

Riemann curvature

Levi-Civita connection ∇ in TM

because ∇ is orthogonal, (see exercises)

Denote its curvature by $R = F_{\nabla} \in \mathcal{E}(\Lambda^2 T^*M \otimes \mathfrak{o}(TM)) \subset \mathcal{E}(\Lambda^2 T^*M \otimes \text{End}(TM))$

$$R(X,Y)Z = \nabla_X \nabla_Y Z - \nabla_Y \nabla_X Z - \nabla_{[X,Y]} Z$$

Components: $R(\partial_i, \partial_j) \partial_k = R_{ijk}{}^l \partial_l$

Lower the last index to get R_m :

$$R_m(X,Y,Z,W) = \langle R(X,Y)Z, W \rangle = -\langle R(X,Y)W, Z \rangle = -R_m(X,Y,W,Z)$$

$$R_m(\partial_i, \partial_j, \partial_k, \partial_l) = R_{ijkl}$$

Symmetries: 1) $R_m(X,Y,Z,W) = -R_m(Y,X,Z,W) = -R_m(X,Y,W,Z)$

$$2) R_m(X,Y,Z,W) + R_m(Y,Z,X,W) + R_m(Z,X,Y,W) = 0$$

$$3) R_m(X,Y,Z,W) = R_m(Z,W,X,Y) \quad \text{since } R \text{ is } \mathfrak{o}(TM)\text{-valued}$$

Pf 1) $R_m(X,Y,Z,W) = \langle R(X,Y)Z, W \rangle = -\langle R(X,Y)W, Z \rangle = -R_m(X,Y,W,Z)$

2) Let $\sigma T(X,Y,Z) = T(X,Y,Z) + T(Y,Z,X) + T(Z,X,Y)$ for any T .

$$\text{Then } \sigma R(X,Y)Z = \sigma \nabla_X \nabla_Y Z - \sigma \nabla_Y \nabla_X Z - \sigma \nabla_{[X,Y]} Z$$

$$= \sigma \nabla_Z \nabla_X Y - \sigma \nabla_Z \nabla_Y X - \sigma \nabla_{[X,Y]} Z$$

$$= \sigma \nabla_Z (\nabla_X Y - \nabla_Y X) - \sigma \nabla_{[X,Y]} Z$$

$$= \sigma (\nabla_Z [X,Y] - \nabla_{[X,Y]} Z)$$

$$= \sigma [Z, [X,Y]]$$

$$= 0 \text{ by Jacobi}$$

since ∇ is torsion-free

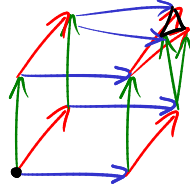
3) follows from 1), 2): $R(X,Y,Z,W) = -R(Z,X,Y,W) - R(Y,Z,X,W)$

$$= R(Z,X,W,Y) + R(Y,Z,W,X)$$

$$\begin{aligned}
&= -R(X,W,Z,Y) - R(W,Z,X,Y) - R(Z,W,Y,X) - R(W,Y,Z,X) \\
&= 2R(Z,W,X,Y) + R(X,W,Y,Z) + R(W,Y,X,Z) \\
&= 2R(Z,W,X,Y) - R(Y,X,W,Z) \\
&= 2R(Z,W,X,Y) - R(X,Y,Z,W)
\end{aligned}$$

so $R(X,Y,Z,W) = R(Z,W,X,Y)$

Rk 1) Geometric interpretation of 2):



2) Every R_m with these symmetries can be realized: if C_{ijkl} obeys 1), 2) then

$$g_{ik} = \delta_{ik} - \frac{1}{6} (C_{ijkl} + C_{iljk}) x^j x^l = \delta_{ik} - \frac{1}{3} C_{ijkl} x^j x^l \quad (*)$$

has $R_m(x=0)_{ijkl} = C_{ijkl}$. [We'll prove it below.]

3) If M is a Lie gp G with a bi-inv^t metric, then $\mathfrak{g} \cong T_e G$ and $R(X,Y)Z = \frac{1}{4} [Z, [X,Y]]$ where $[,]$ means the bracket in \mathfrak{g} (or the bracket of left-inv^t vector fields). So here the symmetry 2) becomes literally the Jacobi identity!

Def d_∇ acts on 2-forms by $d_\nabla G(X,Y,Z) = \nabla_X G(Y,Z) + \nabla_Y G(Z,X) + \nabla_Z G(X,Y) - G([X,Y],Z) - G([Y,Z],X) - G([Z,X],Y)$

(Differential) Bianchi identity

For any connection one has $d_\nabla F_\nabla = 0$.

(cf. abelian case, where locally $F_\nabla = dA$ so $dF_\nabla = d^2A = 0$)

Pf Recall if $\nabla = \nabla' + A$, ∇' flat, $A \in \mathcal{E}(T^* \otimes \text{End } E)$

$$F_\nabla = d_{\nabla'} A + A \wedge A \in \mathcal{E}(\text{End } E)$$

and $d_\nabla G = d_{\nabla'} G + A \wedge G - G \wedge A$ for $G \in \mathcal{E}(\wedge^2 T^* \otimes \text{End } E)$ [Exercise]

$$\begin{aligned}
d_\nabla F_\nabla &= d_{\nabla'}(d_{\nabla'} A + A \wedge A) + A \wedge (d_{\nabla'} A + A \wedge A) - (d_{\nabla'} A + A \wedge A) \wedge A \\
&= d_{\nabla'} A \wedge A - A \wedge d_{\nabla'} A + A \wedge d_{\nabla'} A - d_{\nabla'} A \wedge A + A \wedge A \wedge A - A \wedge A \wedge A \\
&= 0
\end{aligned}$$

In particular this applies to R . Moreover, using the torsion-free condition it can be rewritten purely in terms of the tensor ∇R :

$$\text{Prop} \quad (\nabla_x R)(y, z) + (\nabla_y R)(z, x) + (\nabla_z R)(x, y) = 0$$

$$\begin{aligned} \text{Pf } e \nabla_x (R(y, z)) &= e [(\nabla_x R)(y, z) + R(\nabla_x y, z) + R(y, \nabla_x z)] \\ &= e [(\nabla_x R)(y, z) + R(\nabla_x y, z) - R(\nabla_y x, z)] \\ &= e [(\nabla_x R)(y, z) + R([x, y], z)] \end{aligned}$$

$$\text{so } e [(\nabla_x R)(y, z)] = e [\nabla_x (R(y, z)) - R([x, y], z)] = 0$$

$$\text{Prop } \varphi: M \rightarrow \tilde{M} \text{ isometry: } \varphi^* \tilde{R} = R.$$

$$\text{Prop [Riemann]} \quad \text{Say } R=0.$$

Then any $p \in M$ has a nbhd isometric to an open subset of \mathbb{R}^n .

Pf Fix $p \in M$ and take an orthonormal frame at p . Fix a simply connected nbhd U of p . Since $R=0$, frame can be extended to a basis $\{e^i\}$ of sections of TM over U , with $\nabla e^i = 0$. In p^* , $\nabla_{e^i} e^j - \nabla_{e^j} e^i = [e^i, e^j] = 0$.

\implies can find coordinates x^i such that $e^i = \frac{\partial}{\partial x^i}$ (perhaps after shrinking U) [Frobenius]

In these coordinates $g_{ij} = \delta_{ij}$.

Computing in local coordinates

$$R_{ijk}{}^l \partial_l = R(\partial_i, \partial_j) \partial_k$$

$$= \nabla_{\partial_i} \nabla_{\partial_j} \partial_k - \nabla_{\partial_j} \nabla_{\partial_i} \partial_k$$

$$= \nabla_{\partial_i} (T_{jk}{}^l \partial_l) - \nabla_{\partial_j} (T_{ik}{}^l \partial_l)$$

$$= (\partial_i T_{jk}{}^l - \partial_j T_{ik}{}^l + T_{jk}{}^m T_{im}{}^l - T_{ik}{}^m T_{jm}{}^l) \partial_l$$

[an instance of our formula $F = dA + A \wedge A$]

In normal coordinates this simplifies to

$$R_{jkl}^i = (\partial_i T_{jk}^l - \partial_j T_{ik}^l) = \frac{1}{2} (\partial_i (g^{lm} (\partial_j g_{km} + \partial_k g_{jm} - \partial_m g_{jk}))) - (i \leftrightarrow j)$$

$$= \frac{1}{2} g^{lm} (\partial_i \partial_j g_{km} + \partial_i \partial_k g_{jm} - \partial_i \partial_m g_{jk}) - (i \leftrightarrow j)$$

ie, $R_{ijkl}(x=0) = \frac{1}{2} [\partial_i \partial_k g_{jl} - \partial_i \partial_l g_{jk} - \partial_j \partial_k g_{il} + \partial_j \partial_l g_{ik}]$ (From this, can prove the Taylor exp. (*) of g.)

How to think about R?

One approach: look for its "irreducible pieces."

Let V be an \langle, \rangle space and $\underline{Riem} \subset (V^*)^4$ consist of tensors with the symmetries of R.

$O(V) \curvearrowright \underline{Riem}$ and we can ask for its irreducible pieces.

Fact: $\underline{Riem} \simeq \underline{Ric} \oplus \underline{Weyl} \oplus \underline{S}$ for $n \geq 3$; all irreducible for $n > 4$
 $\frac{1}{2}n(n+1)-1$ $\frac{1}{12}n(n+1)(n+2)(n-3)$ 1 (Weyl = Weyl⁺ \oplus Weyl⁻ for $n=4$)

Ricci curvature: $Ric_{ij} = R_{kij}^k$ (symmetric)

Scalar curvature: $S = R_i^i$

Weyl curvature: $C = R - \frac{1}{n-2} (Ric - \frac{S}{n} g) \otimes g - \frac{S}{2n(n-1)} g \otimes g$

where $(g \otimes h)(v_1, v_2, v_3, v_4) = g(v_1, v_3)h(v_2, v_4) + g(v_2, v_4)h(v_1, v_3) - g(v_1, v_4)h(v_2, v_3) - g(v_2, v_3)h(v_1, v_4)$


$S^2 T^* \otimes S^2 T^* \rightarrow S^2(A^2 T^*)$

Interpretations:

1) In normal coordinates, let $h =$ Euclidean metric,

$dvol(g) = \left[1 - \frac{1}{6} Ric_{kl} x^k x^l + \dots \right] dvol(h)$

(pf easy beginning from $g_{ij} = \delta_{ij} - \frac{1}{3} R_{ijkl} x^k x^l$: just expand $\sqrt{\det g}$, using $\det(\mathbb{1} + \epsilon A) = 1 + \epsilon \text{tr} A$)

This tells us about the volume of wedges of a geodesic ball.  $\frac{vol}{vol_{Euc}} = \left[1 - \frac{\epsilon^2}{6} Ric(v, v) \right]$

2) Integrating over the ball: $\int_{B_\epsilon^{(0)}} x^k x^l dvol(h) = \begin{cases} 0 & [k \neq l] \\ \frac{1}{n} \int_0^\epsilon r^2 dvol(h) = \frac{1}{n} \int_0^\epsilon r^{n+1} vol(S^{n-1}) dr & \\ = \frac{\epsilon^{n+2}}{n(n+2)} vol(S^{n-1}) = \frac{\epsilon^{n+2}}{n(n+2)} vol(S^{n-1}) & [k=l] \\ = \frac{\epsilon^2}{n+2} vol(B_\epsilon^n) & \end{cases}$

$$\text{So, } \frac{\text{vol}(B_\epsilon(p))}{\text{vol}(B_\epsilon(o))_{\mathbb{R}^n}} = 1 - \frac{S(p)}{6(n+2)} \epsilon^2 + \dots$$

Special cases: if $n=1$: $R=0$ identically

$n=2$: $R_{12}{}^{21} = \frac{1}{2}S$, other components related by symmetries
