_{R'-2\delta} \le 2\xi^2 \|H\|_{R'}$$

(A.8)

and the maps $G \mapsto g$ and $G \mapsto H \circ g$ are analytic.

Proof. Define the map $T: g \mapsto \operatorname{Id} + \mathbb{J}\nabla G \circ \widehat{g}$ on the open ball B in $\mathcal{G}_{R'-2\delta}$ centred at the identity and with radius $\xi\delta$. It is simple to check that $T(B) \subset B$, in particular T(g) for $g \in B$ is symplectic. We now show that T is a contraction on B and thus its unique fixed point is the map we are looking for. In fact, whenever $g \in B$ we obtain

$$\|DT(g)\| \leq \|D\nabla G \circ \widehat{g}\|_{R'-2\delta} \leq \|D\nabla G\|_{R'-\delta}$$

$$\leq \frac{2\pi+1}{\delta} \|\nabla G\|_{R'} \leq \frac{2\pi+1}{\delta} \|G\|'_{R'} < \xi.$$
(A.9)

For the estimates in (A.8) (the first is now immediate) we introduce the differentiable function

$$f: \{z \in \mathbb{C}: |z| < \zeta\} \to \mathcal{A}_{R'}$$

$$z \mapsto H \circ (\mathrm{Id} + z \mathbb{J} \nabla G (\mathrm{Id} + z (\widehat{g} - \mathrm{Id})))$$
(A.10)

where $\zeta = 1/\xi \ge 2$. Cauchy's integral formula yields that

$$\begin{aligned} \|H \circ g - H\|_{R'-2\delta} &= \|f(1) - f(0)\|_{R'-2\delta} \\ &\leq \frac{1}{2\pi} \oint_{|z|=\zeta} \frac{\|f(z)\|_{R'-2\delta}}{|z(z-1)|} dz \\ &\leq \frac{1}{\zeta - 1} \sup_{|z|=\zeta} \|f(z)\|_{R'-2\delta} \leq 2\xi \|H\|_{R'}. \end{aligned}$$
(A.11)

and

$$\begin{aligned} \|H \circ g - H - \{H, G\}\|_{R'-2\delta} &= \|f(1) - f(0) - f'(0)\|_{R'-2\delta} \\ &\leq \frac{1}{2\pi} \oint_{|z|=\zeta} \frac{\|f(z)\|_{R'-2\delta}}{|z^2(z-1)|} dz \\ &\leq \frac{1}{\zeta(\zeta-1)} \sup_{|z|=\zeta} \|f(z)\|_{R'-2\delta} \leq 2\xi^2 \|H\|_{R'}. \end{aligned}$$
(A.12)

By the implicit function theorem the maps $G \mapsto g$ and $G \mapsto H \circ g$ are analytic. \Box Lemma A.6. Let $\sigma > 2R_2 ||Q||$, $\varepsilon' > 0$ and $H \in \mathcal{A}'_R$ such that

$$||H - H^0||_R < \varepsilon' \le \frac{\sigma\delta}{(2\pi + 1)[1 + 2\pi + (\tau + 1)/R_2]}.$$
 (A.13)

Then there is $G \in \mathbb{I}^-\mathcal{A}'_{R'}$ such that

$$\mathbb{I}^{-}(H + \{H, G\}) = 0 \quad and \quad \|G\|'_{R'} \le \frac{\delta}{(2\pi + 1)\varepsilon'} \|\mathbb{I}^{-}H\|_{R'}.$$
(A.14)

Moreover, the map $H \mapsto G$ is analytic.

Proof. Consider the linear operator associated to H:

$$\mathcal{F} \colon \mathbb{I}^{-}\mathcal{A}'_{R'} \to \mathbb{I}^{-}\mathcal{A}_{R'}, \qquad K \mapsto \mathbb{I}^{-}\{H, K\}.$$
(A.15)

It is well-defined since

$$\|\mathbb{I}^{-}\{H,K\}\|_{R'} \leq \|\nabla H\|_{R'} \|\nabla K\|_{R'}$$
$$\leq \|H\|'_{R'} \|K\|'_{R'}.$$

We will show that $\mathcal{F}^{-1} \colon \mathbb{I}^- \mathcal{A}_{R'} \to \mathbb{I}^- \mathcal{A}'_{R'}$ is bounded and

$$\|\mathcal{F}^{-1}\| < \frac{1}{\frac{\pi R_2 \sigma}{(2\pi+1)R_2 + \tau + 1} - 2\frac{2\pi+1}{\delta}\varepsilon'} \le \frac{\delta}{(2\pi+1)\varepsilon'}.$$
 (A.16)

We start by decomposing any Hamiltonian $H = H^0 + F$ as

$$H(\boldsymbol{x}, \boldsymbol{y}) = \sum_{\boldsymbol{k}} H_{\boldsymbol{k}}(\boldsymbol{y}) e^{2\pi i \boldsymbol{k} \cdot \boldsymbol{x}} \text{ with } H_{\boldsymbol{k}}(\boldsymbol{y}) = \sum_{\boldsymbol{\nu}} H_{\boldsymbol{k}, \boldsymbol{\nu}} \boldsymbol{y}^{\boldsymbol{\nu}}.$$

Write $D_0 = \nabla_2 H^0 \cdot \nabla_1$, with ∇_1 and ∇_2 standing for the derivatives with respect to \boldsymbol{x} and \boldsymbol{y} . The definition of \mathcal{F} in (A.15) yields

$$\mathcal{F}(K) = \mathbb{I}^{-}(\widehat{F} - D_0) K = -\left(\mathbb{I} - \mathbb{I}^{-}\widehat{F}D_0^{-1}\right) D_0 K,$$

where $\widehat{F}(K) = \{F, K\}$. If the inverse of \mathcal{F} exists is given by

$$\mathcal{F}^{-1} = -D_0^{-1} \left(\mathbb{I} - \mathbb{I}^- \widehat{F} D_0^{-1} \right)^{-1}.$$
 (A.17)

The map $D_0^{-1} \colon \mathbb{I}^- \mathcal{A}_{R'} \to \mathbb{I}^- \mathcal{A}'_{R'}$ is linear and given by

$$D_0^{-1} W(\boldsymbol{x}, \boldsymbol{y}) = \sum_{\boldsymbol{k} \in \mathbb{Z}^d - \{0\}} \frac{W_{\boldsymbol{k}}(\boldsymbol{y})}{2\pi \mathrm{i}(\boldsymbol{k} \cdot \nabla_2 H^0(\boldsymbol{y}))} \mathrm{e}^{2\pi \mathrm{i} \boldsymbol{k} \cdot \boldsymbol{x}}, \quad W \in \mathbb{I}^- \mathcal{A}_{R'}.$$

For each $\boldsymbol{k} \in I^-$, using (4.33) and $\|Q\| < \sigma/(2R_2)$ thus $|\boldsymbol{k} \cdot Q\boldsymbol{y}/\boldsymbol{k} \cdot \boldsymbol{\omega}| < 1/2$,

$$\frac{W_{\boldsymbol{k}}(\boldsymbol{y})}{\boldsymbol{k}\cdot\boldsymbol{\omega}\left(1+\frac{\boldsymbol{k}\cdot\boldsymbol{Q}\boldsymbol{y}}{\boldsymbol{k}\cdot\boldsymbol{\omega}}\right)} = \frac{W_{\boldsymbol{k}}(\boldsymbol{y})}{\boldsymbol{k}\cdot\boldsymbol{\omega}}\sum_{n\geq 0}\left(-\frac{\boldsymbol{k}\cdot\boldsymbol{Q}\boldsymbol{y}}{\boldsymbol{k}\cdot\boldsymbol{\omega}}\right)^{n},\tag{A.18}$$

we get the estimate

$$\left\| \frac{W_{\boldsymbol{k}}}{\boldsymbol{k} \cdot \nabla_{2} H^{0}} \right\|_{R_{2}} \leq \sum_{n \geq 0} \sum_{\boldsymbol{\nu}} \frac{|W_{\boldsymbol{k},\boldsymbol{\nu}}| R_{2}^{\|\boldsymbol{\nu}\|} \|Q\|^{n} R_{2}^{n}}{\sigma^{n+1} \|\boldsymbol{k}\|}$$
$$< \sum_{n \geq 0} \sum_{\boldsymbol{\nu}} \frac{|W_{\boldsymbol{k},\boldsymbol{\nu}}| R_{2}^{\|\boldsymbol{\nu}\|}}{\sigma \|\boldsymbol{k}\|} \left(\frac{1}{2}\right)^{n}$$
$$= \frac{2}{\sigma \|\boldsymbol{k}\|} \|W_{\boldsymbol{k}}\|_{R_{2}}.$$
(A.19)

Similarly, we find the bound

$$\left\| \frac{\nabla_2 W_{\boldsymbol{k}}}{\boldsymbol{k} \cdot \nabla_2 H^0} \right\|_{R_2} \le \sum_{\boldsymbol{\nu}} \frac{2 \|\boldsymbol{\nu}\| \|W_{\boldsymbol{k},\boldsymbol{\nu}} |R_2^{\|\boldsymbol{\nu}\|-1}}{\sigma \|\boldsymbol{k}\|} < \frac{2\tau}{\sigma R_2} \|W_{\boldsymbol{k}}\|_{R_2}.$$
(A.20)

Finally,

$$\left\| \frac{W_{\boldsymbol{k}}Q\boldsymbol{k}}{(\boldsymbol{k}\cdot\nabla_{2}H^{0})^{2}} \right\|_{R_{2}} < \frac{2}{\sigma R_{2}\|\boldsymbol{k}\|} \|W_{\boldsymbol{k}}\|_{R_{2}}.$$
(A.21)

It is now immediate to see that

$$\|D_0^{-1}W\|_{R'} \le \frac{2}{2\pi\sigma} \|W\|_{R'}$$
, and $\|\nabla_1(D_0^{-1}W)\|_{R'} \le \frac{2}{\sigma} \|W\|_{R'}$.

Moreover,

$$\nabla_2 \left(\frac{W_{\boldsymbol{k}}(\boldsymbol{y})}{\boldsymbol{k} \cdot \nabla_2 H^0(\boldsymbol{y})} \right) = \frac{\nabla_2 W_{\boldsymbol{k}}(\boldsymbol{y})}{\boldsymbol{k} \cdot \nabla_2 H^0(\boldsymbol{y})} - \frac{W_{\boldsymbol{k}}(\boldsymbol{y}) Q \boldsymbol{k}}{(\boldsymbol{k} \cdot \nabla_2 H^0(\boldsymbol{y}))^2}$$

which implies

$$\|\nabla_2(D_0^{-1}W)\|_{R'} < \frac{\tau+1}{\pi\sigma R_2} \|W\|_{R'}.$$
(A.22)

Hence,

$$\|D_0^{-1}\| < \frac{2}{\sigma} \left(1 + \frac{1}{2\pi} + \frac{\tau + 1}{2\pi R_2}\right)$$

As $\widehat{F} \colon \mathbb{I}^- \mathcal{A}'_{R'} \to \mathcal{A}_{R'}$ with $\|\widehat{F}\| \leq 2 \|\nabla F\|_{R'} \leq 2\frac{2\pi+1}{\delta} \|F\|_R$ (by Cauchy's estimate),

$$\|\mathbb{I}^{-}\widehat{F} D_{0}^{-1}\| < \frac{4}{\sigma} \left(1 + \frac{1}{2\pi} + \frac{\tau + 1}{2\pi R_{2}}\right) \|\nabla F\|_{R'} < 1,$$

and

$$\left\| \left(\mathbb{I} - \mathbb{I}^{-} \widehat{F} D_{0}^{-1} \right)^{-1} \right\| < \left[1 - \frac{4}{\sigma} \left(1 + \frac{1}{2\pi} + \frac{\tau + 1}{2\pi R_{2}} \right) \| \nabla F \|_{R'} \right]^{-1}$$

Thus \mathcal{F}^{-1} exists given by (A.17) and the estimate (A.16) on its norm follows immediately.

ately. A solution of (A.14) is simply given by $G = \mathcal{F}^{-1}(-\mathbb{I}^-H)$. Therefore, $\|G\|'_{R'} \leq \|\mathcal{F}^{-1}\| \|\mathbb{I}^-H\|_{R'}$.

Consider the pairs $R = (\rho_n + \nu, r')$ and $R' = (\rho_n, r)$, $\sigma > 2r' ||Q||$ and $H_0 = H$ as given in Theorem 4.5. We are going to iterate the procedure indicated in the previous lemmas. Let a sequence of Hamiltonians be given by

$$H_k = H_{k-1} \circ g_k, \quad k \in \mathbb{N},$$

where G_k and g_k are determined for H_{k-1} by Lemmas A.6 and A.5, respectively. In addition, denote by

$$g^{(k)} = g_1 \circ \dots \circ g_k \tag{A.23}$$

the composition of all symplectomorphisms up to the kth-step so that $H_k = H \circ g^{(k)}$. In order to determine the right domains of H_k , G_k and g_k , define the sequences

$$R_{k} = R_{k-1} - 4\delta_{k} = R - 4\sum_{i=1}^{k} \delta_{i},$$

$$R'_{k} = R_{k-1} - \delta_{k},$$
(A.24)

with $R_0 = R$ and

$$\delta_k = \frac{1}{2^k} \min\left\{1, \frac{\nu}{2\pi}, r' - r\right\} \le \frac{1}{2^k}.$$
(A.25)

From now on, assume that

$$\varepsilon' = \min\left\{\frac{1}{2} \|H^0\|_R, \frac{\sigma\delta_1}{(2\pi+1)(1+2\pi+\frac{\tau+1}{r'})}\right\}.$$
 (A.26)

Lemma A.7. If for every $k \in \mathbb{N}$, $\|\mathbb{I}^- H_{k-1}\|_{R_{k-1}} \leq \varepsilon'/2$ and

$$||G_k||'_{R'_k} < \frac{\delta_k}{(2\pi+1)\varepsilon'} ||\mathbb{I}^- H_{k-1}||_{R_{k-1}}$$

then $g_k(\mathcal{D}_{R_k}) \subset \mathcal{D}_{R_{k-1}}$ and

$$\|g^{(k)} - \operatorname{Id}\|_{R_{k}} \leq \sum_{i=1}^{k} \frac{\delta_{i}}{\varepsilon'} \|\mathbb{I}^{-}H_{i-1}\|_{R_{i-1}}$$

$$\|g^{(k)} - g^{(k-1)}\|_{R_{k}} \leq \frac{3}{2\varepsilon'} \|\mathbb{I}^{-}H_{k-1}\|_{R_{k-1}}.$$
(A.27)

Proof. Recall Lemma A.5 for $\xi = \|\mathbb{I}^- H_{k-1}\|_{R_{k-1}}/\varepsilon'$ and check that

$$||g_k - \mathrm{Id}||_{R_k} \le ||g_k - \mathrm{Id}||_{R'_k - 2\delta_k} < \frac{\delta_k}{\varepsilon'} ||\mathbb{I}^- H_{k-1}||_{R_{k-1}}$$

and $R_k + \delta_k < R_{k-1}$ componentwise. Now,

$$g^{(k)} - \mathrm{Id} = \sum_{i=1}^{k} g_i \circ \cdots \circ g_k - \sum_{i=2}^{k} g_i \circ \cdots \circ g_k - \mathrm{Id}$$
$$= \sum_{i=1}^{k-1} (g_i \circ \cdots \circ g_k - g_{i+1} \circ \cdots \circ g_k) + g_k - \mathrm{Id}$$
$$= \sum_{i=1}^{k-1} (g_i - \mathrm{Id}) \circ g_{i+1} \circ \cdots \circ g_k + g_k - \mathrm{Id}.$$
(A.28)

Thus,

$$\|g^{(k)} - \operatorname{Id}\|_{R_k} \le \sum_{i=1}^k \|g_i - \operatorname{Id}\|_{R_i} \le \sum_{i=1}^k \frac{\delta_i}{\varepsilon'} \|\mathbb{I}^- H_{i-1}\|_{R_{i-1}}.$$
 (A.29)

Furthermore, as

$$g^{(k)} - g^{(k-1)} = (g^{(k-1)} - \mathrm{Id}) \circ g_k - (g^{(k-1)} - \mathrm{Id}) + (g_k - \mathrm{Id})$$
 (A.30)

we get

$$\begin{aligned} \|g^{(k)} - g^{(k-1)}\|_{R_{k}} &\leq (\|Dg^{(k-1)} - I\|_{R_{k}} + 1) \|g_{k} - \operatorname{Id}\|_{R_{k}} \\ &\leq \frac{\delta_{k}}{\varepsilon'} \|\mathbb{I}^{-}H_{k-1}\|_{R_{k-1}} \left(\frac{2\pi + 1}{4\delta_{k}} \sum_{i=1}^{k-1} \frac{\delta_{i}}{\varepsilon'} \|\mathbb{I}^{-}H_{i-1}\|_{R_{i-1}} + 1\right) \\ &\leq \frac{1}{\varepsilon'} \left(1 + \frac{1}{2}\right) \|\mathbb{I}^{-}H_{k-1}\|_{R_{k-1}}. \end{aligned}$$
(A.31)

Notice that since $\varepsilon' \leq \frac{1}{2} \|H^0\|_R$, we have

$$\varepsilon' \le \|H^0\|_R - \varepsilon' \le \|H\|_R \le \|H^0\|_R + \varepsilon' \tag{A.32}$$

and also

$$\frac{1}{2} \|H^0\|_R \le \|H\|_R \le \frac{3}{2} \|H^0\|_R.$$
(A.33)

Lemma A.8. For any $k \in \mathbb{N}$, if $\|\mathbb{I}^- H\|_R \leq {\varepsilon'}^2/(8\|H\|_R)$, then

$$\|\mathbb{I}^{-}H_{k}\|_{R_{k}} \leq \left(\frac{4\|H\|_{R}}{\varepsilon'^{2}}\right)^{2^{k}-1} \|\mathbb{I}^{-}H\|_{R}^{2^{k}} \leq \frac{\varepsilon'}{2}, \qquad (A.34)$$

$$|H_{k} - H_{k-1}||_{R_{k}} \leq \frac{4||H||_{R}}{\varepsilon'} ||\mathbb{I}^{-}H_{k-1}||_{R_{k-1}},$$
(A.35)

$$\|H_k\|_{R_k} \leq 2\|H\|_{R}. \tag{A.36}$$

Proof. We will prove the above inequalities by induction. The generating Hamiltonian G_1 given by Lemma A.6 and the symplectomorphism g_1 by Lemma A.5 satisfy $||G_1||'_{R'_1} \leq \delta_1 ||\mathbb{I}^- H|| / [(2\pi + 1)\varepsilon'], ||g_1 - \operatorname{Id}||_{R'_1 - 2\delta_1} < ||\mathbb{I}^- H||_R \delta_1 / \varepsilon'$ and $\mathbb{I}^- H_1 = \mathbb{I}^- H \circ g_1 - \mathbb{I}^- (H + \{H, G_1\})$. Hence,

$$\|\mathbb{I}^{-}H_{1}\|_{R_{1}} \leq \|H \circ g_{1} - H_{k} - \{H, G_{1}\}\|_{R_{1}} \leq 2\left(\frac{\|\mathbb{I}^{-}H\|_{R}}{\varepsilon'}\right)^{2} \|H\|_{R}.$$
 (A.37)

and

$$\|H_1 - H\|_{R_1} \le \|\nabla H\|_{R_1} \|g_1 - \operatorname{Id}\|_{R_1} \le \frac{2\pi + 1}{4\varepsilon'} \|\mathbb{I}^- H\|_R \|H\|_R \le \frac{2}{\varepsilon'} \|\mathbb{I}^- H\|_R \|H\|_R.$$
(A.38)

Thus, (A.34) and (A.35) are valid for k = 1 and so is (A.36) because $||H_1||_{R_1} \le ||H_1 - H||_{R_1} + ||H||_{R_1}$.

Now, assume that the inequalities are true for k. Under these conditions, Lemma A.6 guarantees the existence of G_{k+1} so that

$$\|G_{k+1}\|'_{R_{k+1}} \le \frac{\delta_{k+1}}{(2\pi+1)\varepsilon'} \|\mathbb{I}^- H_k\|_{R_k}$$
(A.39)

and Lemma A.5 yields g_{k+1} . Therefore, $\mathbb{I}^- H_{k+1} = \mathbb{I}^- H_k \circ g_{k+1} - \mathbb{I}^- (H_k + \{H_k, G_{k+1}\})$ and

$$\|\mathbb{I}^{-}H_{k+1}\|_{R_{k+1}} \leq \|H_{k} \circ g_{k+1} - H_{k} - \{H_{k}, G_{k+1}\}\|_{R_{k+1}}$$

$$\leq 2\left(\frac{\|\mathbb{I}^{-}H_{k}\|_{R_{k}}}{\varepsilon'}\right)^{2} \|H_{k}\|_{R_{k}}$$

$$\leq \left(\frac{4\|H\|_{R}}{\varepsilon'^{2}}\right)^{2^{k+1}-1} \|\mathbb{I}^{-}H\|_{R}^{2^{k+1}}.$$
(A.40)

Similarly,

$$\begin{aligned} \|H_{k+1} - H_k\|_{R_{k+1}} &\leq \|\nabla H_k\|_{R_{k+1}} \|g_{k+1} - \operatorname{Id}\|_{R_{k+1}} \\ &\leq \frac{2\pi + 1}{4\delta_{k+1}\varepsilon'} \|\mathbb{I}^- H_k\|_{R_k} \delta_{k+1} \|H_k\|_{R_k} \\ &\leq \frac{4}{\varepsilon'} \|\mathbb{I}^- H_k\|_{R_k} \|H\|_{R}. \end{aligned}$$
(A.41)

Finally, making use of the above inequality,

$$\begin{aligned} \|H_{k+1}\|_{R_{k+1}} &\leq \|H\|_{R} + \sum_{i=1}^{k+1} \|H_{i} - H_{i-1}\|_{R_{k+1}} \\ &\leq \|H\|_{R} + \frac{4\|H\|_{R}}{\varepsilon'} \sum_{i=1}^{k+1} \|\mathbb{I}^{-}H_{i-1}\|_{R_{i-1}} \\ &\leq \|H\|_{R} + \|H\|_{R} \sum_{i=1}^{k+1} \left(\frac{4\|H\|_{R}\|\mathbb{I}^{-}H\|_{R}}{\varepsilon'^{2}}\right)^{2^{i-1}} \\ &\leq \left(1 + \frac{1}{2} + \sum_{i=1}^{k} \frac{1}{4^{i}}\right) \|H\|_{R} < 2\|H\|_{R}. \end{aligned}$$
(A.42)

Theorem 4.5 will now be a consequence of the following result.

Theorem A.9. If

$$||H - H^0||_R < \varepsilon = \frac{\varepsilon'^2}{12||H^0||_R} \le \frac{\varepsilon'^2}{8||H||_R},$$
 (A.43)

then there exists $g = \lim_{k \to +\infty} g^{(k)} \in \mathcal{G}_{R'}$ such that $\mathbb{I}^- H \circ g = 0$ on $\mathcal{D}_{R'}$. Furthermore, the maps $\mathfrak{G} \colon H \mapsto g$ and $\mathcal{U} \colon H \mapsto H \circ g$ are analytic, and

$$\|g - \operatorname{Id}\|_{R'} \leq \frac{1}{\varepsilon} \|\mathbb{I}^- H\|_R \tag{A.44}$$

$$||H \circ g - H^0||_{R'} \leq \left(1 + \sqrt{\frac{12||H^0||_R}{\varepsilon}}\right) ||H - H^0||_R.$$
 (A.45)

Proof. Lemmas A.7 and A.8 imply that the sequence $g^{(k)}$ converges to a map $g: \mathcal{D}_{R'} \to \mathcal{D}_R$ which is analytic and symplectic, and $H_{\infty} = \lim_{k \to +\infty} H_k = H \circ g$. Moreover, $\mathbb{I}^- H \circ g = \mathbb{I}^- H_{\infty} = 0$. Since the convergence is uniform, the maps $H \mapsto g$ and $H \mapsto H \circ g$ are analytic.

Notice that

$$\sum_{i=1}^{+\infty} \left(\frac{4\|H\|_R \|\mathbb{I}^- H\|_R}{\varepsilon'^2} \right)^{2^{i-1}} \leq \frac{4\|H\|_R \|\mathbb{I}^- H\|_R}{\varepsilon'^2} + \sum_{i=1}^{+\infty} \left(\frac{4\|H\|_R \|\mathbb{I}^- H\|_R}{\varepsilon'^2} \right)^{2i}$$
$$\leq \left(1 + \frac{16\|H\|_R}{3\varepsilon'^2} \|\mathbb{I}^- H\|_R \right) \frac{4\|H\|_R}{\varepsilon'^2} \|\mathbb{I}^- H\|_R$$
$$\leq \frac{20\|H\|_R}{3\varepsilon'^2} \|\mathbb{I}^- H\|_R \leq \frac{1}{\varepsilon} \|\mathbb{I}^- H\|_R.$$
(A.46)

The inequality in (A.44) follows by taking the limit $k \to +\infty$ in (A.27). That is,

$$||g - \operatorname{Id}||_{R'} \le \sum_{i=1}^{+\infty} \frac{\delta_i}{\varepsilon'} ||\mathbb{I}^- H_{i-1}||_{R_{i-1}} \le \frac{1}{\varepsilon} ||\mathbb{I}^- H||_R.$$
 (A.47)

Now,

$$\|H_{\infty} - H^{0}\|_{R_{k}} \leq \|H - H^{0}\|_{R} + \sum_{i=1}^{+\infty} \|H_{i} - H_{i-1}\|_{R_{i}}$$
$$\leq \left(1 + \sqrt{\frac{12\|H^{0}\|_{R}}{\varepsilon}}\right) \|H - H^{0}\|_{R},$$

where we have used Lemma A.8 and the fact that $\|\mathbb{I}^-H\|_R \leq \|H-H^0\|_R$.

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