

## Higgs mechanism

Consider a gauge theory coupled to matter:

$$\mathcal{L} = \begin{cases} (P, \nabla): & \text{principal } G\text{-bundle w/conn} \\ \phi: & \text{section of } V_P \end{cases}$$

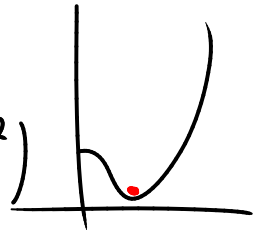
$$S = \frac{1}{4g^2} \int_X F \wedge *F + \frac{1}{2} \int_X \|D\phi\|^2 + \int_X H(\phi)$$

where  $H$  is some  $G$ -invariant function.

Suppose moreover that  $H$  has minima other than the origin.

For simplicity, say  $H$  has a single  $G$ -orbit worth of minima.

e.g.  $G = SU(2)$ ,  $V = \text{fundamental}$ ,  $H(\phi) = (\|\phi\|^4 - 2m^2\|\phi\|^2)$



In this case, the perturbative physics at low energies will be best understood as an expansion around the orbit  $\|\phi\| = m$ , not around  $\phi = 0$ .

In this pert. expansion, we can gauge-fix by taking  $P$  trivial and

choosing say  $\phi = \begin{pmatrix} 1 \\ 0 \end{pmatrix} f$ ,  $f \in \mathbb{R}$ . Then our theory is  $\approx$  to one with  $\mathcal{L} = \begin{cases} \text{conn } A \text{ on triv. } P \\ \phi = \begin{pmatrix} 1 \\ 0 \end{pmatrix} f \end{cases}$  and no gauge group  $\mathcal{G}$  anymore.

Write  $\phi_0 = m \begin{pmatrix} 1 \\ 0 \end{pmatrix}$  and  $\phi = \phi_0 + \delta\phi$ ;

$$\text{then } \|D\phi\|^2 = \|D(\phi_0 + \delta\phi)\|^2 = \|A\phi_0\|^2 + (\text{terms involving } \delta\phi)$$

Let's see concretely what  $\|A\phi_0\|^2$  looks like:

$$\text{if } A = A^1 \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} + A^2 \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix} + A^3 \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}, \quad A^i \in \Omega^1(X)$$

$$\text{get } \|A\phi_0\|^2 = m^2 (\|A^1\|^2 + \|A^2\|^2 + \|A^3\|^2)$$

i.e. all 3 components of  $A$  become massive. At energies  $E \ll gm$  they should be integrated out.  $\delta\phi$  is also massive.

If instead we take  $V = \text{adjoint}$ , then the picture is different: fix  $\phi = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} f$  say, then we have a residual gauge symmetry left over:

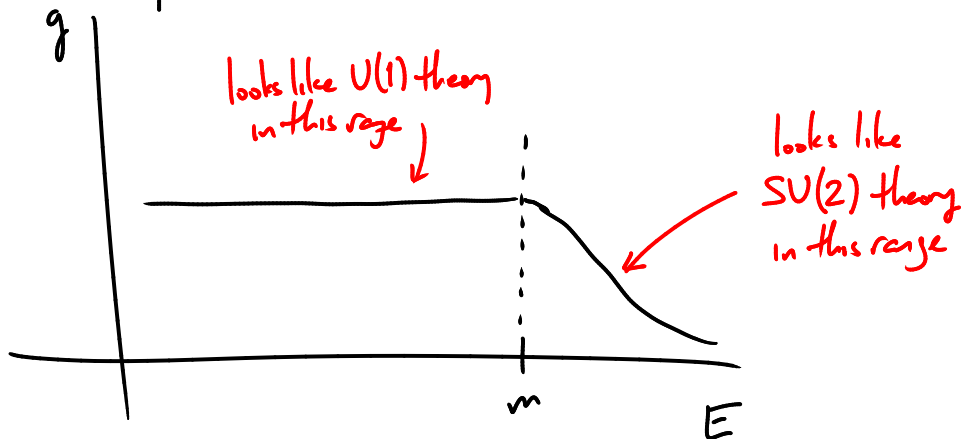
$$G_{\text{eff}} = \text{Map} \{ X, Z(\phi_0) \} \cong \text{Map} \{ X, U(1) \}$$

Have  $\| [A, \phi_0] \|^2 = m^2 (\|A^2\|^2 + \|A^3\|^2)$  so only 2 of the 3 components of  $A$  become massive. ("W bosons")

Set massive fields to 0  $\Rightarrow S_{\text{eff}} = \frac{1}{4g^2} \int F^{\frac{1}{2}} * F^{\frac{1}{2}}$

How reliable is this picture?

It would predict that the couplings run like:



It will be most reliable when  $m \gg \Lambda_s$  with  $\Lambda_s$  the strong-coupling scale of the nonabelian theory. (In that case the effective coupling is weak everywhere along the flow.)

For  $m \lesssim \Lambda_s$ , quantum corrections can change the picture even qualitatively.