Lifshitz tails for the Interband Light Absorption Coefficient

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Abstract

In this paper we consider the Interband Light Absorption Coefficient for various models. We show that at the lower and upper edges of the spectrum the Lifshitz tails behaviour of the density of states implies similar behaviour for the ILAC at appropriate energies. The Lifshitz tails property is also exhibited at some points corresponding to the internal band edges of the density of states.

1 Introduction

In this work we look for Lifshitz tails behaviour of the Interband Light Absorption Coefficient (ILAC) defined in equation (eqnl4). The standard definition of the ILAC involves considering a pair of operators of the form $H^{\pm}_{\omega} = \Delta \pm V^{\omega}$, with Δ the Laplacian on either $\ell^2(\mathbb{Z}^d)$, in the discrete case or on $L^2(\mathbb{R}^d)$ in the continuous case, and taking a random potential V^{ω} . Restricting these operators H^{\pm}_{ω} to boxes Λ gives operators with discrete spectra so that in any finite region of energy these operators have only finitely many eigenvalues. Using this fact one can define the quantity

$$\frac{1}{Vol(\Lambda)}\sum_{\lambda_{\omega}^{-}+\lambda_{\omega}^{+}\leq E}|\langle\phi_{\omega,\lambda_{\omega}^{-}},\psi_{\omega,\lambda_{\omega}^{+}}\rangle|^{2}$$

where $\phi_{\omega,\lambda_{\omega}^{-}}, \psi_{\omega,\lambda_{\omega}^{+}}$ are the eigen functions of the operators H_{ω}^{\mp} restricted to the box Λ , corresponding to the eigenvalues $\lambda_{\omega}^{-}, \lambda_{\omega}^{+}$ respectively.

The limit of the above quantity, when it exits, gives the ILAC.

We consider a correlation measure (mentioned also in [12]) ρ and identify the ILAC as the distribution function of a marginal of the measure ρ in a diagonal direction. This identification enables us to prove theorems on the Lifshitz tails behaviour of the ILAC more easily since it involves only comparing the marginal of ρ with the density of states of either of the operators H_{ω}^{\pm} . We also do not need to approximate to define the ILAC, but can obtain the function directly.

In the next section, we present an abstract version of the correlation measure ρ and the density of states n for a pair of random covariant operators and obtain relations between the two.

2 General Covariant Operators

We start with a definition of a random family of self adjoint operators which are covariant under a group action.

Hypotheses 1.

- 1. \mathcal{H} is a (separable, complex) Hilbert space, $(\Omega, \mathcal{F}, \mathbb{P})$ a probability space.
- 2. There is a locally compact abelian group G and $\{U_x\}_{x\in G}$ is a group of unitary operators on \mathcal{H} , i.e. the U_x are unitary and $U_{x+y} = U_x U_y$, $U_0 = \mathrm{Id}, U_{-x} = U_x^{-1} = U_x^*$
- 3. There is a discrete subgroup L of G and an orthogonal projection Pon \mathcal{H} such that $\{U_n^* P U_n\}_{n \in L}$, $\{U_n P U_n^*\}_{n \in L}$ are orthogonal partitions of unity on \mathcal{H} . We set $P_n = U_n^* P U_n$, $\tilde{P}_n = U_n P U_n^*$.
- 4. $\{T_n\}_{n\in L}$ is a group of probability preserving transformations on Ω .

Definition 1. A family $\{A_{\omega}\}_{\omega\in\Omega}$ of self-adjoint operators on \mathcal{H} is called measurable if the family $\{(A_{\omega}+i)^{-1}\}_{\omega\in\Omega}$ is measurable

It is known (see [3], [2] and section 2.4 of [25]) that a family of *bounded* self-adjoint operators is measurable iff it's weakly measurable.

Moreover, if $\{A_{\omega}\}$ is a measurable family of self-adjoint operators then for any bounded measurable function f the operator family $f(A_{\omega})$ is weakly measurable. (also in [3], [2], section 2.4 [25]).

Finally, the product of weakly measurable families is weakly measurable (see [2]).

Definition 2. A weakly measurable family A_{ω} of bounded operators is called covariant (with respect to U_x, T_x) if

$$A_{T_x\omega} = U_x^* A_\omega U_x$$
 for all $x \in G$

Also, a measurable family A_{ω} of self adjoint operators is called covariant (with respect to U_x, T_x) if

$$A_{T_x\omega} = U_x^* A_\omega U_x$$
 for all $x \in G$

If A_{ω} is a covariant family of self-adjoint operators and f is a bounded measurable function, then the family $f(A_{\omega})$ is covariant (also in [3], [2]). Moreover, if both A_{ω} and B_{ω} are covariant families of bounded operators, then $A_{\omega} B_{\omega}$ is a covariant family. We denote by $||B||_1$ the trace norm of a trace class operator B.

Proposition 1. Let A_{ω} and B_{ω} be covariant families of bounded operators and assume that $A_{\omega}P$ and $B_{\omega}P$ are trace class and

$$\mathbb{E}(\|A_{\omega}P\|_{1}) < \infty \text{ and } \mathbb{E}(\|B_{\omega}P\|_{1}) < \infty.$$
(1)

Then:

$$\mathbb{E}(Tr(PA_{\omega}B_{\omega}P)) = \mathbb{E}(Tr(PB_{\omega}A_{\omega}P))$$
(2)

Proof.

$$Tr(PA_{\omega}B_{\omega}P) = Tr(PA_{\omega}B_{\omega}P) \tag{3}$$

$$=\sum_{n} Tr(PA_{\omega}P_{n}B_{\omega}P) \tag{4}$$

since P_n is a partition of unity of orthogonal projections.

$$=\sum_{n} Tr(P_{n}B_{\omega}PA_{\omega}P_{n})$$
(5)

using the Tr(AB) = Tr(BA) and the invariance of trace $Tr(U^*CU) = Tr(C)$,

$$=\sum_{n} Tr(PA_{T_n^{-1}\omega}\tilde{P}_n B_{T_n^{-1}\omega}P)$$
(6)

from the covariance of A_{ω} and B_{ω}

$$= Tr(PB_{T_n^{-1}\omega}A_{T_n^{-1}\omega}P)\rangle \tag{7}$$

In the last step we used the fact that \tilde{P}_n is a partition of unity also. Now we take expectations of either side of the above equation and obtain

$$\mathbb{E}(Tr(PA_{\omega}B_{\omega}P) = \mathbb{E}(\sum_{n} Tr(PB_{T_{n}^{-1}\omega}\tilde{P}_{n}A_{T_{n}^{-1}\omega}P))$$
(8)

$$= \sum_{n} \mathbb{E}(Tr(PB_{T_{n}^{-1}\omega}\tilde{P}_{n}A_{T_{n}^{-1}\omega}P))$$
(9)

We have used Fubini's theorem to interchange expectation and sum, allowed because of (22)

$$= \sum_{n} \mathbb{E}(Tr(PB_{\omega}\tilde{P}_{n}A_{\omega}P))$$
(10)

since T_n^{-1} is probability preserving

$$= \mathbb{E}(\sum_{n} Tr(PB_{\omega}\tilde{P}_{n}A_{\omega}P))$$
(11)

 \tilde{P}_n is a partition of unity.

$$= \mathbb{E}(Tr(B_{\omega}A_{\omega}P)) = \mathbb{E}(Tr(PB_{\omega}A_{\omega}P)) \qquad (12)$$

Corollary 2. 1. If A_{ω} , B_{ω} , C_{ω} are covariant families of bounded operators satisfying the condition (1) then:

$$\mathbb{E}(Tr(PA_{\omega}B_{\omega}C_{\omega}P)) = \mathbb{E}(Tr(PC_{\omega}A_{\omega}B_{\omega}P))$$
(13)

2. If A_{ω}, B_{ω} are covariant families of bounded, positive (i.e. ≥ 0) operators satisfying the conditions (1) then

$$\mathbb{E}(TrPA_{\omega}B_{\omega}P)) \ge 0 \tag{14}$$

Proof. The first assertion is clear as we can apply the proposition to the covariant families $A_{\omega}B_{\omega}$ and C_{ω} .

For the second claim we observe that $B_{\omega} = C_{\omega}C_{\omega}$ with a $C_{\omega} = \sqrt{B_{\omega}}$. C_{ω} as a function of the covariant family B_{ω} is covariant as well. Moreover, since A_{ω} is positive (and the Hilbert space is complex), A_{ω} is self-adjoint and so is C_{ω} .

By part (i) of the corollary we have:

$$\mathbb{E}(Tr(PA_{\omega}B_{\omega}P)) = \mathbb{E}(PA_{\omega}C_{\omega}C_{\omega}P))$$
(15)

$$= \mathbb{E}(Tr(PC_{\omega}A_{\omega}C_{\omega}P)) \tag{16}$$

$$\geq 0$$
 since A_{ω} is positive (17)

Hypotheses 2. Let H_{ω} be family of self-adjoint operators, which are bounded below, on a Hilbert space \mathcal{H} . Let $E_{H_{\omega}}(\cdot)$ be the (projection-valued) spectral measure of H_{ω} such that for any bounded borel set A, the operators $PE_{H_{\omega}}(A), E_{H_{\omega}}(A)P$ are trace class for a.e. ω and form a **covariant** family of operators.

For operators H_{ω} satisfying the above hypothesis, it is clear that for any finite x, the spectral measure $E_{H_{\omega}}((-\infty, x]) = E_{H_{\omega}}([c, x])$, with c finite and smaller than the infimum of the spectrum of H_{ω} . Therefore the hypothesis implies that for any finite x, the operators $PE_{H_{\omega}}((-\infty, x]), E_{H_{\omega}}((-\infty, x])P$ are trace class. Therefore we can now define the density of states for such operators.

Definition 3. Let H_{ω} be a family of self adjoint operators satisfying Hypothesis 2. Then the density of states of this family is defined to be the unique

 σ -finite measure n associated with the monotone right continuous function F,

$$F(x) = \mathbb{E}\left(Tr(PE_{H_{\omega}}((-\infty, x])P)\right),\$$

via $n((a,b]) = F(b) - F(a), \ a, b \in \mathbb{R}.$

Thus for any bounded borel set A, n(A) agrees with the right hand side of the above relation with A replacing $(-\infty, x]$.

In the above framework we define another measure that is used to define the Interband Light Absorption Coefficient (ILAC). To do this we need a pair H^{\pm}_{ω} of self-adjoint operators as in the Hypothesis 2 and consider the associated projection valued measures $E_{H^{\pm}_{\omega}}(\cdot)$. We then define the density of states of these operators by,

$$n_{\pm}(A) = \mathbb{E}\left(Tr(PE_{H^{\pm}}(A)P)\right). \tag{18}$$

Consider the semi algebra $\mathcal{I}\times\mathcal{I}$ of subsets of \mathbb{R}^2 where

$$\mathcal{I} = \mathbb{R} \cup \{(a, b] : a, b \in \mathbb{R}\} \cup \{(a, \infty) : a \in \mathbb{R}\} \cup \{(-\infty, a] : a \in \mathbb{R}\}.$$

We define the correlation measure ρ on $\mathcal{I} \times \mathcal{I}$ as

$$\rho(A \times B) = \mathbb{E}\left(Tr(PE_{H^+_{\omega}}(A)E_{H^-_{\omega}}(B)P)\right),\tag{19}$$

where ρ is set to be ∞ if either A or B is an unbounded element of \mathcal{I} .

This set function takes values in [0, 1] if P is trace class and in $[0, \infty]$, if $PE_{H^{\pm}_{\omega}((a,b])}$ are trace class only for bounded intervals (a, b], in view of Proposition 3. We set

$$\rho(A) = \sum_{i=1}^{\infty} \rho(A_i \times B_i), \text{ if } A = \sqcup_{i=1}^{\infty} A_i, A_i \in \mathcal{I}.$$

It is a simple exercise to see that this ρ is well defined on $\mathcal{I} \times \mathcal{I}$ and via standard measure theory extends as a σ -finite measure to the whole borel σ -algebra of \mathbb{R}^2 .

Using the Hypothesis 2, and Proposition 1 we see that the following is valid.

Proposition 3. Consider the operators H_{ω}^{\pm} satisfying Hypothesis 2 and let n_{\pm} and ρ be as in equation (19). Then for any $B, C \in \mathcal{I}$ bounded,

- 1. $\rho(B \times C) = \mathbb{E} \left(Tr(PE_{H_{e}^{-}}(C)E_{H_{e}^{+}}(B)E_{H_{e}^{-}}(C)P) \right)$
- 2. $\rho(B \times C) = \mathbb{E}\left(Tr(PE_{H_{\alpha}^+}(B)E_{H_{\alpha}^-}(C)E_{H_{\alpha}^+}(B)P)\right)$
- 3. The following inequalities are valid

$$\rho(B \times C) \le n_+(B), \ \rho(B \times C) \le n_-(C).$$

Proof: Since the subsets B, C are bounded the operators $PE_{H_{\omega}^{-}}(C), PE_{H_{\omega}^{+}}(B)$ are covariant trace class operators satisfying the inequality (1). Therefore the result follows by an application of Proposition 1 and Corollary 2.

We collect the arguments about ρ in a proposition.

Proposition 4. Consider a pair of covariant operators H^{\pm}_{ω} satisfying the Hypothesis 2 and consider the correlation measures ρ extended to the borel σ -algebra on \mathbb{R}^2 from that given by equation (19). Then the following are valid.

- 1. If P is trace class, then ρ is a probability measure on \mathbb{R}^2 , with support in the closure of $\bigcup_{\omega} \sigma(H^+_{\omega}) \times \sigma(H^-_{\omega})$.
- 2. If P is not trace class but, $PE_{H^{\pm}_{\omega}}((a, b])P$ is trace class, for bounded intervals (a, b], then ρ is a positive σ -finite measure on \mathbb{R}^2 , with support in the closure of $\cup_{\omega} \sigma(H^+_{\omega}) \times \sigma(H^-_{\omega})$.

Remark 5. Typically the first case occurs for operators on $\ell^2(\mathbb{Z}^d)$ and the second case occurs in $L^2(\mathbb{R}^d)$.

We take the transformation T on \mathbb{R}^2 given by

$$T\begin{pmatrix}\lambda_1\\\lambda_2\end{pmatrix} = \begin{pmatrix}\frac{\lambda_1 + \lambda_2}{\sqrt{2}}\\\frac{\lambda_1 - \lambda_2}{\sqrt{2}}\end{pmatrix}.$$

Using this T we define the Interband Light Absorption Coefficient (ILAC) A as the distribution function,

$$A(\lambda) - A(\lambda') = \nu\left(\frac{1}{\sqrt{2}}(\lambda',\lambda]\right), \text{ where } \nu(B) = \rho \circ T^{-1}(B \times \mathbb{R})$$
(20)

In the above equation the factor $\frac{1}{\sqrt{2}}$ is because of the normalisation we used for T, so that this definition of ILAC agrees with the standard one in the case of finite box operators. We also note that since the operators H^{\pm}_{ω} are assumed to be bounded below $A(-\infty) = 0$.

In the case when \mathbb{P} in Hypothesis 1 is ergodic with respect to the action of G on Ω , then, the spectra $\sigma(H_{\omega}^{\pm})$ of covariant families of operators H_{ω}^{\pm} are almost everywhere constant sets. In such a case we can talk about the infimum of spectra of H_{ω}^{\pm} without reference to ω . In this context we have the following theorem.

Theorem 2.1. Suppose H_{ω}^{\pm} are a pair of random families of self-adjoint operators satisfying Hypothesis 1. Assume further that \mathbb{P} is ergodic with respect to the action of G on Ω .

- 1. Let $E_{\pm} = \inf \sigma(H_{\omega}^{\pm})$. Then $A(E_{+} + E_{-} + a) A(E_{+} + E_{-} a) \le n_{\pm}((E_{\pm} 2a, E_{\pm} + 2a)), a > 0.$
- 2. Let $E'_{\pm} = \sup \sigma(H^{\pm}_{\omega})$. Then $A(E'_{+} + E'_{-} + a) A(E'_{+} + E'_{-} a) \le n_{\pm}((E'_{\pm} 2a, E'_{\pm} + 2a)), a > 0.$

Proof: We shall prove the first case, the other proof is similar (where one has to use the fact that $\lambda_1 \leq E'_+, \lambda_2 \leq E'_-$ respectively for the other case and work it out). Let E_+, E_- to be the infima of the spectra $\sigma(H^+_{\omega}), \sigma(H^-_{\omega})$ of $H^+_{\omega}, H^-_{\omega}$. We consider the closure of the Cartesian product $\Sigma = \sigma(H^+_{\omega}) \times$ $\sigma(H^-_{\omega})$ of the spectra of H^\pm_{ω} , which is the support of the measure ρ . Therefore if we denotes points of Σ by (λ_1, λ_2) , so that $\lambda_1 \geq E_+, \lambda_2 \geq E_-$, then the possible values of $\lambda_1 + \lambda_2$ have a lower bound $E_- + E_+$, so $\lambda_1 + \lambda_2 \in (E_- + E_+, E_- + E_+ + a)$ implies $\lambda_1 \in (E_+ - 2a, E_+ + 2a)$ and $\lambda_2 \in (E_- - 2a, E_- + 2a)$, (see Figure 2). This immediately implies the inclusions (the first inclusion is clear and the second one uses the above):

$$\{(\lambda_1, \lambda_2) : \lambda_2 \in (E_-, E_- + (a/2)) \text{ and } \lambda_1 \in (E_+, E_+ + (a/2))\} \\ \subset \{(\lambda_1, \lambda_2) : \lambda_1 + \lambda_2 \in (E_- + E_+ - a, E_- + E_+ + a)\} \\ \subset \{(\lambda_1, \lambda_2) : \lambda_2 \in (E_-, E_- + 2a) \text{ and } \lambda_1 \in (E_+, E_+ + 2a)\}.$$

This then would lead to the inequalities that

$$\begin{aligned} A(E_{+} + E_{-} + a) - A(E_{+} + E_{-}) \\ &= \rho \circ T^{-1}(\frac{1}{\sqrt{2}}(E_{+} + E_{-}, E_{+} + E_{-} + a] \times \mathbb{R}) \\ &= \rho \left(\{ (\lambda_{1}, \lambda_{2}) : E_{-} + E_{+} \leq \lambda_{1} + \lambda_{2} \leq E_{-} + E_{+} + a \} \right) \\ &\leq \rho \left((E_{-}, E_{-} + 2a) \times (E_{+}, E_{+} + 2a) \right) \\ &\leq \min\{\rho \left((E_{-}, E_{-} + 2a) \times \mathbb{R} \right), \rho \left((E_{+}, E_{+} + 2a) \times \mathbb{R} \right) \} \\ &\leq \min\{n_{-} \left((E_{-} - 2a, E_{-} + 2a) \right), n_{+} \left((E_{+} - 2a, E_{+} + 2a) \right) \}, \end{aligned}$$

where the last inequality comes from Proposition 3(3) and enlarging the intervals slightly, which only increases the bound since n_{\pm} are measures. \Box

Remark 6. If the density of states n_{\pm} have Lifshitz tails behaviour $n_{\pm}((E_{\pm} - a, E_{\pm} + a)) \approx e^{-Ca^{\alpha}}$ as a goes to zero, for an appropriate α , at E_{\pm} respectively, then we have

$$\limsup_{a>0} \frac{1}{h(2a)} n_{\pm} \left((E_{-} - 2a, E_{-} + 2a) \right) < \infty,$$

for $h(a) = e^{-Ca^{\alpha}}$ for some α , so, using the above inequalities,

$$\limsup_{a>0} \frac{1}{h(2a)} (A(E_{-} + E_{+} + a) - A(E_{-} + E_{+} - a))$$

$$\leq \limsup_{a>0} \frac{1}{h(2a)} n_{+} ((E_{+} - a, E_{+} + 2a)) < \infty.$$

In the case when the density of states n_{\pm} have Lifshitz tails behaviour at other internal band edges, the same behaviour is valid for ILAC under some conditions. Suppose the spectra of H^{\pm}_{ω} consist of bands $\bigcup_{i=1}^{N} [a_i^{\pm}, b_i^{\pm}]$. Then the product of the spectra is $\bigcup_{i=1,j}^{N} [a_i^{\pm}, b_i^{\pm}] \times [a_j^{-}, b_j^{-}]$. Let us denote $R_{ij} = [a_i^{+}, b_i^{+}] \times [a_j^{-}, b_j^{-}]$. Then, the measure ρ is supported on the set $\bigcup_{i=1,j}^{N} R_{ij}$.

We index the pairs (ij) by β and use R_{β} to denote a rectangle forming part of Σ henceforth. So we have $\Sigma = \bigcup_{\beta} R_{\beta}$.

The central point in the proof of Theorem 2.1 is that if (c, d) is a corner of the rectangle formed by the lowest bands of the spectra of H^{\pm}_{ω} , then the strip $\{(\lambda_1, \lambda_2) : c + d \leq \lambda_1 + \lambda_2 \leq c + d + a\}$ intersected with the support of ρ is a triangle of side length $\sqrt{2a}$, (see Figure 2), hence its ρ measure is smaller than that of the square with the corner (c, d) and side length 2a, as can be seen in the Figure 2. As we see in Figure 1, there may be some rectangles in the support of ρ , with this property. Those rectangles in Figure 1, where this is not true are marked by X and the solid lines are those lines $\lambda_1 + \lambda_2 = const$ for which this feature is valid and the dashed lines are those for which this is not true.

In the definition below the sets $R_{\beta} \subset \mathbb{R}^2$ and we denote the coordinates of \mathbb{R}^2 by (λ_1, λ_2) .

Definition 4. Let the support of ρ be $\Sigma = \bigcup_{\beta} R_{\beta}$, with $R_{\beta} = [a_i^+, b_i^+] \times [a_j^-, b_j^-]$, $\beta = (ij)$. Then we call a corner (c, d) of a rectangle R_{β} good, if the intersection of the line $\lambda_1 + \lambda_2 = c + d$ with Σ consists of finitely many points and all of them are corners of rectangles forming Σ . Given a corner (c, d) in Σ we shall denote by $K_{c,d}$ the set of corners that lie on the line $\lambda_1 + \lambda_2 = c + d$.

Theorem 2.2. Let spectra of H_{ω}^{\pm} be as in theorem 2.1 and let Σ be the support of the measure ρ given in equation 19. Let A, as given in equation (20) be the corresponding ILAC. If (c, d) is a good corner in Σ . Denote the elements of $K_{c,d}$ by $\{(c_{\gamma}, d_{\gamma})\}$. Then we have

$$A(c+d+a) - A(c+d-a) \le \sum_{(c_{\gamma},d_{\gamma})\in K_{c,d}} \min \left\{ n_{+}((c_{\gamma}-2a,c_{\gamma}+2a)), n_{-}((d_{\gamma}-2a,d_{\gamma}+2a)) \right\}.$$

Proof: Firstly we note that if we take a rectangle, R_{β} , then only the lower-left and the top-right corners are candidates of being *good* corners, since for the other two corners, the line $\lambda_1 + \lambda_2 = const$ that contains the said corner will pass through the rectangle and hence has infinitely many points. We will prove the theorem for a good corner (c, d) which is a lower left corner of a rectangle, the proof for the case of a top-right good corner is similar. In this case we see immediately that if (c, d) is a good corner in Σ , then the intersection of the strip $S_a((c, d)) = \{(\lambda_1, \lambda_2) : c + d \leq \lambda_1 + \lambda_2 \leq c + d + a\}$ with Σ is contained in finitely many rectangles R_{β} forming Σ . Further $S_a((c, d)) \cap R_{\beta}$ is contained in a square of side length 2a contained in R_{β} and having one corner common with a corner of R_{β} . Given a good corner (c, d) and the associated strip $S_a((c, d))$, let $(c_{\gamma}, d_{\gamma}) \in K_{c,d}$ denote the corner of rectangle R_{γ} that has nonempty intersection with it. (Note that this corner satisfies $c_{\gamma} + d_{\gamma} = c + d$).

Then whenever (c, d) is a good corner we have the inequality, with γ ranging over a finite set,

$$S_a\left((c,d)\right) \cap \Sigma \subset \bigcup_{(c_\gamma,d_\gamma) \in K_{c,d}} [c_\gamma, c_\gamma + 2a] \times [d_\gamma, d_\gamma + 2a].$$

$$(21)$$

This inequality implies immediately that:

$$A(c + d + a) - A(c + d - a)$$

$$\leq A(c + d + a) - A(c + d) = \rho(S_a((c, d)) \cap \Sigma)$$

$$\leq \sum_{(c_{\gamma}, d_{\gamma}) \in K_{c,d}} \rho([c_{\gamma}, c_{\gamma} + 2a) \times [d_{\gamma}, d_{\gamma} + 2a))$$

$$\leq \sum_{(c_{\gamma}, d_{\gamma}) \in K_{c,d}} \min\{n_+((c_{\gamma} - 2a, c_{\gamma} + 2a)), n_-((d_{\gamma} - 2a, d_{\gamma} + 2a))\}\} (22)$$

where in the last inequality we enlarged the sets using the fact that n_{\pm} are measures.

This shows that at ILAC has the same continuity property as the density of states at the band edges. \Box

In the theorem below we identify good corners for a simple case of spectra having two bands.

Theorem 2.3. Consider a pair of self adjoint operators H_{ω}^{\pm} as in Theorem 2.1. Suppose a.e. ω , the spectra of $H_{\omega}^+, H_{\omega}^-$ are given by $\cup_{i=1}^2 [a_i^+, b_i^+]$ and $\cup_{i=1}^2 [a_i^-, b_i^-]$, respectively, where a_i^{\pm}, b_j^{\pm} are listed in the increasing order. Then the corners

$$\{(a_1^+, a_1^-), (b_1^+, b_1^-), (a_2^+, a_2^-), (b_2^+, b_2^-)\}$$

are good whenever a_i^{\pm}, b_i^{\pm} satisfy,

$$a_1^+ + a_1^- < b_1^+ + b_1^- < max(a_2^+ + a_1^-, a_1^+ + a_2^-) < max(b_2^+ + b_1^-, b_1^+ + b_2^-) < a_2^+ + a_2^- < b_2^+ + b_2^-.$$

In the case $a_i^+ = a_i^-, b_i^+ = b_i^-, i = 1, 2$, even the corners

$$\{(a_2^+, a_1^-), (a_1^+, a_2^-), (b_2^+, b_1^-), (b_1^+, b_2^-)\}$$

are good.

Proof: This is direct verification to see that the diagonal lines $\lambda_1 + \lambda_2 = const$ passing through the respective corners do not intersect any other rectangle. In the latter case when the spectra are the same, we have $a_2^+ + a_1^- = a_1^+ + a_2^-$ and $b_2^+ + b_1^- = b_1^+ + b_2^-$, hence the stated result.

Remark 7. In the symmetric case $a_i^{\pm} = a_i, b_j^{\pm} = b_j$, however, the rectangles $R_{\beta}, R_{\gamma} \subset S_{\beta}$ if $\beta = (ij), \gamma = (ji)$. In this case the above assumption still ensures that the lower-left and top-right corners of the rectangles are good.



Figure 1: Products of Spectra



Figure 2: A corner of a rectangle

3 Discrete Models:

Consider $\ell^2(\mathbb{Z}^d)$ and the discrete Laplacian $(\Delta u)(n) = \sum_{|n-i|=1} u(i)$. Consider real valued i.i.d random variables $\{q(n)\}$ with common distribution μ . Let V_{ω} denote the operator of multiplication by the sequence $q_{\omega}(n)$. Consider the operators

$$H^{\pm}_{\omega} = \Delta \pm q_{\omega}.$$

Taking $G = L = \mathbb{Z}^d$, it is known that operators $E_{H_{\omega}}(A)$ are covariant. The projection P is taken to be the projection $|\delta_0\rangle\langle\delta_0|$ onto the subspace generated by the vector δ_0 , which is an element of the standard basis for $\ell^2(\mathbb{Z}^d)$.

Then the density of states in these models are given by

$$n_{\pm}((a,b)) = \mathbb{E}\left(Tr(PE_{H_{\omega}^{\pm}}((a,b))\right) = \mathbb{E}\left(\langle \delta_{0}, E_{H_{\omega}^{\pm}}((a,b))\delta_{0}\rangle\right)$$

and the correlation measure ρ is given by

$$\rho((a,b) \times (c,d)) = \mathbb{E}\left(\langle \delta_0, E_{H^+_{a}}((a,b)) E_{H^-_{a}}((c,d)) \delta_0 \rangle\right)$$

and is a probability measure as per Proposition 4 (1), since P is trace class in this case.

In this model the density of states of H^{\pm}_{ω} are shown to have Lifshitz tails behaviour at the bottom of the spectra [21], under the condition that μ satisfies $\mu((a, a + \epsilon)) \geq C\epsilon^N$, where a is the infimum of the support of μ .

In the case when the support of μ has two closed intervals $[a_1, b_1] \cup [a_2, b_2], (a_i, b_i \text{ arranged in an increasing order so that <math>a_{i+1} > a_i$ for all i) and such that $b_i + 2d < a_{i+1} - 2d$, Simon [22] proves the Lifshitz tails behaviour at the internal band edges, if μ satisfies $\mu((a_i, a_i + \epsilon)) \geq C\epsilon^N$ and $\mu((b_i - \epsilon, b_i)) \geq C\epsilon^N$ for all i. When $[a_1 - 2d, b_1 + 2d]$ and $[a_2 - 2d, b_2 + 2d]$ are disjoint, Lifshitz tails behaviour at the band edges is also shown for the associated density of states n. That is at any of the band edges one has $n((E - \delta, E + \delta)) = O(e^{-C\delta^{-\frac{d}{2}}})$ as $\delta \to 0$.

An application of Theorems 2.1, 2.2 shows that the results are true for the ILAC A, namely

Theorem 3.1. Consider the Anderson models as above on $\ell^2(\mathbb{Z}^d)$. If $[a_1^{\pm}, b_1^{\pm}] \cup [a_2^{\pm}, b_2^{\pm}]$, are $\pm(supp(\mu))$. Then, for some C > 0,

• (External band edge case) For $E \in \{a_1^+ + a_1^-, b_2^+ + b_2^-\}$, one has

$$A(E+\delta) - A(E-\delta) = o(e^{-C\delta^{-\frac{d}{2}}}), \text{ as } \delta \to 0.$$

• (Internal band edge case) If the gap between the intervals $[a_1^{\pm} - 2d, b_1^{\pm} + 2d]$ and $[a_2^{\pm} - 2d, b_2^{\pm} + 2d]$ is large enough, then

$$A(E+\delta) - A(E-\delta) = o(e^{-C\delta^{-\frac{d}{2}}}), \text{ as } \delta \to 0,$$

for any $E \in \{b_1^+ + a_1^-, a_2^+ + a_1^-, b_2^+ + a_1^-, a_1^+ + a_2^-, b_1^+ + b_2^-, a_2^+ + a_2^-\}.$

Remark 8. We gave a simple example of a discrete model, however there are many more, those with periodic backgrounds [11], those which are unbounded [16] and so on. We refer to the review [7] for the numerous cases where the Lifshitz tails for the density of states is proved and for which our theorem applies to the ILAC.

4 Continuous Models:

Let us start by stating a theorem which is essentially a very weak version of the uncertainty principle.

We take $H_0 = -\Delta = -\sum_{i=1}^d D_j^2$, $D_j = i\frac{\partial}{\partial x_j}$, is self adjoint on its natural domain in $L^2(\mathbb{R}^d)$ and its spectrum is $[0,\infty)$.

We start with a couple of lemmas. We recall the definition of the trace ideal \mathscr{I}_p to be those bounded operators A with the property $|A|^p$ is trace class. Recall that elements of \mathscr{I}_1 are called Trace class operators.

Lemma 9. Consider $L^2(\mathbb{R}^d)$ and the operator $M = |-i\nabla|$. Then the operator $(|x|+i)^{-1}(M+i)^{-1} \in \mathscr{I}_{d+1}$.

Proof: Since the function $f(x) = (|x|+i)^{-1}$ is $L^{d+1}(\mathbb{R}^d)$ and the operator in question is just $f(x)f(-i\nabla)$, the result follows by an application of Theorem 4.1 in [23], which gives an estimate

$$||f(x)g(-i\nabla)||_{\mathscr{I}_p} \le 2\pi^{-\frac{d}{p}}||f||_p ||g||_p.$$

Let V be an operator of multiplication by a function V(x) on $L^2(\mathbb{R}^d)$ on its natural domain and such that V is bounded with respect to H_0 having relative bound smaller than 1. This means the operator $V(H+i)^{-1}$ is bounded. Then $H = H_0 + V$ is also self adjoint (Kato-Rellich theorem) on the domain of H_0 and its spectrum is also bounded below. Writing $(H_0 + i)(H + i)^{-1} =$ $I - V(H+i)^{-1}$, we see that $(H_0+i)(H+i)^{-1}$ is also bounded. Let P denote the operator of multiplication by the indicator function χ_{Λ} of a bounded region $\Lambda \subset \mathbb{R}^d$ on $L^2(\mathbb{R}^d)$. Let $E_H(A)$ denote the spectral measure of a bounded borel set A, with respect to the (projection valued) spectral measure of H. Then,

Theorem 4.1. Consider $L^2(\mathbb{R}^d)$ and the operator $H_0 = -\Delta$. Let V be an operator of multiplication by a function V(x), such that V is relatively bounded w.r.t. H_0 with relative bound c < 1 and consider $H = H_0 + V$. Suppose either

- 1. $d \leq 3$, then $PE_H(A)$ and E_HP are Hilbert-Schmidt, so $PE_H(A)P$ is trace class for any bounded borel set A.
- 2. Suppose $d \ge 1$ and suppose V is bounded or $\frac{\partial}{\partial x_j}V, j = 1, \ldots, d$ are relatively bounded w.r.t H_0 . Then $PE_H(A)$ and $E_H(A)P$ are trace class.

Proof: (1) Writing $PE_H(A) = P(|x|+i)^d(|x|+i)^{-d}(H_0+i)^{-1}(H_0+i)(H+i)^{-1}(H+i)E_h(A)$, we see that since all the factors are bounded, it is enough to show that $(|x|+i)^{-d}(H_0+i)^{-1}$ is Hilbert-Schmidt. The operator $(H_0+i)^{-1}$ is multiplication by $(|\xi|^2+i)^{-1}$ after taking Fourier transforms and hence is in $L^2(\mathbb{R}^d), d \leq 3$. Therefore an application of Lemma 9, shows that the product is Hilbert-Schmidt.

(2) We will prove that $PE_H(A) \in \mathscr{I}_1$, the proof for $E_H(A)P$ is similar. By taking a compactly supported smooth function ϕ which is value 1 on the closure of A, we have $\phi(H)E_H(A) = E_H(A)$. We will therefore show that $P\phi(H)$ is trace class for any compactly supported smooth function ϕ . We also note that the function $H\phi(H)$ is again a function of the same type as ϕ .

Further since P is multiplication by compactly supported function of x, $P(x^2 + i)^d$ is bounded. Therefore we will show that $(x^2 + i)^{-d}\phi(H) \in \mathscr{I}_1$.

We prove this by induction. Before we start, we note that if $M \in \mathscr{I}_p$ and N is a bounded operator then $MN \in \mathscr{I}_p$.

First consider $(x^2 + i)^{-1}\phi(H)$. We write this product as $(x^2 + i)^{-1}(H + i)^{-1}(H + i)\phi(H)$ and consider (recalling $M = |-i\nabla|$),

$$(x^{2}+i)^{-1}(H+i)^{-1} = (x^{2}+i)^{-1}(M+i)^{-1}(M+i)(H_{0}+i)^{-1}(H_{0}+i)(H+i)^{-1}.$$
(23)

The product of the first two factors is in \mathscr{I}_{d+1} (since $(|\xi| + i) - d - 1$ is integrable), by Lemma 9, the next two factors form a bounded operator (which can be seen by taking Fourier transforms). The final two factors form a bounded operator as argued before the lemma. Therefore the entire product is in \mathscr{I}_{d+1} . Since $(H + i)\phi(H)$ is bounded also, we get that $(x^2 + i)^{-1}\phi(H) \in \mathscr{I}_{d+1}$.

Now assume that $(x^2+i)^{-n}\phi(H) \in \mathscr{I}_{\frac{d+1}{n}}$, and show that $(x^2+i)^{-n-1}\phi(H) \in \mathscr{I}_{\frac{d+1}{n+1}}$. We write, $\psi(H) = (H+i)\phi(H)$, then

$$(x^{2}+i)^{-d-1}\phi(H)$$
(24)
= $(x^{2}+i)^{-d-1}(H+i)^{-1}(H+i)\phi(H)$
 $(x^{2}+i)^{-1}[(x^{2}+i)^{-n},(H+i)^{-1}]\psi(H) + (x^{2}+i)^{-1}(H+i)^{-1}(x^{2}+i)^{-n}\psi(H).$

Using Theorem 2.8 (2.5b) in [23], (which says $M \in \mathscr{I}_q, N \in \mathscr{I}_r \implies MN \in \mathscr{I}_p$ with $\frac{1}{p} = \frac{1}{q} + \frac{1}{r}$), so using the induction hypothesis and the already proved fact that $(x^2 + i)^{-1}(H + i)^{-1} \in \mathscr{I}_{d+1}$, the last term is seen to be in $\mathscr{I}_{\frac{d+1}{n+1}}$.

So we concentrate on the first term.

$$(x^{2}+i)^{-1}[(x^{2}+i)^{-n},(H+i)^{-1}]\psi(H)$$
(25)

$$= (x^{2}+i)^{-1}(H+i)^{-1}[H,(x^{2}+i)^{-n}](H+i)^{-1}\psi(H)$$

$$= (x^{2}+i)^{-1}(H+i)^{-1}[H_{0},(x^{2}+i)^{-n}](H+i)^{-1}\psi(H)$$

$$= (x^{2}+i)^{-1}(H+i)^{-1} \times$$

$$\left(-4ni\sum_{j=1}^{d}P_{j}x_{j}(x^{2}+i)^{-1} - 2d(x^{2}+i)^{-1} + 4d(n+1)x^{2}(x^{2}+i)^{-2}\right) \times (x^{2}+i)^{-n}\psi_{1}(H)$$
(26)

where we set $(H+i)^{-1}\psi(H) = \psi_1(H)$, where $P_j = -i\nabla_j$.

$$(x^{2}+i)^{-1}(H+i)^{-1}P_{j}$$

= $(x^{2}+i)^{-1}(H_{0}+1)^{-1/2}(H_{0}+1)^{\frac{1}{2}}(H+i)^{-1}(H_{0}+1)^{\frac{1}{2}}(H_{0}+1)^{-\frac{1}{2}}P_{j},$

and using Lemma 9, Lemma 10 below, we see that this expression is in \mathscr{I}_{d+1} . Induction hypothesis gives $(x^2 + i)^{-n}\psi_1(H) \in \mathscr{I}_{\frac{d+1}{n}}$. Therefore combining these two facts we see that $(x^2 + i)^{-n-1}\phi(H) \in \mathscr{I}_{\frac{d+1}{n+1}}$.

Lemma 10. Suppose either V is bounded or $(\frac{\partial}{\partial x_j}V)(H+i)^{-1}$, $j = 1, \ldots, d$ are bounded. Then $(H_0 + 1)^{\frac{1}{2}}(H+i)^{-1}P_j$ is a bounded operator for each $j = 1, \ldots d$.

Proof: Consider the case when $\frac{\partial}{\partial x_j}V(H+i)^{-1}$ is bounded for each j. Then writing the expression using commutators

$$(H_0 + 1)^{\frac{1}{2}}(H + i)^{-1}P_j = (H_0 + 1)^{\frac{1}{2}}P_j(H + i)^{-1} + (H_0 + 1)^{\frac{1}{2}}(H + i)^{-1}[P_j, H](H + i)^{-1} = (H_0 + 1)^{\frac{1}{2}}P_j(H + i)^{-1} + (H_0 + 1)^{\frac{1}{2}}(H + i)^{-1}(\frac{\partial}{\partial x_j}V)(H + i)^{-1}.$$

The boundedness of the first term was seen before since $(H_0 + 1)^{\frac{1}{2}}P_j(H_0 + 1)^{-1}$ and $(H_0 + 1)(H + 1)^{-1}$ are bounded. second term is bounded by the assumption on V and the boundedness of $(H_0 + 1)^{-\frac{1}{2}}$.

Now consider the case when V is bounded, then taking f, g in the domain of H_0 , we have

$$\langle f, (H_0+1)f \rangle = \langle g, (H_0+V+c+1)f \rangle - \langle g, (V+c)f \rangle,$$

where c is a positive constant such that H+c+1 is a positive operator (which is possible since H is bounded below). Since H+c+1 is positive it has a unique square root, so using the boundedness of V and the above inequality, we obtain, for some finite C,

$$\|(H_0+1)^{\frac{1}{2}}f\|^2 \le \|(H+c)^{\frac{1}{2}}f\|^2 + C\|f\| \le D\|(H+c)^{\frac{1}{2}}f\|^2.$$

Taking $f = (H_0 + 1)^{\frac{1}{2}}g$, ||g|| = 1, for a set of g coming from $C_0^{\infty}(\mathbb{R}^d)$, we see that

$$K \le ||(H+c)^{\frac{1}{2}}(H_0+1)^{-\frac{1}{2}}g||^2, \ K > 0,$$

K independent of g. This shows that $(H+c)^{\frac{1}{2}}(H_0+1)^{-\frac{1}{2}}$ has a bounded inverse and that its inverse $(H_0+1)^{-\frac{1}{2}}(H+c)^{\frac{1}{2}}$ and $(H+c)^{\frac{1}{2}}(H_0+1)^{-\frac{1}{2}}$ are both bounded (since M bounded implies M^* is also bounded). Therefore writing

$$(H_0 + 1)^{\frac{1}{2}}(H + i)^{-1}P_j$$

= $(H_0 + 1)^{\frac{1}{2}}(H + c)^{-\frac{1}{2}}(H + c)(H + i)^{-1}(H + c)^{-\frac{1}{2}}(H_0 + 1)^{\frac{1}{2}}(H_0 + 1)^{-\frac{1}{2}}P_j,$

we see that the left hand side is bounded.

Consider $L^2(\mathbb{R}^d)$, $H_0 = \Delta$, $q(n), n \in \mathbb{Z}^d$, i.i.d random variables with distribution μ having compact support. Let Λ denote the unit cube centred at $0 \in \mathbb{R}^d$ and $\Lambda(n)$ denote the unit cube centred at the point $n \in \mathbb{Z}^d$. Let $V_{\omega} = \sum_{n \in \mathbb{Z}^d} q^{\omega}(n) \chi_{\Lambda(n)}$, where χ_A is the operator of multiplication by the indicator function of A. Then taking

$$H^{\pm}_{\omega} = \Delta \pm V_{\omega},$$

we see that, since V_{ω} is bounded for each ω , the conditions of Theorem 4.1 are satisfied. Further taking $G = \mathbb{R}^d$ and $L = \mathbb{Z}^d$, $(U_x f)(y) = f(y-x)$, on $L^2(\mathbb{R}^d)$, $q_{T^m\omega}(n) = q_{\omega}(n+m)$, in Hypothesis 1, the spectral projections $E_{H_{\omega}^{\pm}}((a, b))$ are covariant families in the sense of Definition 1. The Theorem 4.1 shows that $\chi_{\Lambda(0)}E_{H_{\omega}^{\pm}}((a, b))\chi_{\Lambda(0)}$ is trace class whenever (a, b) is a bounded interval. Hence we can define the density of states and the ILAC as in equations (18) and (19), by taking P to be multiplication by $\chi_{\Lambda(0)}$. Therefore the Theorems 2.1, 2.2 are valid in this case.

Our theorem covers models where the random potential has the following forms.

- $V^{\omega}(x) = \sum_{n \in \mathbb{Z}^d} q^{\omega}(n)u(x-n)$, $\{q(n)\}$ i.i.d.random variables whose distribution has compact support and u a nice function with u(x-n) summable.
- An addition of a periodic background potential W to the random potential above.
- Addition of magnetic fields.

If in all these cases the density of states have Lifshitz tails behaviour at the band edges the same is acquired by the ILAC at an appropriate energy level.

Remark 11. Let us remark that in the above examples we can even replace the Laplacian $-\Delta$ with a real polynomial function Q of $-i\nabla$ and the results go through, if for some R > 0, the polynomial satisfies

$$c_1 \|\xi\|^{2n} \le Q(\xi) \le c_2 \|\xi\|^{2n}, \ |\xi| > R, c_1, c_2 > 0.$$

References

- M.S. Atoyan, E.M. Kazaryan, H.A. Sarkisyan: Interband light absorption in parabolic quantum dot in the presence of electrical and magnetic fields, Physica E: Low-dimensional Systems and Nanostructures Vol 31, 83-85 (2006).
- [2] R. Carmona, J. Lacroix: Spectral Theory of Random Schrödinger Operators, (Birkhäuser Verlag, Boston 1990)
- [3] H. Cycone, R. Froese, W. Kirsch and B. Simon: *Schrödinger Operators*, Texts and Monographs in Physic (Springer Verlag, 1985)
- [4] M. Demuth and M. Krishna : *Determining spectra in Quantum Theory*, Progress in Mathematical Physics Vol 44, (Birkhauser, Boston, 2005).
- [5] A.L. Efros: Density of states and interband absorption of light in strongly doped semiconductors, Semiconductors, Vol 16, 1209- (1982).
- [6] A. Figotin, L. Pastur: Spectra of Random and Almost-Periodic Operators (Springer Verlag, Berlin 1992)

- [7] W. Kirsch and B. Metzger: The Integrated Density of States for Random Schr dinger Operators, Spectral theory and mathematical physics: a Festschrift in honor of Barry Simon's 60th birthday, Proc. Sympos. Pure Math., Vol 76 649–696 (2007).
- [8] W. Kirsch and L.A.Pastur : The large-time asymptotics of some Wiener integrals and the interband light absorption coefficient in the deep fluctuation spectrum, Comm. Math. Phys., Vol 132, 365–382 (1990).
- [9] W. Kirsch and L.A.Pastur : The interband light absorption coefficient in the weak disorder regime: an asymptotically exactly solvable model, J. Phys. A, Vol 27, 2527–2543 (1994).
- [10] W. Kirsch, L.A.Pastur and H. Stork : Asymptotics of the interband light absorption coefficient near the band edge for an alloy-type model, J. Statist. Phys, Vol 92, 1173–1191 (1998)
- W. Kirsch, B. Simon: Lifshitz tails for periodic plus random potential, J. Statist. Phys, Vol 42, no 4-6, 799-808 (1986).
- [12] B. Khoruzenko, W. Kirsch and L.A.Pastur: The Interband Light Absorption Coefficient in the weak disorder regime: An asymptotically exactly solvable model, J. Phys. A: Math. Gen. Vol 27, 2527-2543 (1994).
- [13] Klopp, Frederic and Wolff, Thomas: Lifshitz tails for 2-dimensional random Schrödinger operators: Dedicated to the memory of Tom Wolff, J. Anal. Math. Vol 88, 63–147 (2002).
- [14] Klopp, F: Weak disorder localization and Lifshitz tails: continuous Hamiltonians, Ann. Henri Poincare, Vol 3, no. 4, 711–737 (2002).
- [15] Klopp, Frederic: Lifshitz tails for random perturbations of periodic Schrödinger operators, Spectral and inverse spectral theory (Goa, 2000), Proc. Indian Acad. Sci. Math. Sci. Vol 112, no. 1, 147–162 (2002).
- [16] Klopp, Frederic: Precise high energy asymptotics for the integrated density of states for an unbounded random jacobi matrix, Rev. Math. Phys, Vol 12, 575-620 (2000).
- [17] M. Krishna: Continuity of integrated density of states independent randomness, Proc. Ind. Acad. Sci., Vol 173, 401-410 (2007).

- [18] M. Krishna: Regularity of the Interband Light Absorption Coefficient, Preprint (2008).
- [19] K.R. Parthasarathy: Introduction to Probability and Measure, Texts and Readings in Mathematics, Vol 33, Hindustan Book Agency, New Delhi (2005).
- [20] M. Reed and B. Simon: *Methods of Modern Mathematical Physics*, *Functional Analysis*, Academic Press (1972).
- [21] B. Simon: Lifshitz tails for the Anderson Model, J. Stat. Phys, Vol 38, 65-76 (1985).
- [22] B. Simon: Internal Lifshitz tails, J. Stat. Phys, Vol 46, 911-918 (1987).
- [23] B. Simon: Trace Ideals and their applications, Second Edition, Mathematical Surveys and Monographs, Vol 120, American Mathematical Society, Providence RI (2005).
- [24] V. S. Sunder: Functional Analysis Spectral Theory, Texts and Readings in Mathematics, Vol 13, Hindustan Book Agency, New Delhi (1997).
- [25] Ivan Veselic: Existence and Regularity Properties of the Integrated Density of States of Random Schrödinger Operators, Lecture Notes in Mathematics No 1917, Springer Verlag (2008).