

Renormalization of Hamiltonian flows and critical invariant tori

(H. Koch)

- introduction
- formal definition of \mathcal{R}
- near-integrable Hamiltonians
- connection with commuting maps
- normal form: resonant Hamiltonians
- a non-trivial RG fixed point
- on critical invariant tori

→ DCDS 8, #3 (2002)

→ mp_arc 02-175

Collaborators on closely related problems

J.J. Abad, P. Wittwer

Other related work

- **perturbative renormalization**

(Eliasson, Feldman, Trubowitz, Gallavotti, Chierchia, Falcolini, Ecalle, Valet, Gentile, Mastropietro, ...)

- **renormalization à la QFT**

(Bricmont, Gawedzki, Kupiainen, Schenkel, ...)

- **commuting maps**

(Kadanoff, Shenker, MacKay, Morrison, ..., Stirnemann)

- **approx. renormalization of Hamiltonians**

(Escande, Doveil, Jausliu, Govin, Chandre, Benfatto, Celletti, ...)

- **renormalization of torus flows**

(J. Lopes Dias)

See references in mp_arc 02-175

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Consider mainly analytic H 's

$$H(q, p), \quad q \in \mathbb{T}^d, \quad p \in B^d$$

$$\dot{F} = \{F, H\} = \nabla_q F \cdot \nabla_p H - \nabla_q H \cdot \nabla_p F$$

$$\dot{q} = \nabla_p H$$

$$\dot{p} = -\nabla_q H$$

and closed orbits,

or invariant ω -tori "centered at $p=0$ " *

$$\Gamma : \mathbb{T}^d \times \{0\} \rightarrow \mathbb{T}^d \times B^d$$

$$\Phi_t \circ \Gamma = \Gamma \circ \Phi_t^0$$

Integrable example

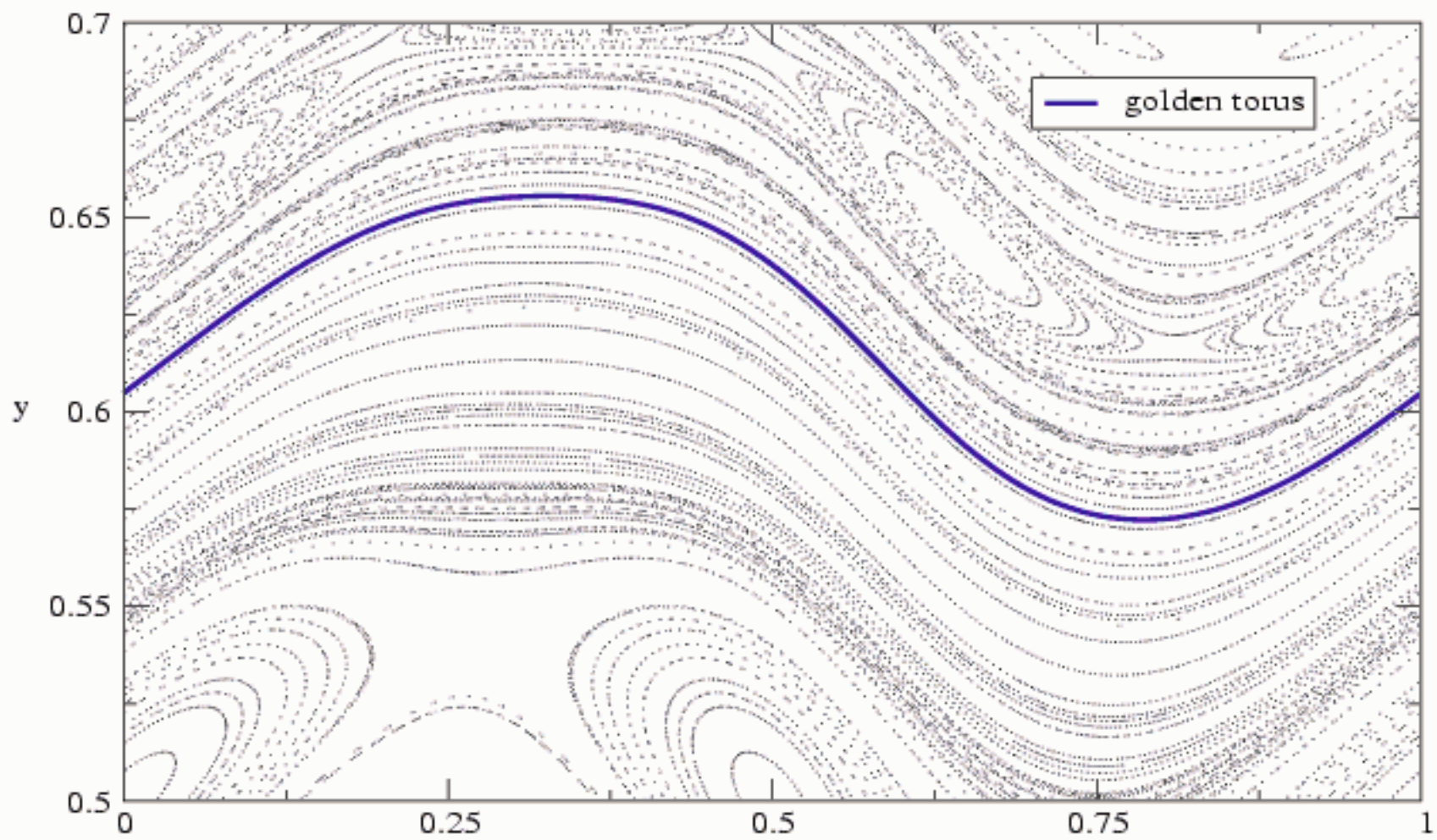
$$H(q, p) = \omega \cdot p + h(p) \quad Dh(0) = 0$$

$$\Phi_t^0(q, 0) = (q + t\omega, 0)$$

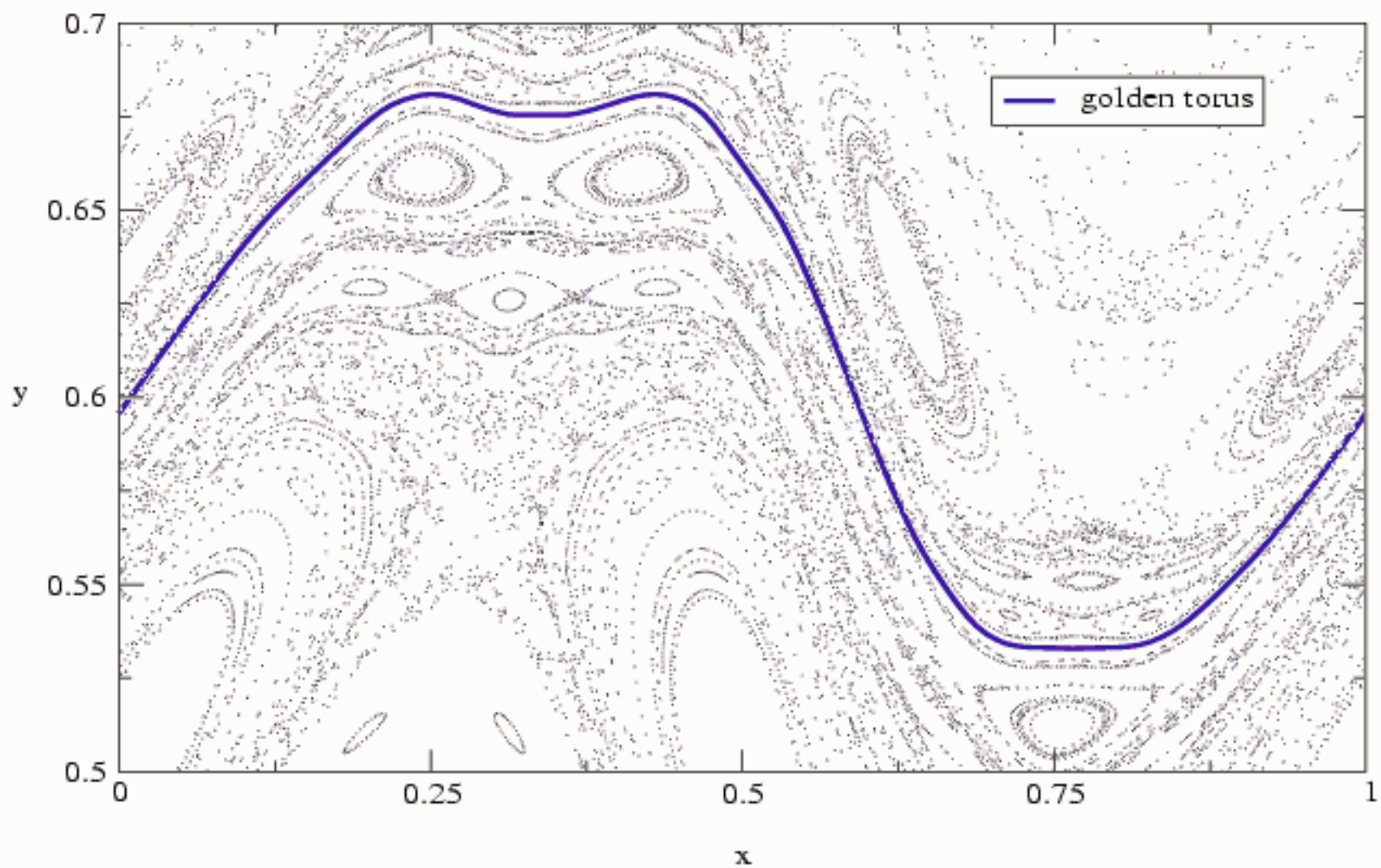
* $\int_{\gamma} p \cdot dq = 0$

standard map orbits (A. Haro)

$K=0.50$

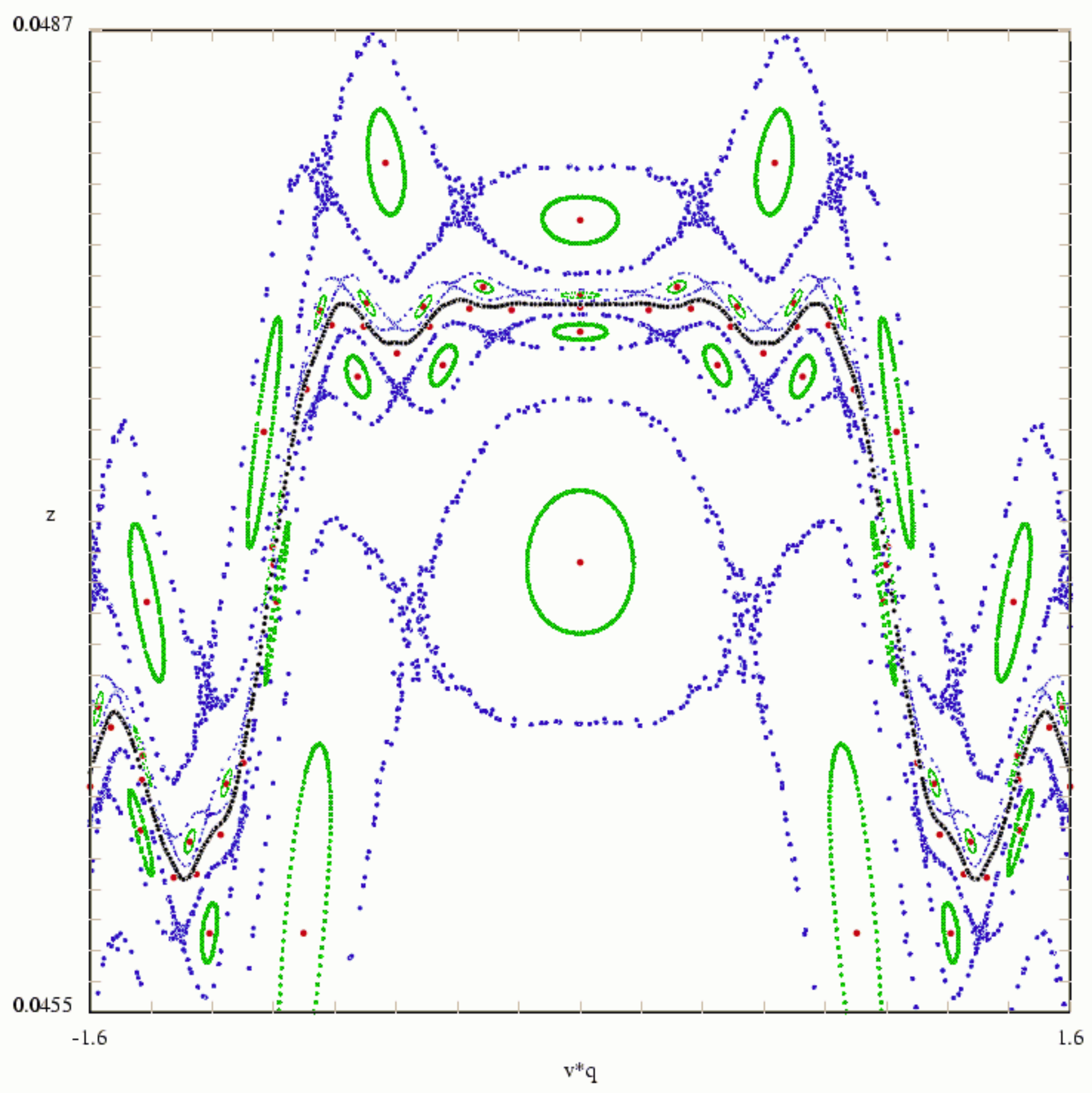


$K=0.93$



Periods $\frac{13}{8}$, $\frac{21}{13}$, $\frac{34}{21}$, ... and critical circle
for the return map of H_*
(J.J. Abad, H.K., P. Wittwer)

Figure 1



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Phenomena considered are invariant under

\mathcal{G} : transformations that do not change
(direction of) rotation vectors, stability, ...

Considering only

- scaling of the energy (or time)

$$H \mapsto \tau H - E \quad (\tau, E \text{ const.})$$

- scaling of the action variables

$$H \mapsto \mu^{-1} H(\cdot, \mu \cdot)$$

- canonical transformations homotopic to Id

$$H \mapsto H \circ U_\phi$$

$U_\phi : (q, p) \mapsto (q + Q, p + P)$ defined implicitly

$$Q(q, p) = \nabla_p \phi(q, p + P(q, p))$$

$$P(q, p) = -\nabla_q \phi(q, p + P(q, p))$$

(periodic in q)

If ϕ is "small" then

$$H \circ U_\phi = H + \{H, \phi\} + \text{small}^2$$

Other canonical transformations can change (direction of) rotation vectors. Consider

$$\mathcal{J}(q, p) = (Tq, (T^*)^{-1}p) \quad T \text{ linear, ...}$$

Correspondence:

$$\begin{array}{ccc} \text{orbit of } H & \text{with rotation vector } \omega & \\ \text{---} & \text{---} & \\ \text{---} & H \circ \mathcal{J} & \text{---} & T^{-1}\omega \end{array}$$

Consider only integer $d \times d$ matrices T with

- $\det(T) = \pm 1$
- one expanding eigenv.: $T\omega = \varphi\omega$, $\varphi > 1$
- $d-1$ contracting eigenvalues

Example: golden mean

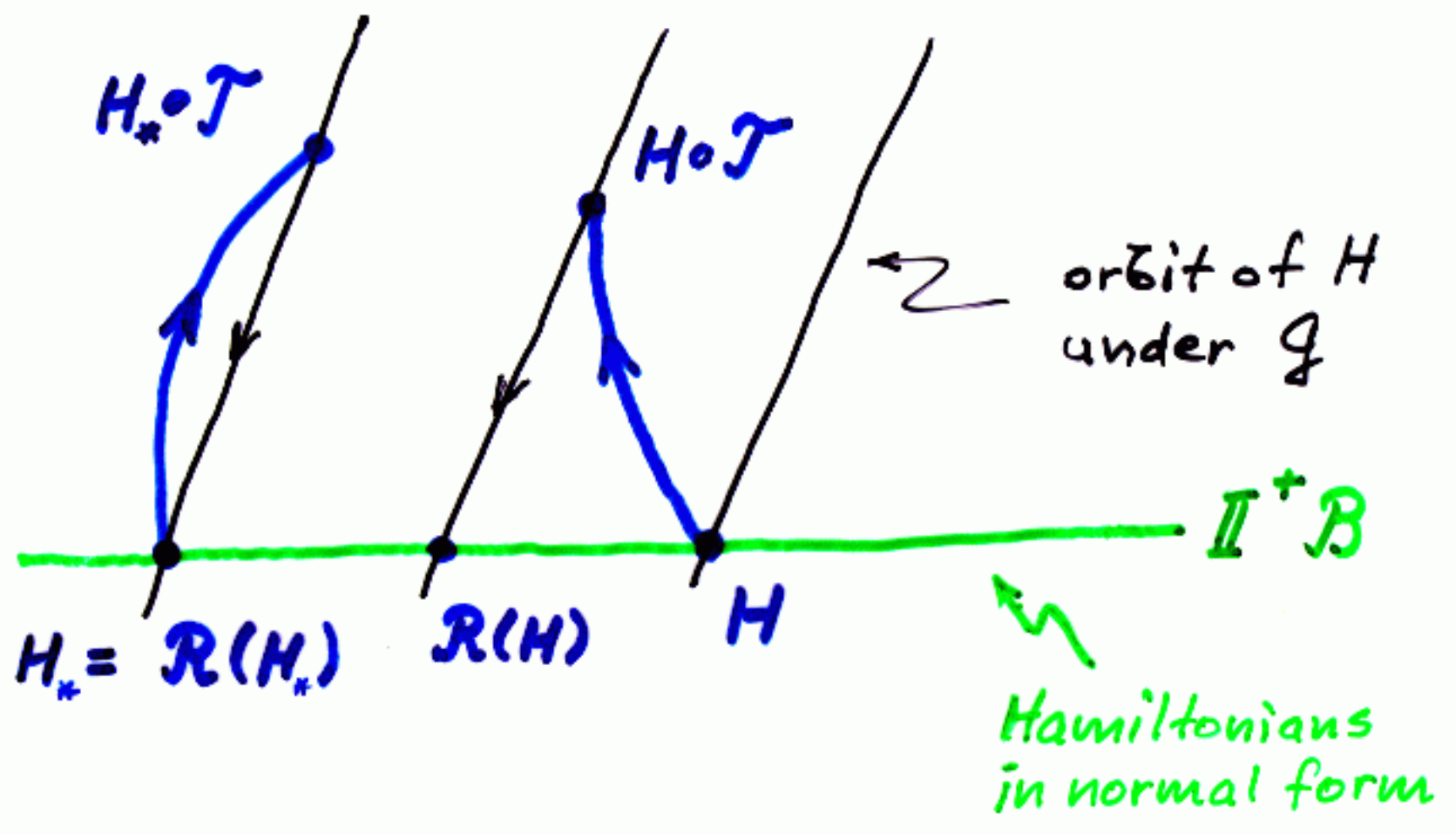
$$T = \begin{bmatrix} 0 & 1 \\ 1 & 1 \end{bmatrix}, \quad \omega = (1, \varphi), \quad \varphi = \frac{\sqrt{5}+1}{2}$$

$$\begin{bmatrix} 1 \\ 1 \end{bmatrix} \rightarrow \begin{bmatrix} 1 \\ 2 \end{bmatrix} \rightarrow \begin{bmatrix} 2 \\ 3 \end{bmatrix} \rightarrow \begin{bmatrix} 3 \\ 5 \end{bmatrix} \rightarrow \dots$$

$$\frac{1}{2}, \frac{2}{1}, \frac{3}{2}, \frac{5}{3}, \dots \rightarrow \varphi$$

$$\mathcal{R}(H) = H \circ \mathcal{I} \pmod{\mathcal{G}}$$

Choose an appropriate* normal form



\mathbb{I}^+ : projection onto resonant subspace

* appropriate:

- \mathcal{R} is a hyperbolic dynamical system on $\mathbb{I}^+ \mathcal{B}$
- the number of expanding directions is finite, and as small as possible

$$\mathcal{R}(H) = \underbrace{\frac{\mathcal{I}_H}{\mu} H \circ \mathcal{T}_\mu \circ \mathcal{U}_{\phi_0} \circ \mathcal{U}_{H'}}_{H'} - E_H$$

$$\mathcal{T}_\mu : (q, p) \mapsto (Tq, \mu(T^*)^{-1} p) \quad \mu = \dots$$

\mathcal{U}_{ϕ_0} : rough normalization
 same for all H 's in a small n'hood
 = Id for analysis of near-integrable H 's

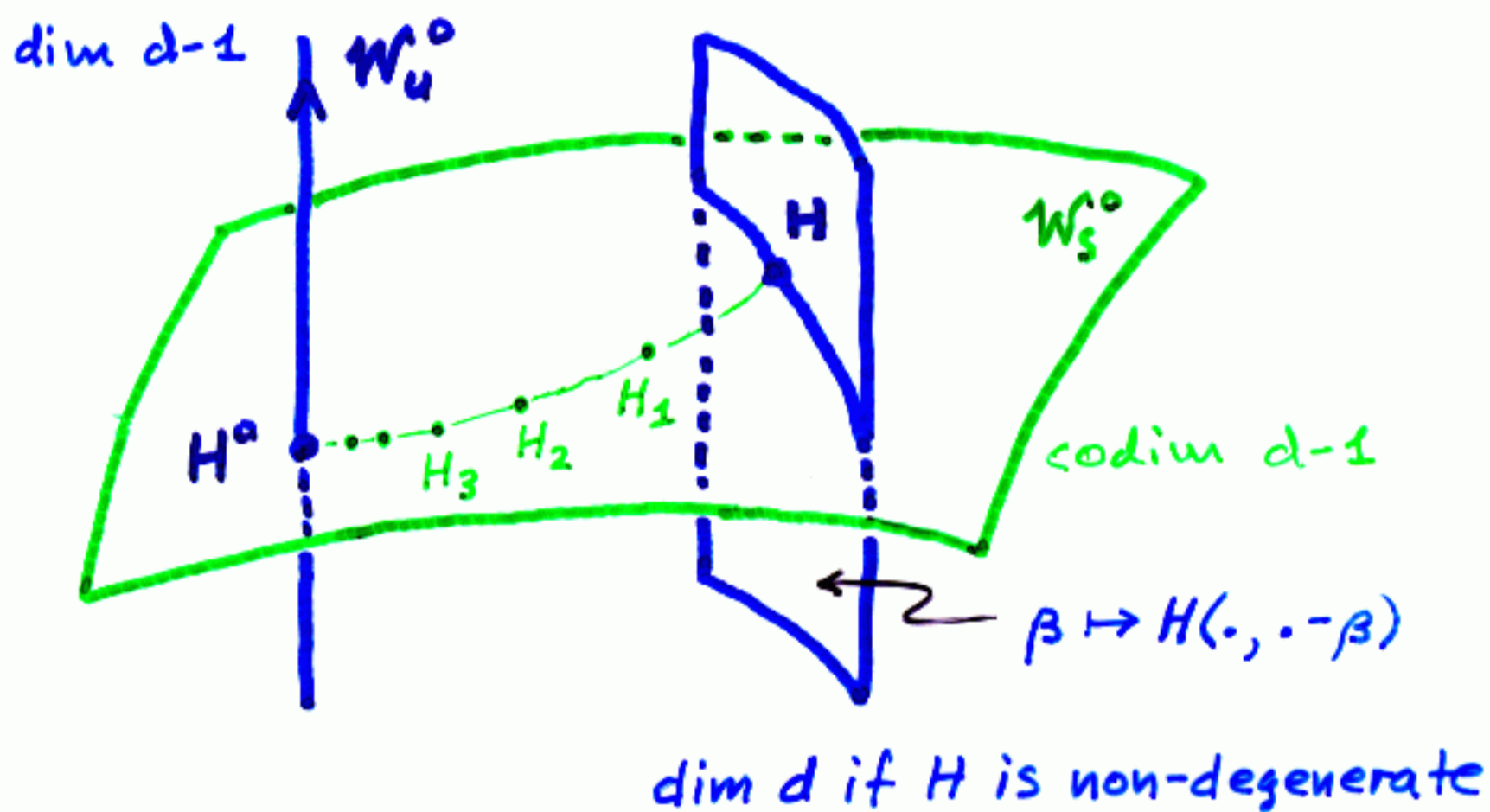
$\mathcal{U}_{H'}$: final normalization
 = $\mathcal{U}_{\phi_1} \circ \mathcal{U}_{\phi_2} \circ \dots$ ϕ_k small, non-resonant
 eliminates non-resonant modes

$$\mathbb{I}^-(H' \circ \mathcal{U}_{H'}) = 0$$

$\mathbb{I}^- = \mathbb{I} - \mathbb{I}^+$: projection onto non-resonant subspace

(precise definitions later)

Theorem: If ... then near $H^0(q,p) = w \cdot p$,
 \mathcal{R} is well defined and analytic and ...



Application 1: invariant ω -torus

$$\Gamma_H = V_0 \circ V_1 \circ V_2 \circ \dots$$

$$V_n = \mathcal{T}_\mu^n \circ \mathcal{U}_{H_n} \circ \mathcal{T}_\mu^{-n}$$

$$H_n = \mathcal{R}^n(H)$$

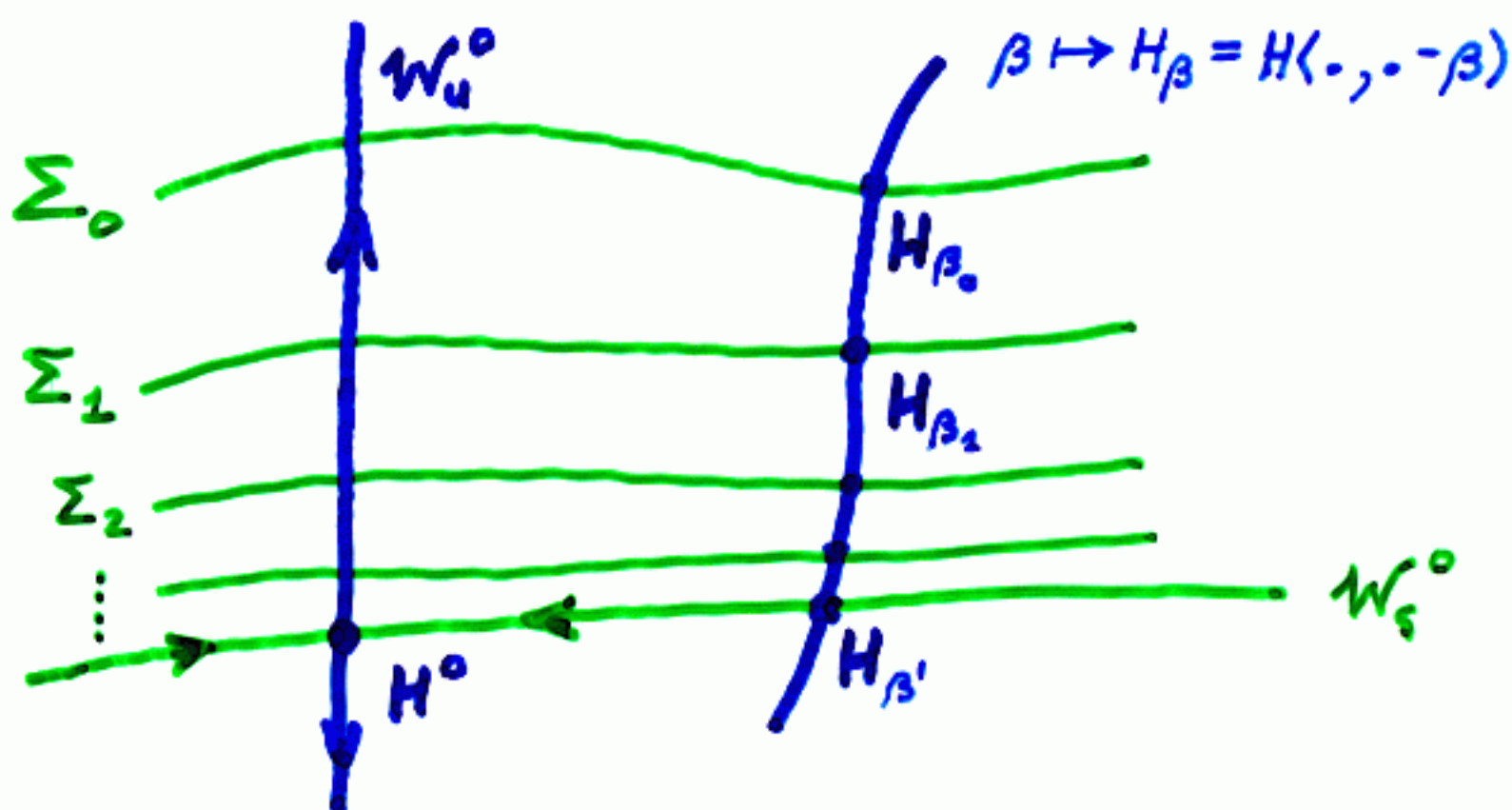
Application 2 (with J.J. Azaad)

Consider H 's even under $q \mapsto -q$.
 Fix w near ω with w_i/w_j rational.
 Define a manifold Σ such that H 's on

$$\Sigma_n = \mathcal{R}^{-n}(\Sigma \cap \text{Ball})$$

have a closed orbit γ_n with frequency
 vector proportional to $w_n = T^n w$

Theorem: If ... then

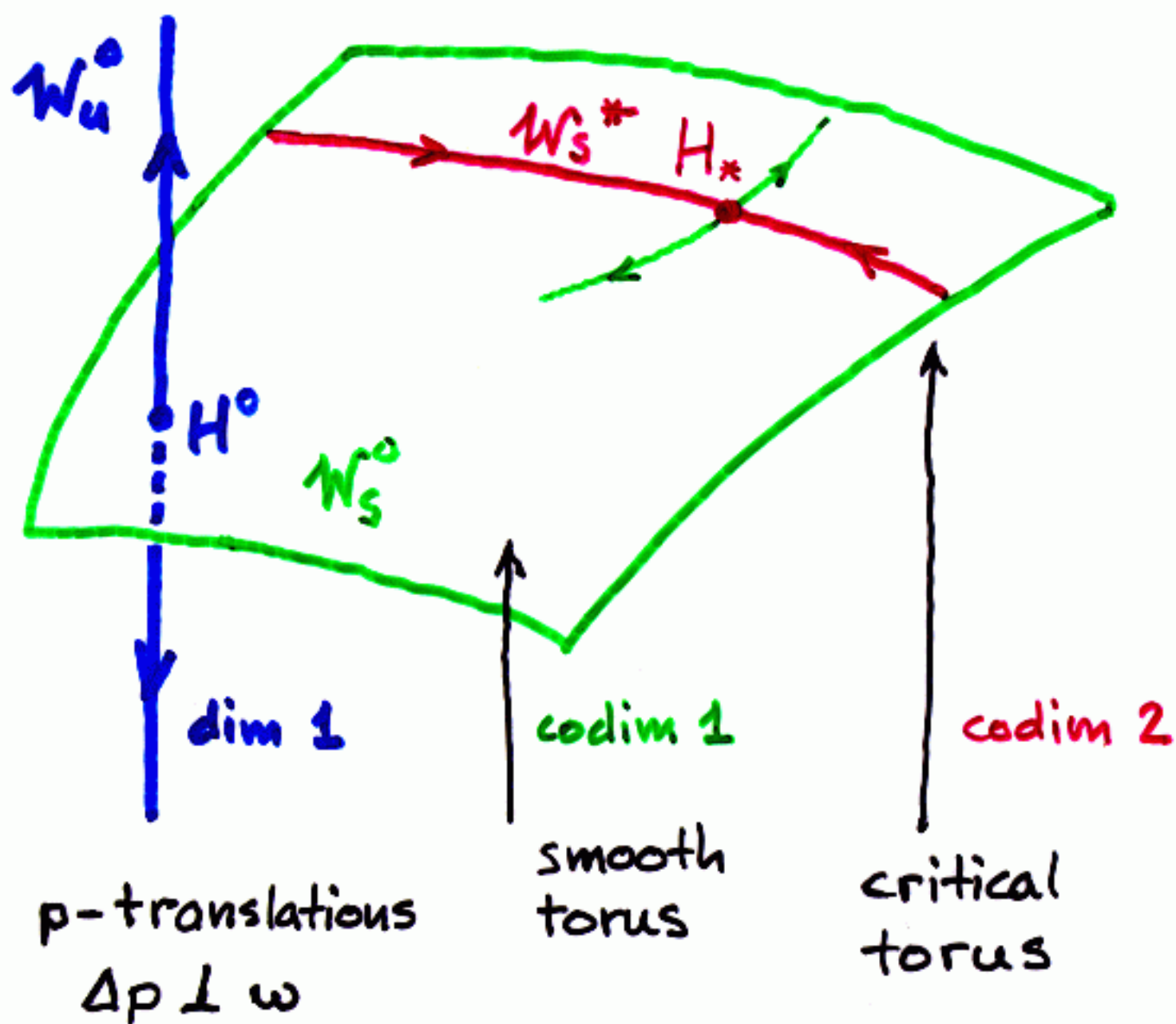


$$-\frac{\varepsilon}{n} \ln |\dot{\gamma}_n(0) - \Gamma'(0)| = \ln |\lambda_2| + O\left(\frac{\varepsilon}{n}\right)$$

↑
 2nd largest eigenvalue of $D\mathcal{R}(H^0)$

For $T = \begin{bmatrix} 0 & 1 \\ 1 & N \end{bmatrix}$ ($\vartheta = \frac{N + \sqrt{N^2 + 1}}{2}$)

expecting



Connection with commuting maps:

Let $T^* \omega' = \mathcal{J} \omega'$, $\omega' \cdot \omega = 1$. Assume


$$\omega' \cdot \nabla_p H = 1 \quad (\omega' \cdot q \text{ is "time"})$$

Consider only coord. changes leaving $\omega' \cdot q$ invar.

$$\mathcal{R}(H) = \frac{\mathcal{J}}{\mu} H \circ \Lambda \quad \Lambda = \mathcal{T}_\mu \circ \mathcal{U}_\phi \circ \mathcal{U}_H$$

Define

$$F_k = \mathbb{I}_{2\pi\omega'_k} \circ V(2\pi\delta_k) \quad k=1,2,\dots,d$$



 q -translations

Theorem: Assume H diff'ble, Λ diffeo, on $T^d \times \mathbb{R}^d$.

Then F_k 's commute with each other, and

leave $\{(q,p) \in H^{-1}(0) : \omega' \cdot q = 0\}$ invariant.

If $\tilde{H} = \mathcal{R}(H)$ and ... then

$$\tilde{F}_k = \Lambda^{-1} \circ F_1^{T_{1,k}} \circ F_2^{T_{2,k}} \circ \dots \circ F_d^{T_{d,k}} \circ \Lambda$$

From now on

$$T = \begin{bmatrix} 0 & 1 \\ 1 & 1 \end{bmatrix}, \quad \vartheta = \frac{\sqrt{5} + 1}{2}$$

$$\omega = (\vartheta^{-1}, 1), \quad \omega' = \text{const } \omega \quad \text{s.t.} \quad \omega' \cdot \omega = 1$$

$$\Omega = (1, -\vartheta^{-1}), \quad \Omega' = \text{const } \Omega \quad \text{s.t.} \quad \Omega' \cdot \Omega = 1$$

Domain D_ϑ and space \mathcal{B}_ϑ

$$D_\vartheta: \text{Im } \omega' \cdot q = 0, \quad |\text{Im } \Omega' \cdot q| < \vartheta_1, \quad |\Omega \cdot p| < \vartheta_2$$

$$H(q, p) = \omega \cdot p + \sum_{\nu, k} h_{\nu, k} \cos(\nu \cdot q) (\Omega \cdot p)^k$$

$$\|h\|_\vartheta = \sum_{\nu, k} |h_{\nu, k}| \cosh(\vartheta_2 \Omega \cdot \nu) \vartheta_2^k$$

Notice

- no decay of coeff's in directions where $\omega \cdot \nu \rightarrow \pm\infty$ with $\Omega \cdot \nu$ and k bounded
(H's can be rough in direction of flow)
- $h \mapsto h \circ \tilde{\gamma}_\mu$ is bounded from \mathcal{B}_ϑ to $\mathcal{B}_{\vartheta\vartheta}$
if $|\mu| \leq \vartheta^{-2}$

Solve $\mathbb{I}^-(H \circ U_H) = 0$ iteratively:

$$H = H_2 + f_2, \quad \mathbb{I}^- H_2 = 0, \quad f_2 \text{ small}$$

First step (ϕ_2 small)

$$H \circ U_{\phi_2} = H_2 + f_2 + \underbrace{\{H_2, \phi_2\}} + \text{small}^2$$

Solve $\mathbb{I}^-(f_2 + \{H_2, \phi_2\}) = 0 \quad (\mathbb{I}^+ \psi_2 = 0)$

$$\psi_2 = [\mathbb{I} + \mathbb{I}^- \mathcal{D}(H_2 - H_0)]^{-1} \mathbb{I}^- f_2 \quad \phi_2 = \frac{1}{\omega \cdot \nabla_q} \psi_2$$

$$\mathcal{D}(h_2) = \left[(\Omega' \cdot \nabla_p h_2) \frac{\Omega \cdot \nabla_q}{\omega \cdot \nabla_q} - (\Omega \cdot \nabla_q h_2) \frac{\Omega' \cdot \nabla_p}{\omega \cdot \nabla_q} \right] \mathbb{I}^-$$

Now

$$\tilde{H} = H \circ U_{\phi_2} = H_2 + f_2, \quad f_2 \text{ small}^2$$

$$\begin{array}{ccc} & \uparrow & \uparrow \\ & \mathbb{I}^+ \tilde{H} & \mathbb{I}^- \tilde{H} \end{array}$$

iterate ...

$$U_H = U_{\phi_1} \circ U_{\phi_2} \circ U_{\phi_3} \circ \dots$$

Require

- (1) $\mathcal{D}(h_n)$ bounded, op-norm < 1
- (2) functions in $I^+ \mathcal{B}_s$ analytic

Define

$$I^+ : |\omega \cdot \nu| \leq \epsilon |\Omega \cdot \nu| \text{ or } |\omega \cdot \nu| < \kappa k$$

$$(I^+ h)(q, p) = \sum_{(\nu, k) \in I^+} h_{\nu, k} \cos(\nu \cdot q) (\Omega \cdot p)^k$$

Then

- (1) on I^- , $\frac{\Omega \cdot \nu}{\omega \cdot \nu}$ and $\frac{k}{\omega \cdot \nu}$ are bounded

$$\|\mathcal{D}(h)\|_s \leq \frac{1}{\epsilon} \|\Omega \cdot \nabla_p h\|_s + \frac{1}{\kappa \Omega_2} \|\Omega \cdot \nabla_q h\|_s$$

- (2) coeff's of $I^+ h$ decrease exponentially:
in I^+ , if $\omega \cdot \nu \rightarrow \pm \infty$ then $|\Omega \cdot \nu| \rightarrow \infty$ or $k \rightarrow \infty$

Specifically, if $S' = (S_1 + \epsilon \epsilon, e^{\epsilon k} S_2)$ then

$$\cosh(S'_2 \Omega \cdot \nu) (S'_2)^k \geq \text{const } e^{\epsilon |\omega \cdot \nu|} \cosh(S_2 \Omega \cdot \nu) S_2^k$$

$$(\nu, k) \in I^+$$

$$\begin{aligned}
(\mathcal{S}H)(q, p) &= c_H H(q, p/c_H), \\
\mathcal{L}H &= \vartheta \mu_0^{-1} H \circ \mathcal{T}_{\mu_0} \circ U_{\phi_0}, \\
\mathcal{N}(H) &= H \circ \mathcal{U}_H,
\end{aligned}$$

where $c_H = 2h_{0,2}$ and $\mu_0 = \vartheta^{-3}$.

$$\mathcal{R} = \mathcal{N} \circ \mathcal{L} \circ \mathcal{S}$$

Theorem. *There exists an even real resonant Fourier-Taylor polynomial h_1 , an odd real Fourier-Taylor polynomial ϕ_0 , and a choice of the parameters σ, κ, ρ , such that \mathcal{R} defines a compact analytic map, from some open neighborhood B of $H_1 = H_0 + h_1$ in \mathcal{B}'_ρ , to the resonant subspace of \mathcal{B}'_ρ . This map \mathcal{R} has a unique fixed point H_* in B , which is real analytic, and non-trivial, in the sense that $c_{H_*} > 1$ ($c_{H_*} = 1.024332969 \dots$).*

For a (computer-assisted) proof,
consider an approx. Newton map \mathcal{M} ,

$$\mathcal{M}(h) = h + \mathcal{R}(H_2 + Mh) - (H_2 + Mh)$$

$$M \approx \mathbb{I} - D\mathcal{R}(H_2)$$

and show that \mathcal{M} contracts a ball $B(r) \dots$

Problem: estimating $D\mathcal{M}$ "as usual"
is computationally prohibitive

Cure: perturbation theory about H_2
(includes the use of U_{ϕ_0})

Affine approximation \mathcal{N}_2 of $\mathcal{N}: H \mapsto H \circ U_H$

$$\mathcal{N}_2(H_2 + f_2) = H_2 + \mathbb{I}^+(f_2 + \{H_2, \phi_2\}) \quad \phi_2 = \dots$$

defines an approximate RG

$$\mathcal{R}_2 = \mathcal{N}_2 \circ \mathcal{L} \circ \mathcal{I}$$

Expecting $\mathcal{N} - \mathcal{N}_2$ and $\mathcal{R} - \mathcal{R}_2$ of order r^2

Define

$$\tilde{f}(H_2, f) = f(H_2 + f) - f(H_2) - Df(H_2)f$$

$$\tilde{R}(H_2, f) = DN_2(H_2) \mathcal{L} \tilde{f}(H_2, f)$$

Then

$$\mathcal{M}(h) = (\mathcal{R}_2(H_2) - H_2) \quad (a)$$

$$+ [\mathbb{I} - (\mathbb{I} - D\mathcal{R}_2(H_2))M] h \quad (b)$$

$$+ \tilde{R}(H_2, Mh) + (\mathcal{R} - \mathcal{R}_2)(H_2 + Mh) \quad (c)$$

To get $\mathcal{M}: B(r) \rightarrow B(r)$, show

(a) $\mathcal{R}_2(H_2) - H_2$ of norm $\varepsilon \ll r$

(b) [...] has op-norm $K < 1$

(c) $\|\tilde{R}(\dots) + (\mathcal{R} - \mathcal{R}_2)(\dots)\|_f \leq Cr^2$ with ...

To get contraction

extend bound (c) to $B(2r)$

Cauchy \rightarrow derivatives of non-linear terms
bounded by Cr on $B(r)$

Chose

$$\xi \approx 0.85001, \quad \kappa \approx \xi/0.4$$

$$\rho \approx (0.85, 0.15), \quad \rho^* > \rho$$

$$r \approx 3 \times 10^{-12}$$

Estimated

$$\varepsilon = \|\mathcal{R}_2(H_2) - H_2\|_{\rho^*} \quad (< 10^{-14})$$

$$K = \|\mathbf{I} - (\mathbf{I} - D\mathcal{R}_2(H_2))M\|_{\rho} \quad (< 0.84)$$

$$K_n = \|\tilde{\mathcal{R}}(H_2, Mh)\|_{\rho^*} \quad h \in B(nr) \quad n = 1, 2$$

$$K'_n = \|(\mathcal{R} - \mathcal{R}_2)(H_2 + Mh)\|_{\rho^*} \quad h \in B(nr) \quad n = 1, 2$$

Verified

$$\varepsilon + Kr + K_2 + K'_2 < r, \quad K + (K_2 + K'_2)/r < 1$$

etc.

- order of operators in \mathcal{R} chosen carefully (e.g. contraction near beginning)
- many "higher order" terms
- variable "degrees"
- optimized first 5 Nash-Moser steps
- parallelized composition estimates
- no switching between FPU modes
- ...

Critical tori :

(applies not only to golden mean, but ...)

consider $H \approx H_*$

$$\frac{\nu}{\mu} H \circ \Lambda = H \quad \Lambda = \mathcal{T}_\mu \circ U_{\phi_0} \circ U_H$$

Then

$$\Lambda \circ \Phi_t = \Phi_{\nu t} \circ \Lambda \quad (\#)$$

If H has an invariant ω -torus Γ
then $(\#)$ and technical assumptions imply that

$$\Lambda \circ \Gamma \circ \mathcal{T}^{-1}$$

is also an invariant ω -torus for H .

Thus, if unique,

$$\Gamma = \Lambda \circ \Gamma \circ \mathcal{T}^{-1}$$

Note: If Λ has a fixed point x , then

$$\Lambda(\Phi_t(x)) = \Phi_{\nu t}(x) \rightarrow \text{invariant manifold}$$

Conversely,

Theorem: Assume $\Gamma: T^2 \times \{0\} \rightarrow D_\varepsilon$ is continuous, "admissible", and satisfies

$$\Gamma = \Lambda \circ \Gamma \circ \mathcal{J}^{-1} \quad (*)$$

- (a) If the derivative of Λ at $x = \Gamma(0)$ has exactly one non-contracting direction, and if $t \mapsto \Gamma(t\omega, 0)$ is C^2 , then Γ is an invariant ω -torus for H .
- (b) If in addition, $-\vartheta^{-1}$ is not an eigenvalue of $D\Lambda(x)$, then Γ is not C^2 .

Proof: Use that Γ maps the unstable manifold of \mathcal{J} at 0 to the unstable manifold of Λ at x .
....

Next goal: Solve (*) and check conditions

Equation (*) is equiv. to $F(\gamma) = \gamma$,

$$F(\gamma) = \Lambda \circ \Gamma \circ \mathcal{J}^{-1} - I$$

$$= \mathcal{J}_\mu \circ [u \circ (I + \gamma) + \delta] \circ \mathcal{J}^{-1}$$

$$\Gamma - I$$

$$u - I$$

$$u = u_{\phi_0} \circ u_{\psi'} = \mathcal{J}_\mu^{-1} \circ \Lambda$$

Crucial: $u = (f\Omega, \dots)$ and $\delta = (g\Omega, \dots)$

have zero component

in the expanding direction of \mathcal{J}_μ

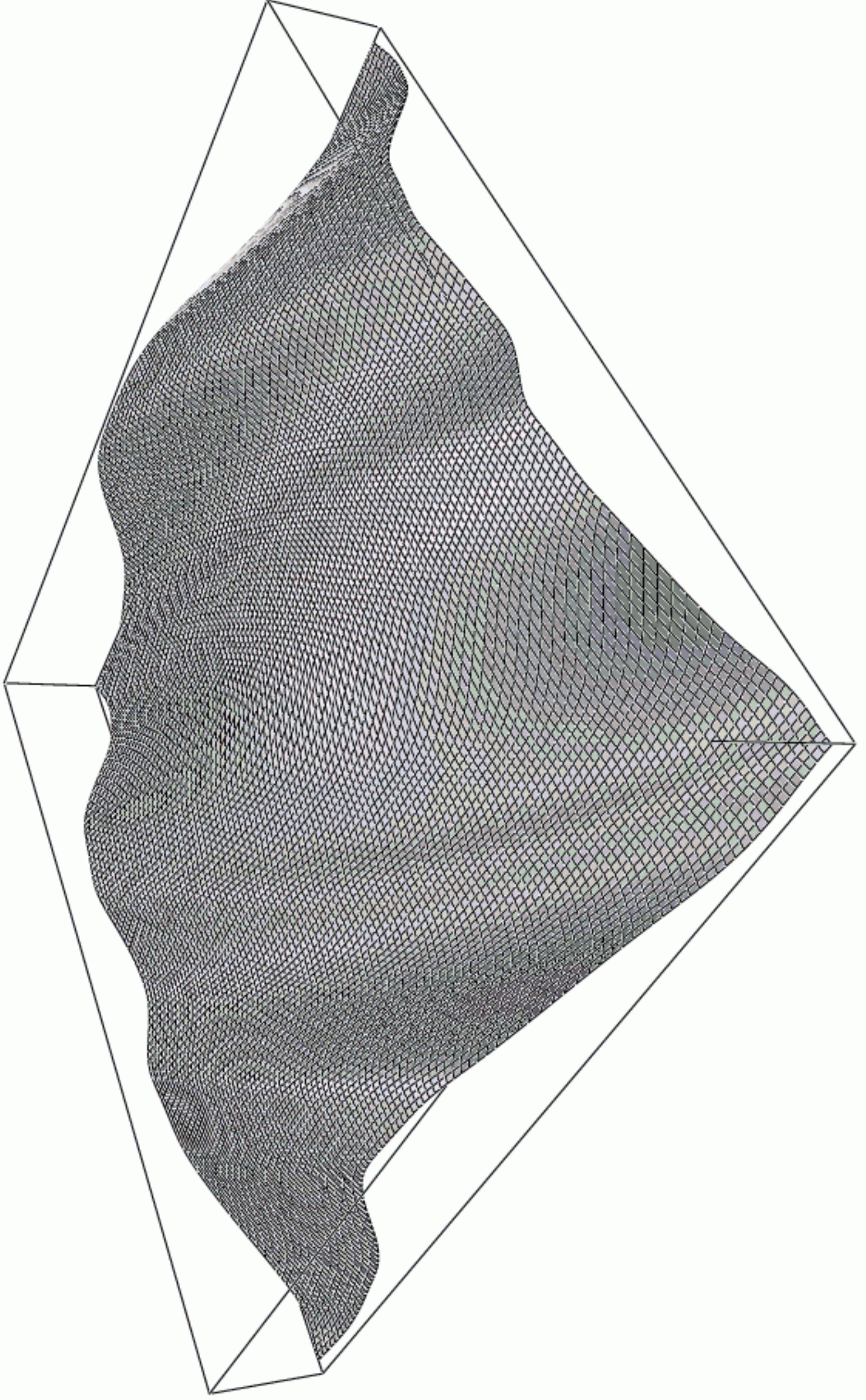
→ Get contraction in a space \mathcal{A}_r

$$\|\gamma\|_r = \sum_{\nu} |\gamma_{\nu}| e^{r|\omega \cdot \nu|} (1 + |\Omega \cdot \nu|)^{\nu}$$

with suff. small $r > 0$.

Starting with an approx. solution γ' ...

**Approximate numerical solution
of torus equation (*)**



Bounds on Λ and δ'
 that hold numerically
 and seem provable
 in the golden mean case

Theorem: If then F has a fixed
 point δ_* in A_r near δ' , and this
 fixed point defines an invariant
 ω -torus $\Gamma_* = I + \delta_*$ for H_* .
 If in addition, $-\nu^{-2}$ is not an eigenvalue
 of $D\Lambda(\Gamma(0))$, then δ_* is not C^2 .